

Trapped Ions and Atoms as Qubits

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National Security Agency



Advanced Research
& Development Activity



Army Research Office

NSF National Science Foundation



Frontiers of Optical
Coherent & Ultrafast Science

Quantum Information and Atomic Physics

$$H = \sum_{i=1}^N \frac{1}{2} \omega_i(t) \hat{\sigma}^{(i)} + \sum_{i,j=1}^N g_{ij}(t) \hat{\sigma}^{(i)} \cdot \hat{\sigma}^{(j)}$$

↑
N qubits

↑
controlled
coupling

... to >99% accuracy*

* provided things have been done right

PERIODIC TABLE

Atomic Properties of the Elements

U.S. DEPARTMENT OF COMMERCE
 Technology Administration
 National Institute of Standards and Technology

Frequently used fundamental physical constants	
For the most accurate values of these and other constants, visit physics.nist.gov/constants	
1 second = 9 192 631 770 periods of radiation corresponding to the transition between the two hyperfine levels of the ground state of ^{133}Cs	
speed of light in vacuum	c 299 792 458 m s ⁻¹ (exact)
Planck constant	h 6.6261 × 10 ⁻³⁴ J s ($h = h/2\pi$)
elementary charge	e 1.6022 × 10 ⁻¹⁹ C
electron mass	m_e 9.1094 × 10 ⁻³¹ kg
	$m_e c^2$ 0.5110 MeV
proton mass	m_p 1.6726 × 10 ⁻²⁷ kg
fine-structure constant	α 1/137.036
Rydberg constant	R_∞ 10 973 732 m ⁻¹
	$R_\infty c$ 3.2898 4 × 10 ¹⁵ Hz
	R_{hc} 13.6057 eV
Boltzmann constant	k 1.3807 × 10 ⁻²³ J K ⁻¹

Period	IA		IIA		IIIB	IVB	VB	VIB	VIIA	VIII	IB	IIB	VIIIA		IB		IIB	
	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18
1	1 H Hydrogen 1.00794	2 He Helium 4.00260																
2	3 Li Lithium 6.941	4 Be Beryllium 9.01218																
3	11 Na Sodium 22.98977	12 Mg Magnesium 24.3050																
4	19 K Potassium 39.0983	20 Ca Calcium 40.078	21 Sc Scandium 44.95591	22 Ti Titanium 47.867	23 V Vanadium 50.9415	24 Cr Chromium 51.9961	25 Mn Manganese 54.93805	26 Fe Iron 55.845	27 Co Cobalt 58.93320	28 Ni Nickel 58.6934	29 Cu Copper 63.546	30 Zn Zinc 65.39	31 Ga Gallium 69.723	32 Ge Germanium 72.61	33 As Arsenic 74.92160	34 Se Selenium 78.96	35 Br Bromine 79.904	36 Kr Krypton 83.80
5	37 Rb Rubidium 85.4678	38 Sr Strontium 87.62	39 Y Yttrium 88.90585	40 Zr Zirconium 91.224	41 Nb Niobium 92.90638	42 Mo Molybdenum 95.94	43 Tc Technetium (98)	44 Ru Ruthenium 101.07	45 Rh Rhodium 102.90550	46 Pd Palladium 106.42	47 Ag Silver 107.8682	48 Cd Cadmium 112.411	49 In Indium 114.818	50 Sn Tin 118.710	51 Sb Antimony 121.760	52 Te Tellurium 127.60	53 I Iodine 126.90447	54 Xe Xenon 131.29
6	55 Cs Cesium 132.90545	56 Ba Barium 137.327																
7	87 Fr Francium (223)	88 Ra Radium (226)																

Atomic Number: 58
 Ground-state Level: $1G_4$
 Symbol: **Ce**
 Name: Cerium
 Atomic Weight: 140.116
 Ground-state Configuration: $[Xe]4f5d6s^2$
 Ionization Energy (eV): 5.5387

- Solids
- Liquids
- Gases
- Artificially Prepared

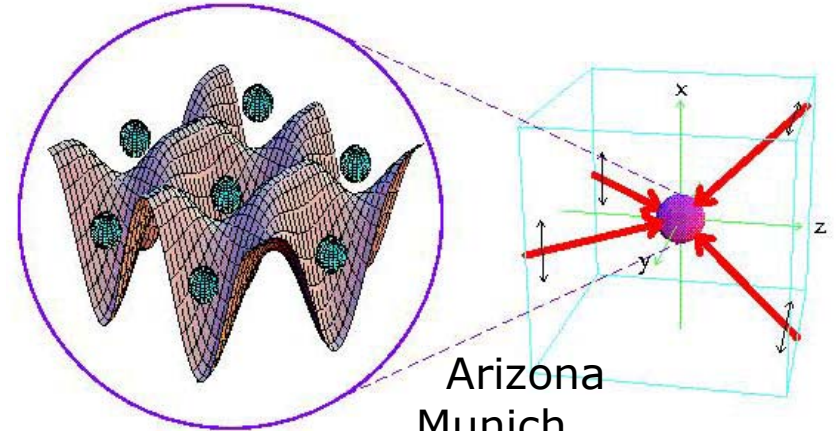
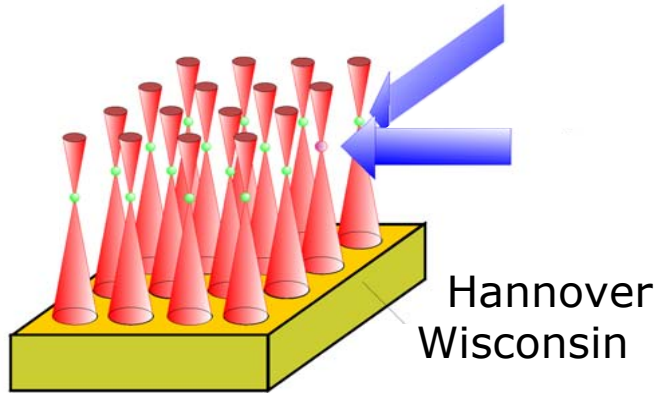
For a description of the atomic data, visit physics.nist.gov/atomic

57 La Lanthanum 138.9055	58 Ce Cerium 140.116	59 Pr Praseodymium 140.90765	60 Nd Neodymium 144.24	61 Pm Promethium (145)	62 Sm Samarium 150.36	63 Eu Europium 151.964	64 Gd Gadolinium 157.25	65 Tb Terbium 158.92534	66 Dy Dysprosium 162.50	67 Ho Holmium 164.93032	68 Er Erbium 167.26	69 Tm Thulium 168.93421	70 Yb Ytterbium 173.04	71 Lu Lutetium 174.967
89 Ac Actinium (227)	90 Th Thorium 232.0381	91 Pa Protactinium 231.03588	92 U Uranium 238.0289	93 Np Neptunium (237)	94 Pu Plutonium (244)	95 Am Americium (243)	96 Cm Curium (247)	97 Bk Berkelium (247)	98 Cf Californium (251)	99 Es Einsteinium (252)	100 Fm Fermium (257)	101 Md Mendelevium (258)	102 No Nobelium (259)	103 Lr Lawrencium (262)

Trapped Neutral Atoms

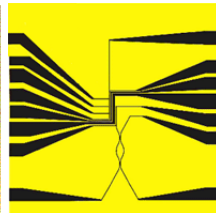
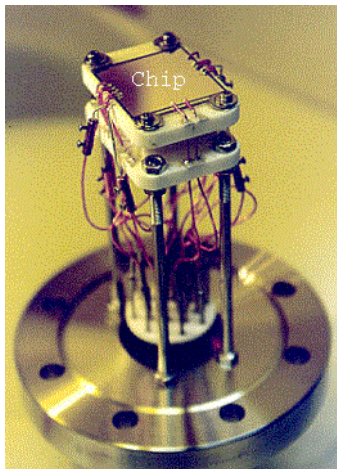
Atom arrangement/control

- optical lattices
- micro-magnetic traps
- dipole-trap arrays



Arizona
Munich
NIST-G'burg
Stanford
Yale

...



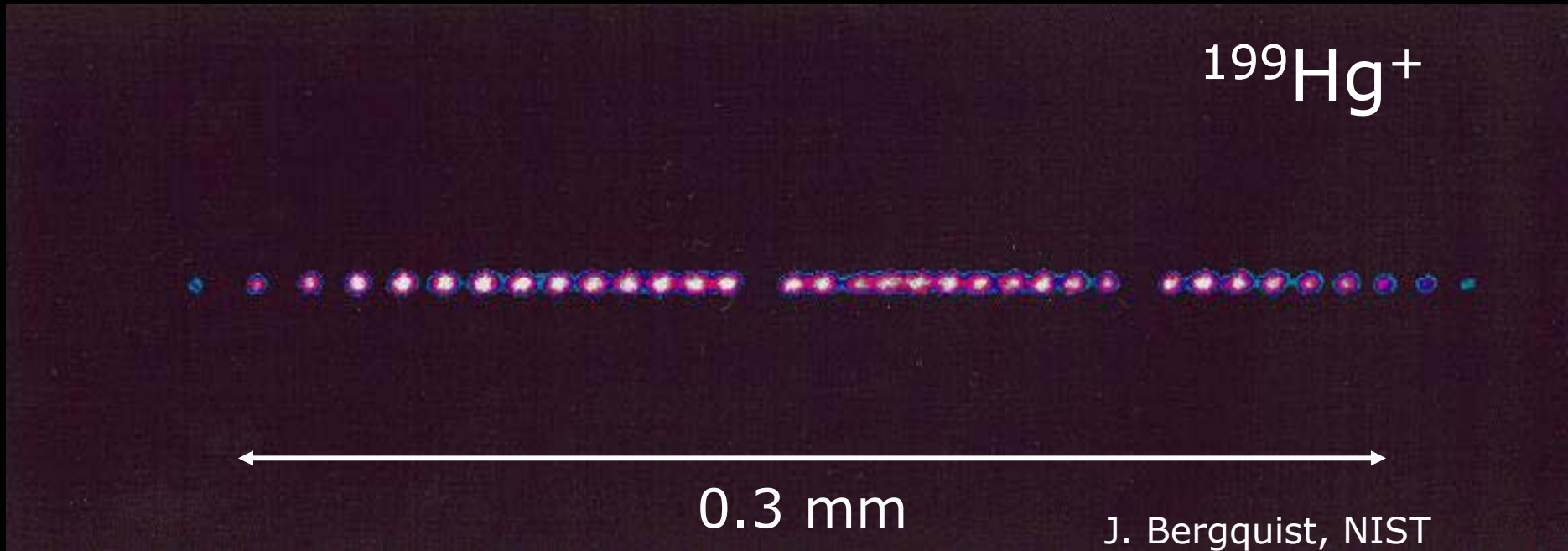
Heidelberg
JILA
Munich
MIT

...

Atom-atom interaction

- dipole-dipole
- s-wave collisions
- Rydberg atoms

Trapped Atomic Ions

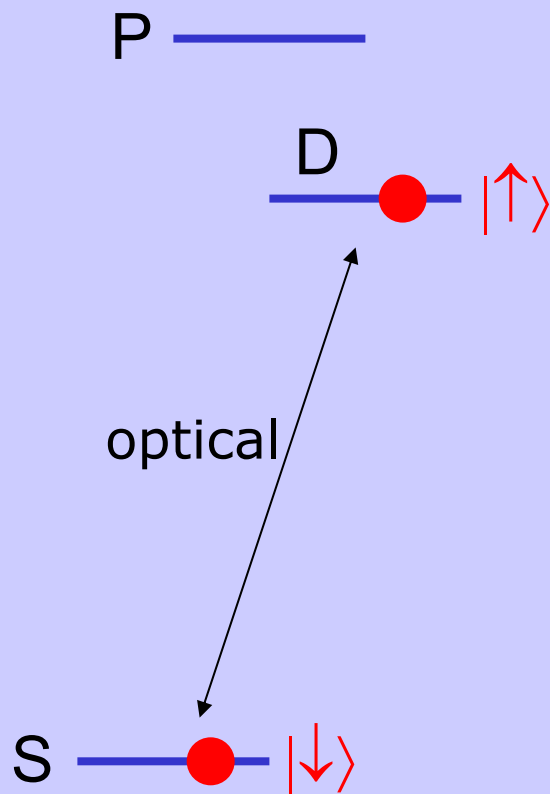


Ion Trap QC
Groups:

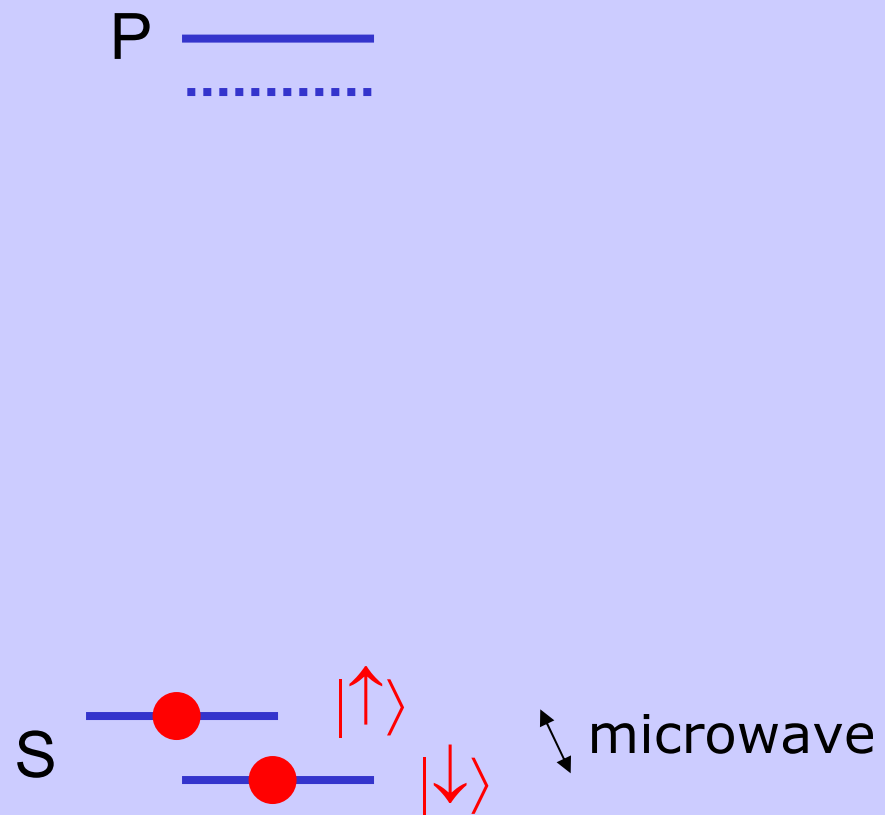
Aarhus
Alamaden (IBM)
Boulder (NIST)
Munich (MPQ)
Hamburg
Innsbruck

Los Alamos
McMaster (Ontario)
Michigan
Oxford
Teddington (NPL)

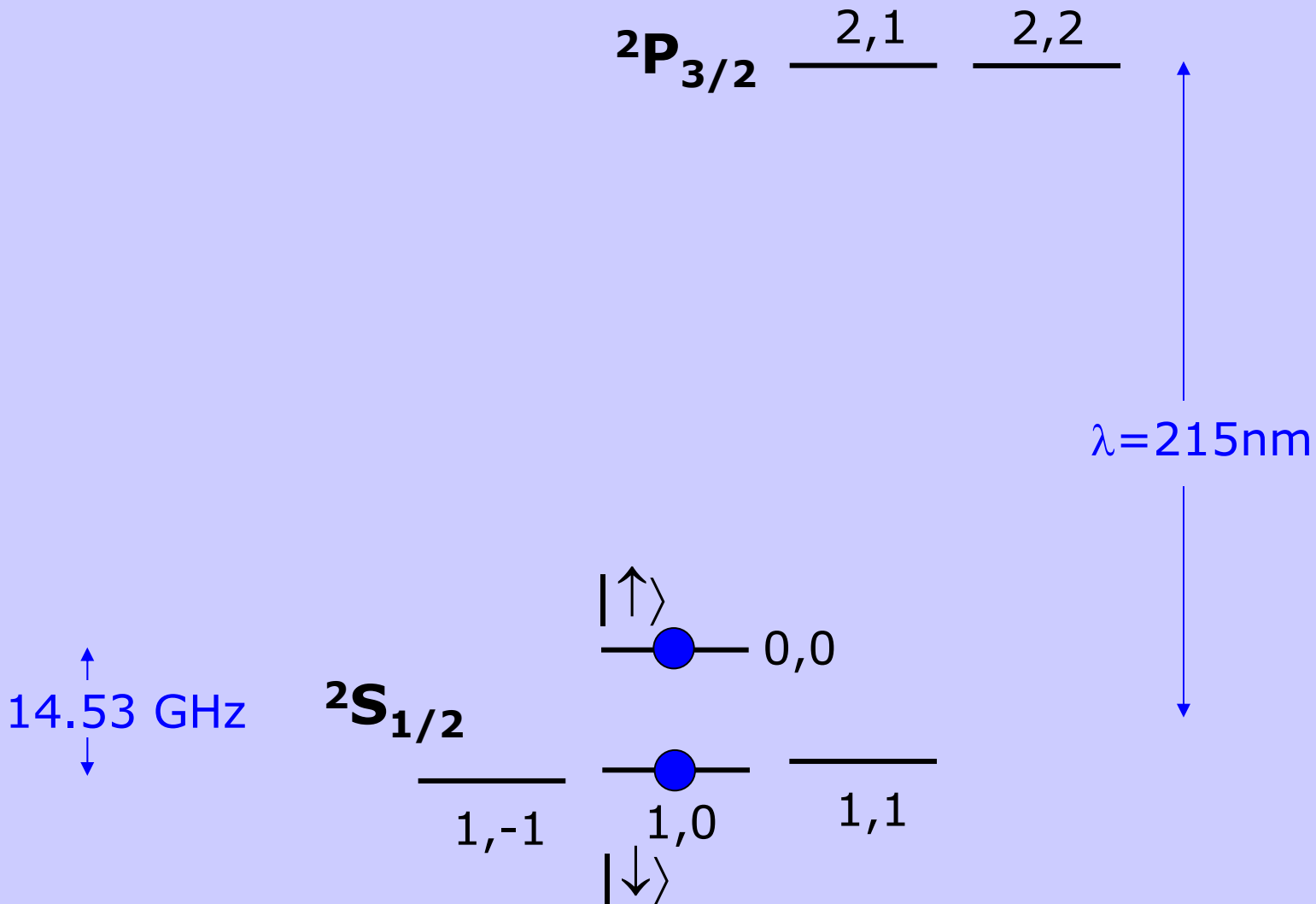
Ca⁺, Sr⁺, Ba⁺, Yb⁺



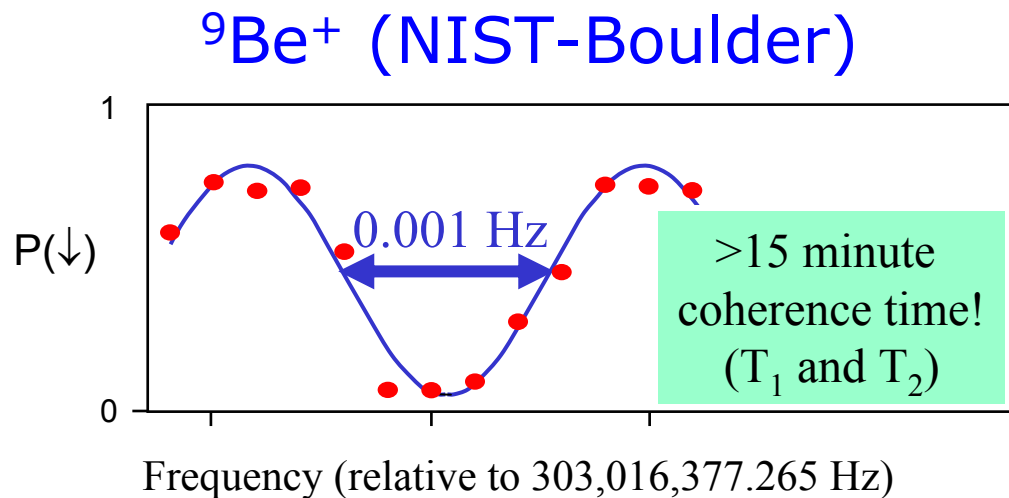
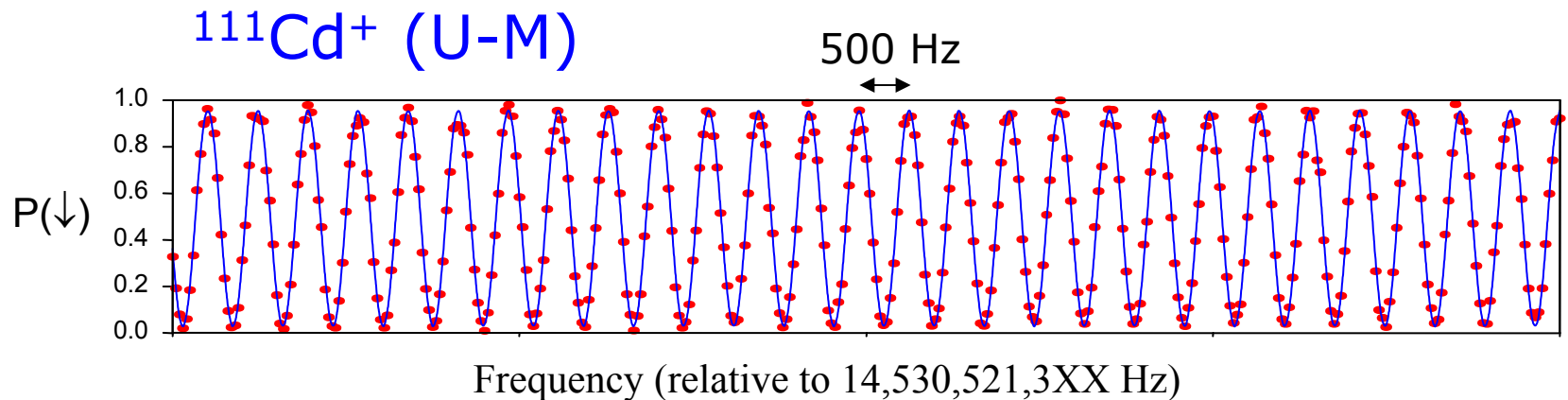
Be⁺, Mg⁺, Hg⁺, Cd⁺, Zn⁺



$^{111}\text{Cd}^+$ atomic structure

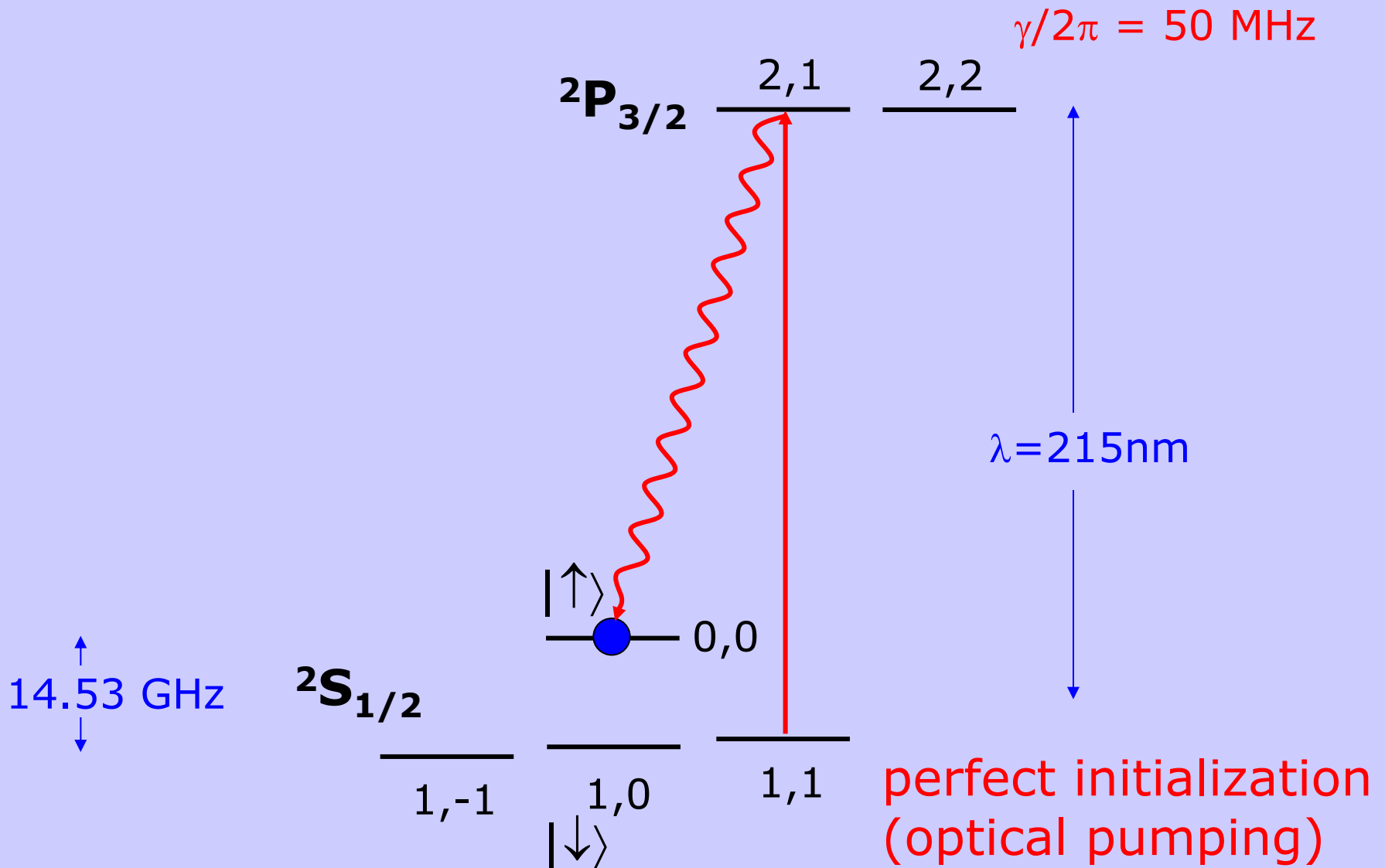


Ramsey interferometry with a trapped ion HF qubit: atomic clockwork

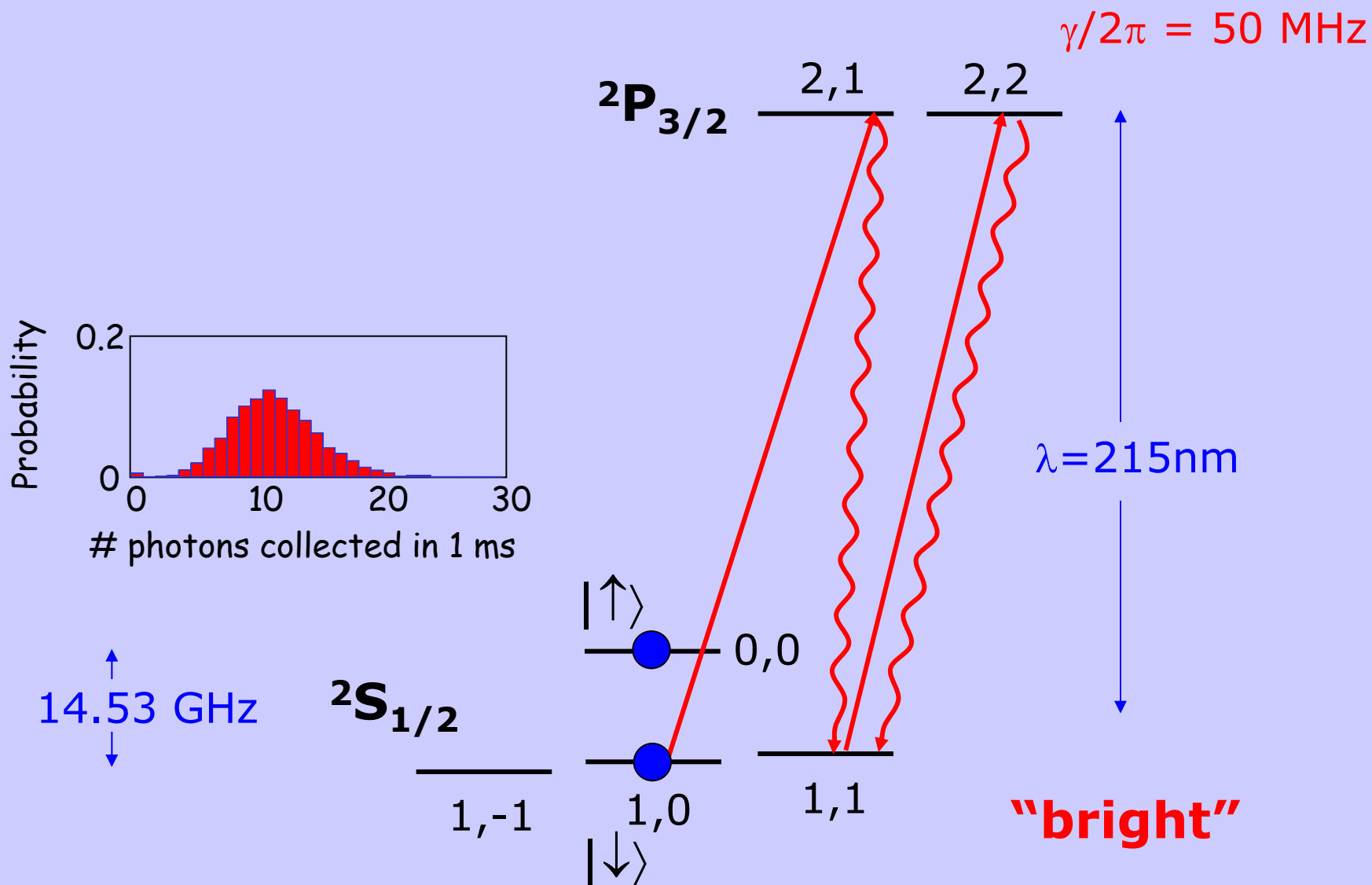


NIST: J. Bollinger, et. al., IEEE
Trans. Instrum. Meas. **40**, 126 (1991)

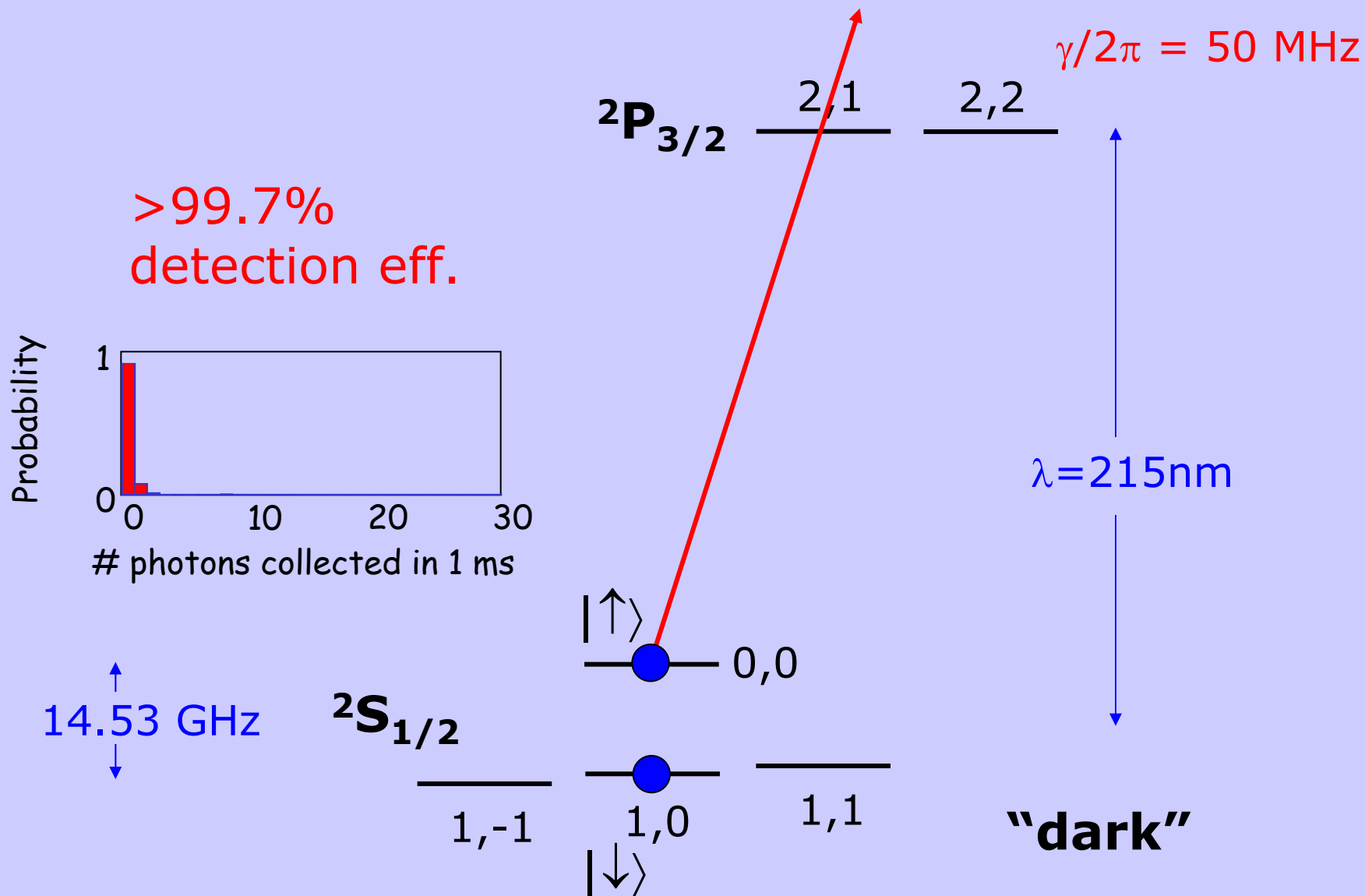
qubit initialization



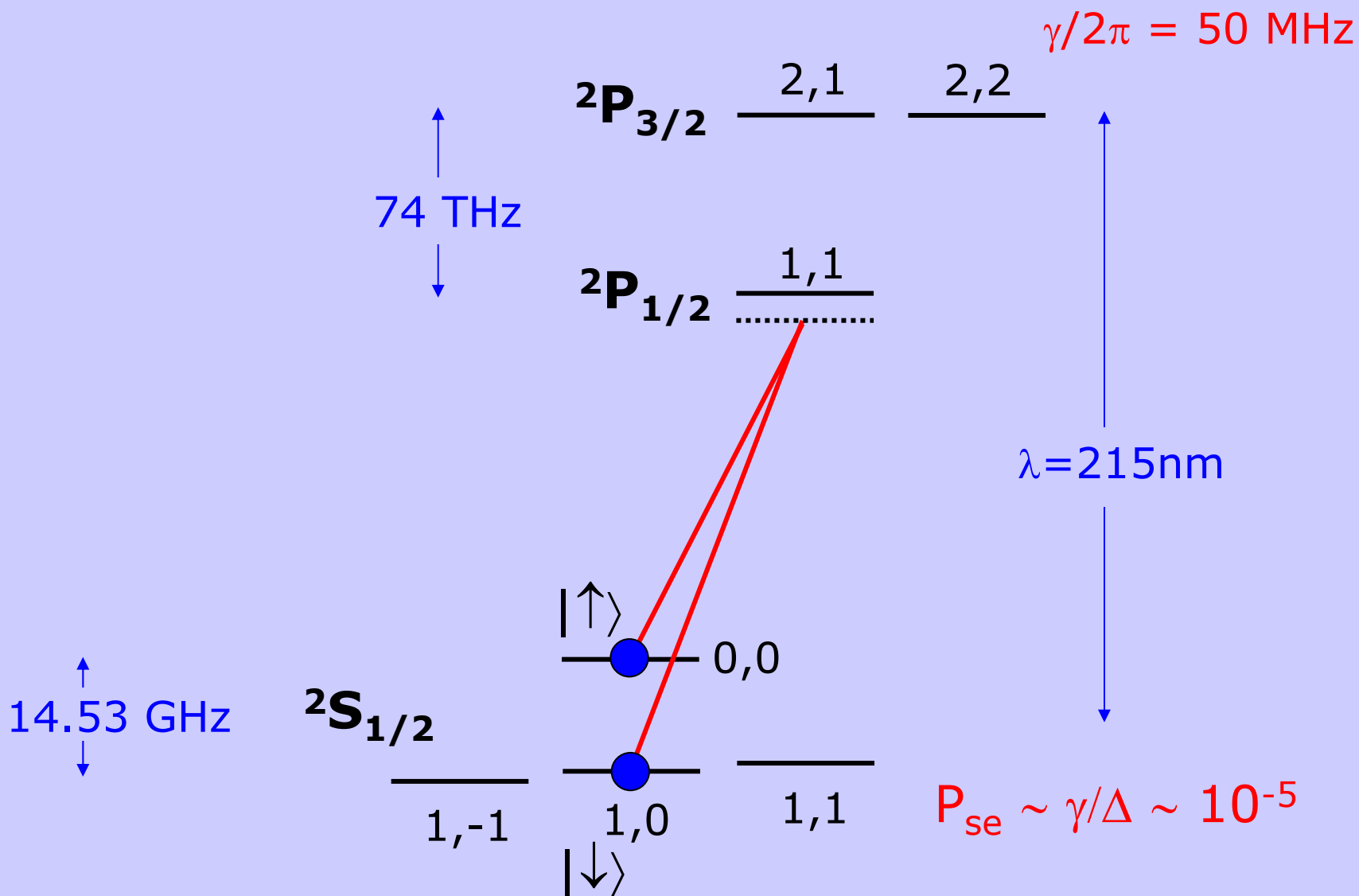
qubit measurement



qubit measurement

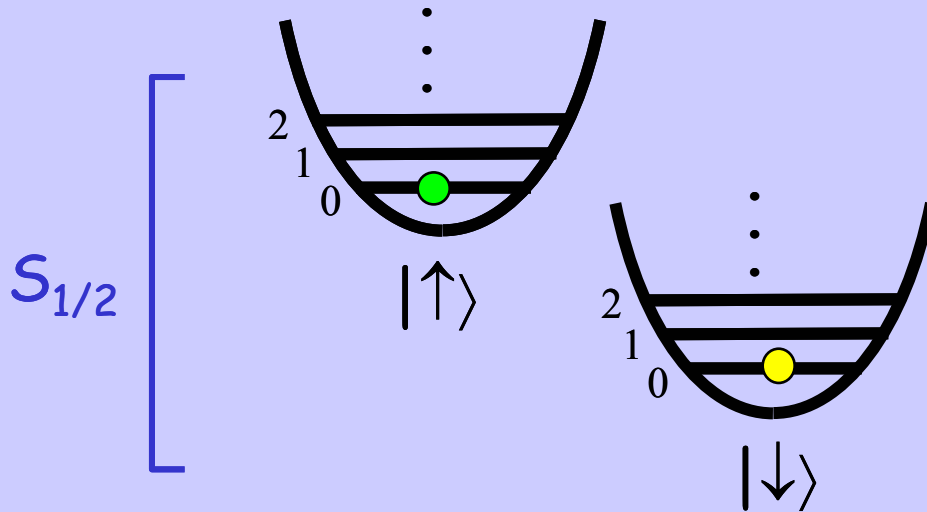


qubit manipulation



More ions: Use collective motion to entangle

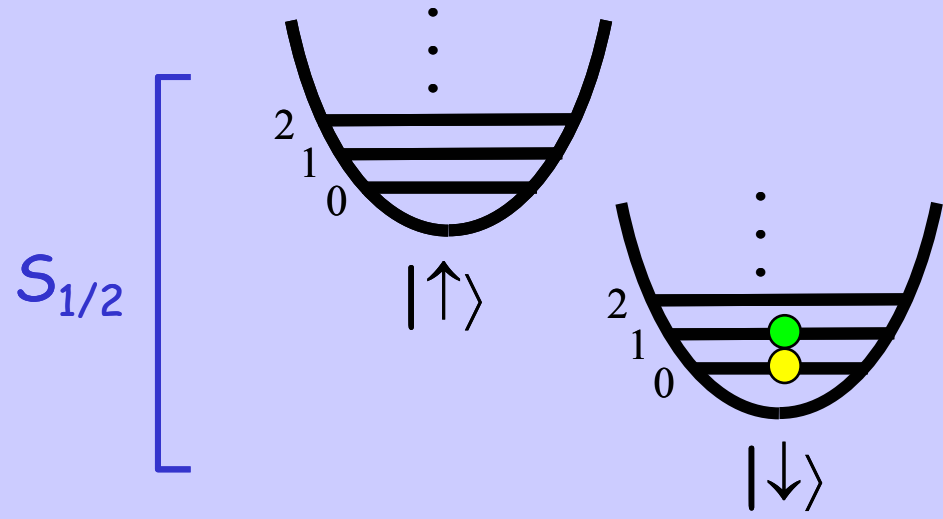
$P_{3/2}$



Mapping: $(\alpha|\downarrow\rangle + \beta|\uparrow\rangle) |0\rangle_m \rightarrow |\downarrow\rangle (\alpha|0\rangle_m + \beta|1\rangle_m)$

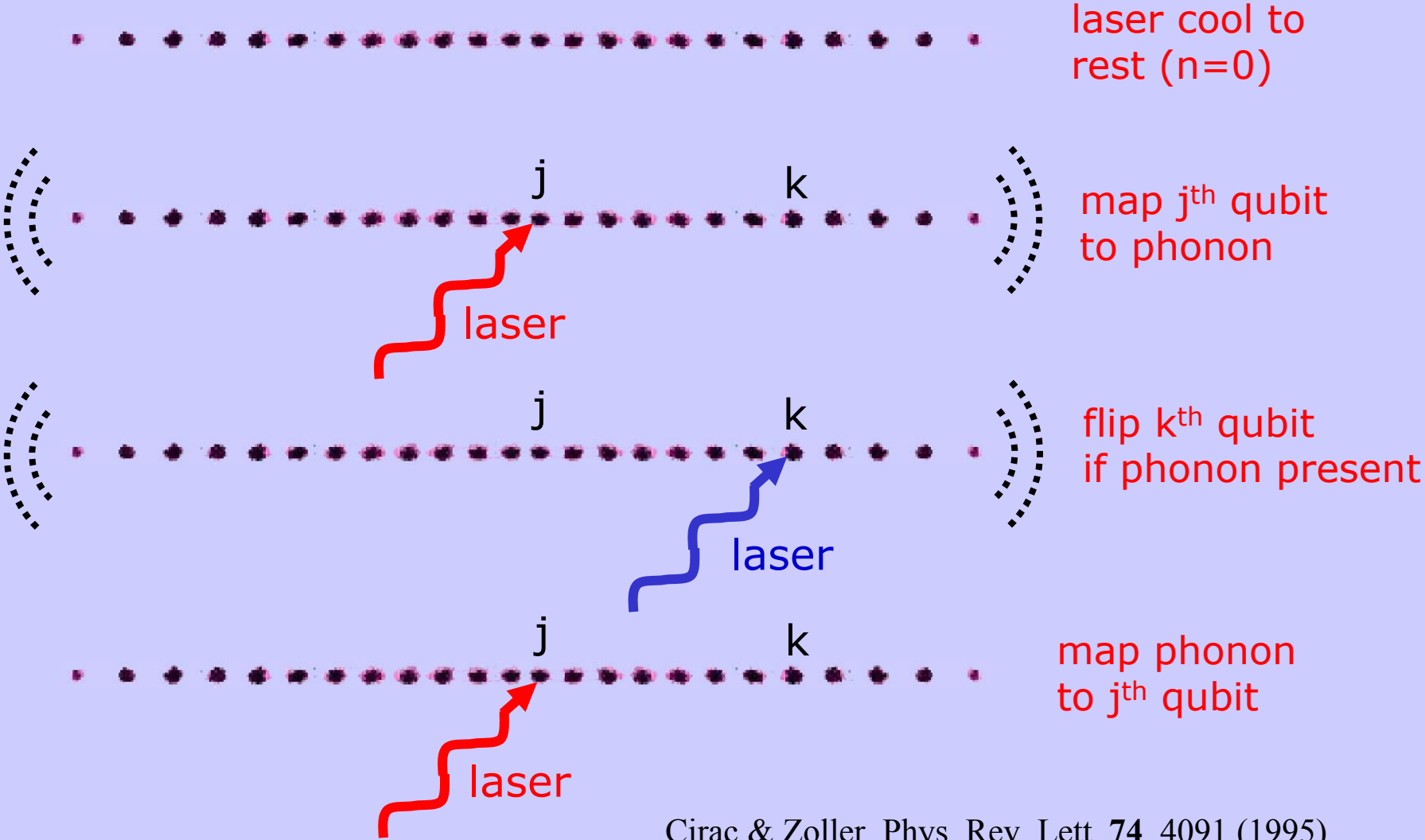
More ions: Use collective motion to entangle

$P_{3/2}$



Mapping: $(\alpha|\downarrow\rangle + \beta|\uparrow\rangle) |0\rangle_m \rightarrow |\downarrow\rangle (\alpha|0\rangle_m + \beta|1\rangle_m)$

Entangling Ions with Collective Phonons



Cirac & Zoller, Phys. Rev. Lett. 74, 4091 (1995)

Cirac-Zoller Scheme

NIST(1995): $N=1$ ion, $F=85\%$

Innsbruck (2003): $N=2$ ions, $F=71\%$

Major improvements:

- Don't need to individually address ions to entangle them
- Don't need a "pure" state of motion

Thy: Molmer & Sorensen (1999)

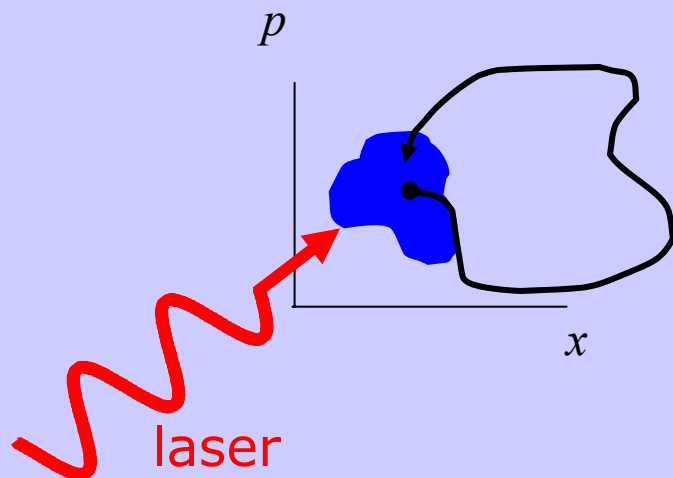
Milburn, Schneider, James (1999)

Exp: NIST (2000): $N=4$ ions, $F=57\%$

NIST (2003): $N=2$ ions, $F=97\%$

“State-dependent force” gates

N=1 ion: Force = $F_0 |\uparrow\rangle\langle\uparrow|$



$$\begin{aligned} \downarrow &\Rightarrow \downarrow \\ \uparrow &\Rightarrow e^{i\phi} \uparrow \end{aligned}$$

($\phi \propto$ enclosed area)

N=2 ions

e.g., force on stretch mode only

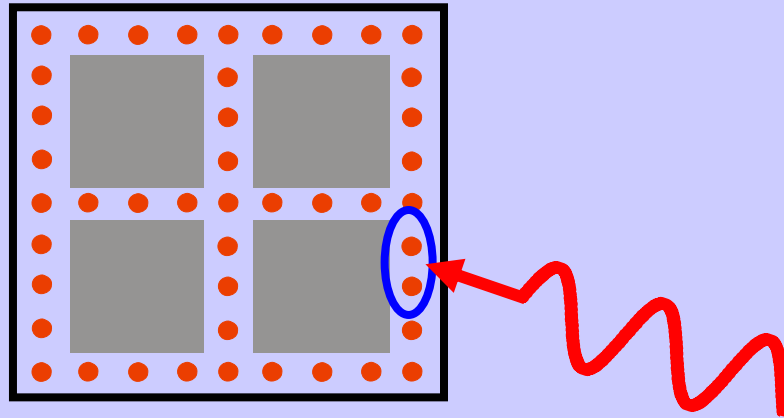


$$\begin{aligned} \downarrow\downarrow &\Rightarrow \downarrow\downarrow \\ \downarrow\uparrow &\Rightarrow e^{i\phi} \downarrow\uparrow \\ \uparrow\downarrow &\Rightarrow e^{i\phi} \uparrow\downarrow \\ \uparrow\uparrow &\Rightarrow \uparrow\uparrow \end{aligned}$$

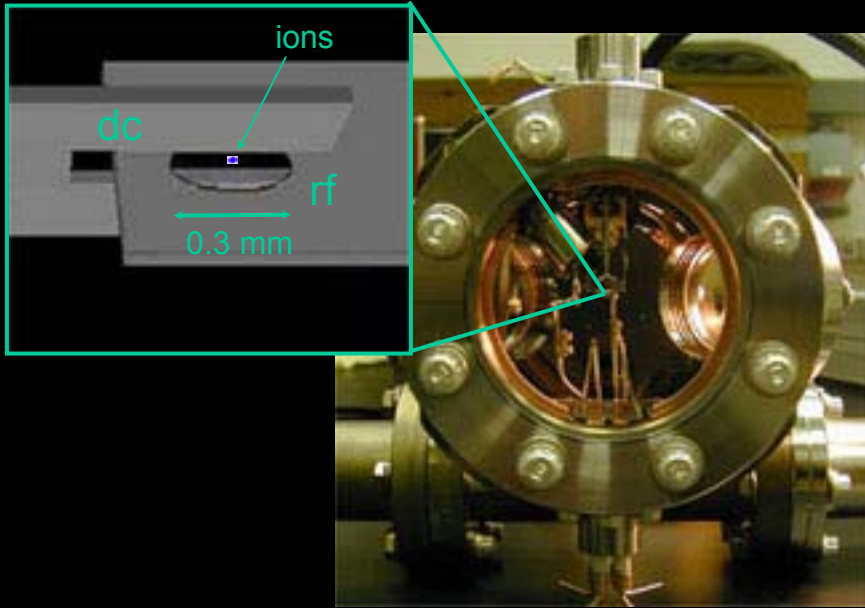
$\phi = \pi/2$: π -phase gate
 NIST (2003): 97% Fidelity

Ultrafast force gates (theoretical):

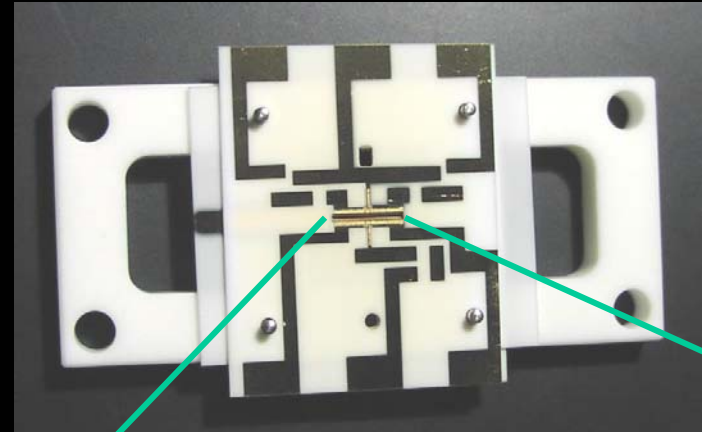
- Don't need $\Delta x < \lambda$ (Lamb-Dicke confinement not required)
- Not speed-limited by motional trap frequency
Thy: Garcia-Ripoll, Zoller, Cirac (2003)
- OK in the presence of other ions
Thy: Duan (2004)



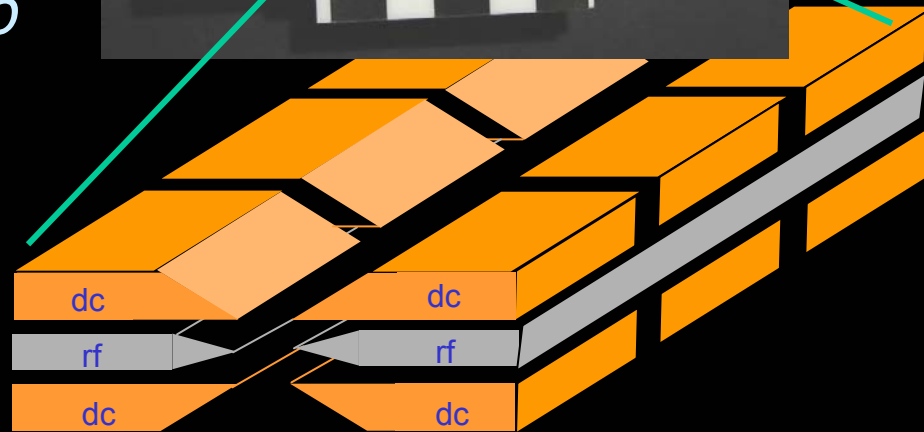
Exp: **Need good pulsed lasers!**
(eg., 1 GHz rep rate, $\tau=1$ psec)



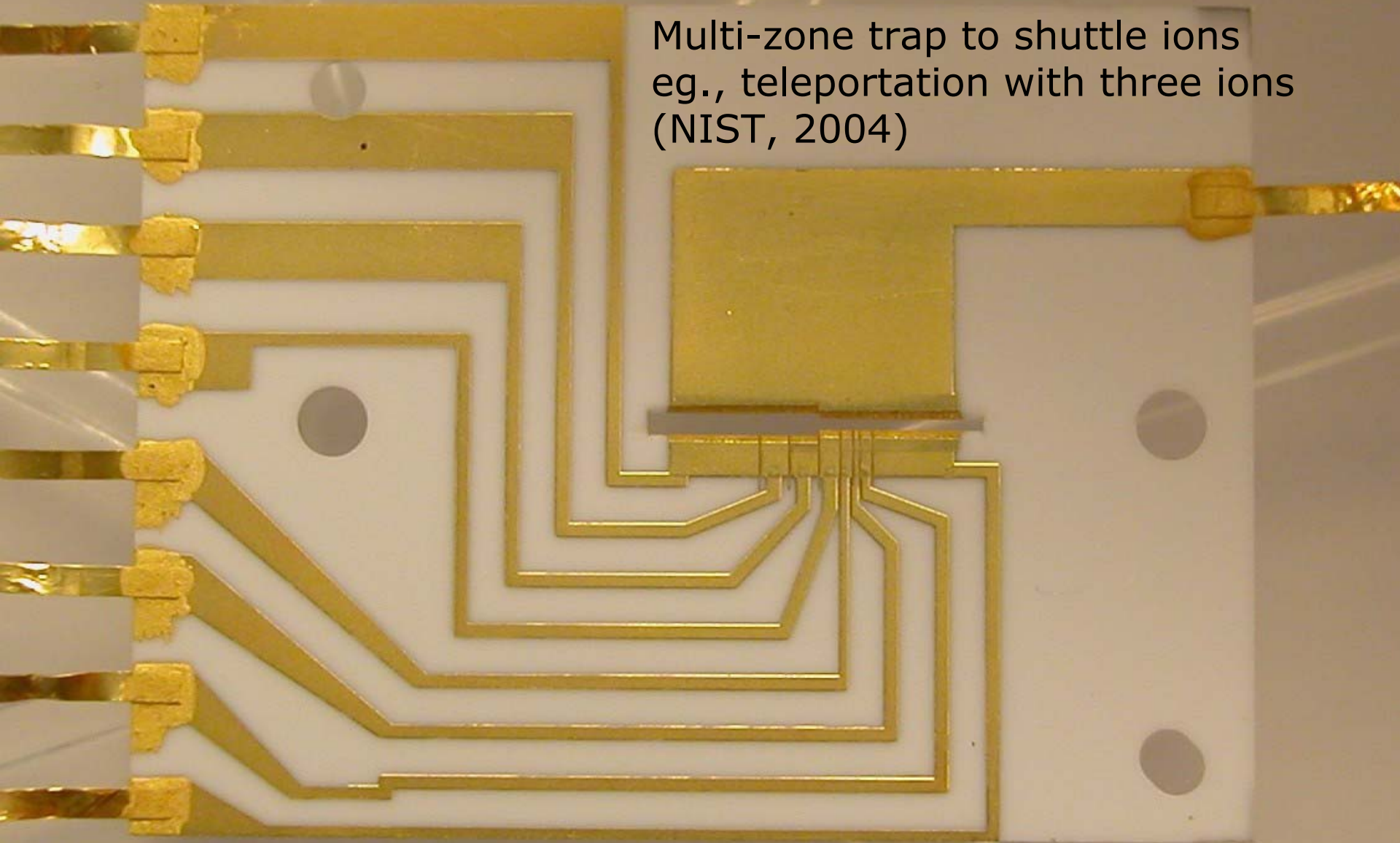
quadrupole rf trap



linear rf trap

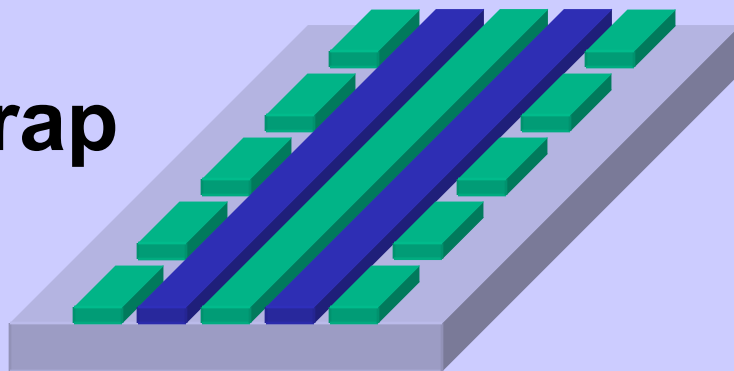


Multi-zone trap to shuttle ions
eg., teleportation with three ions
(NIST, 2004)

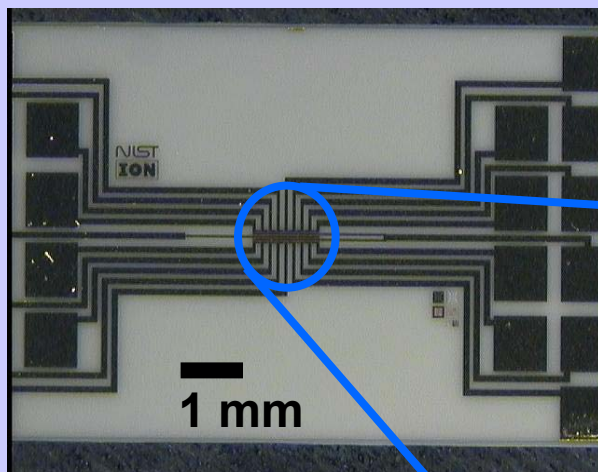
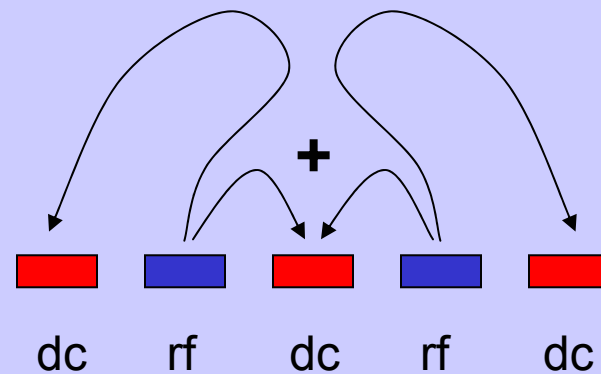


New Trapology (J. Chiaverini, NIST)

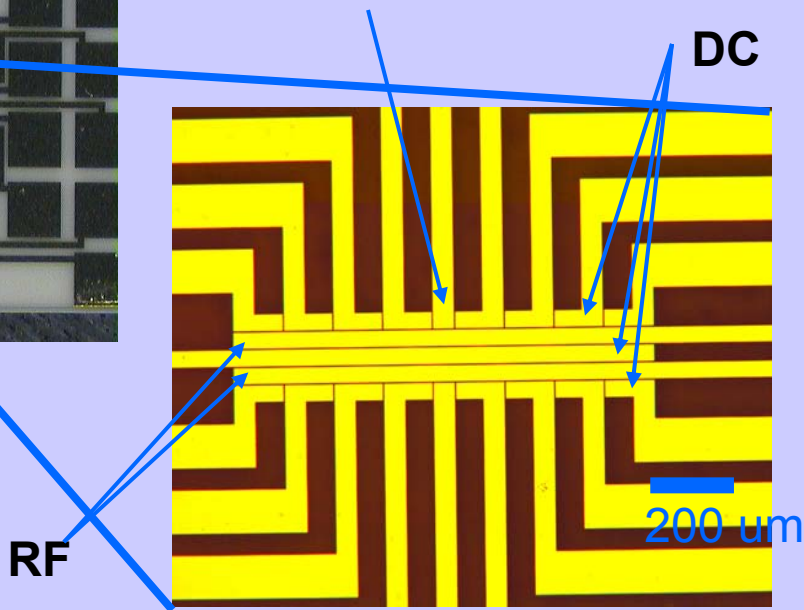
Planar trap



Field lines:



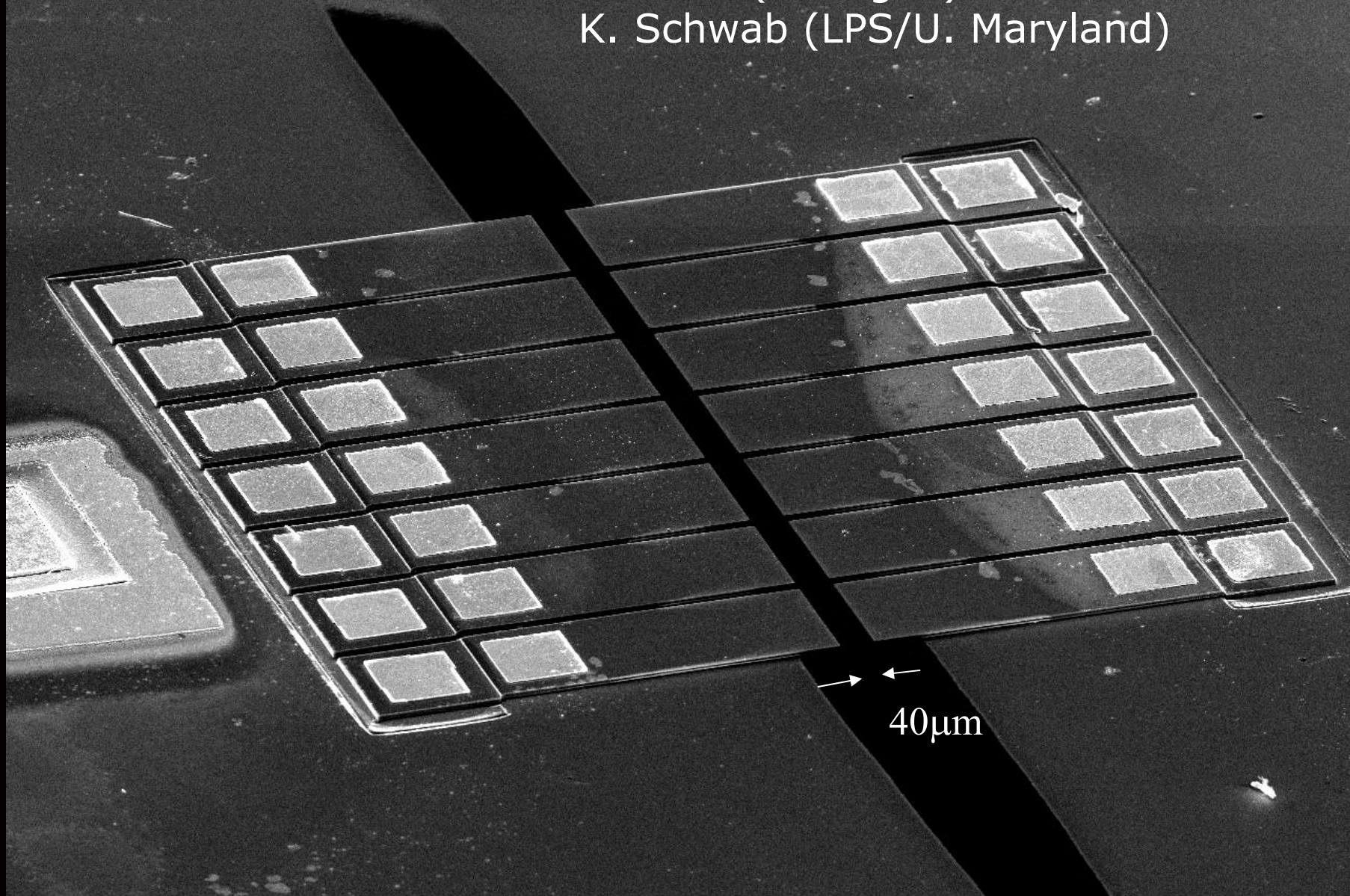
Splitting electrode



GaAs/AlGaAs

D. Stick (Michigan)

K. Schwab (LPS/U. Maryland)



40 μ m

LPS

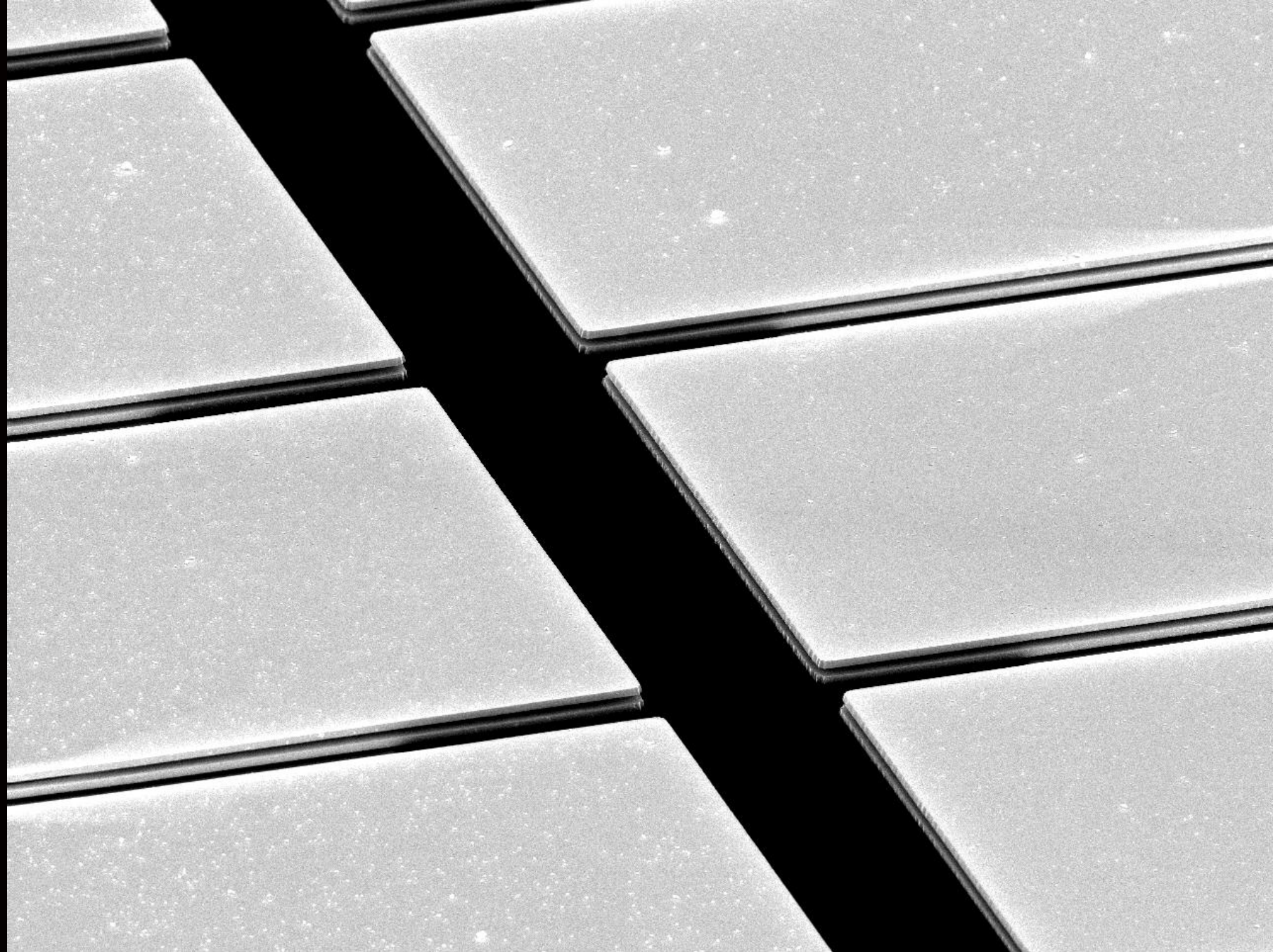
SEI

30.0kV

X80

100 μ m

WD 29.2mm



LPS

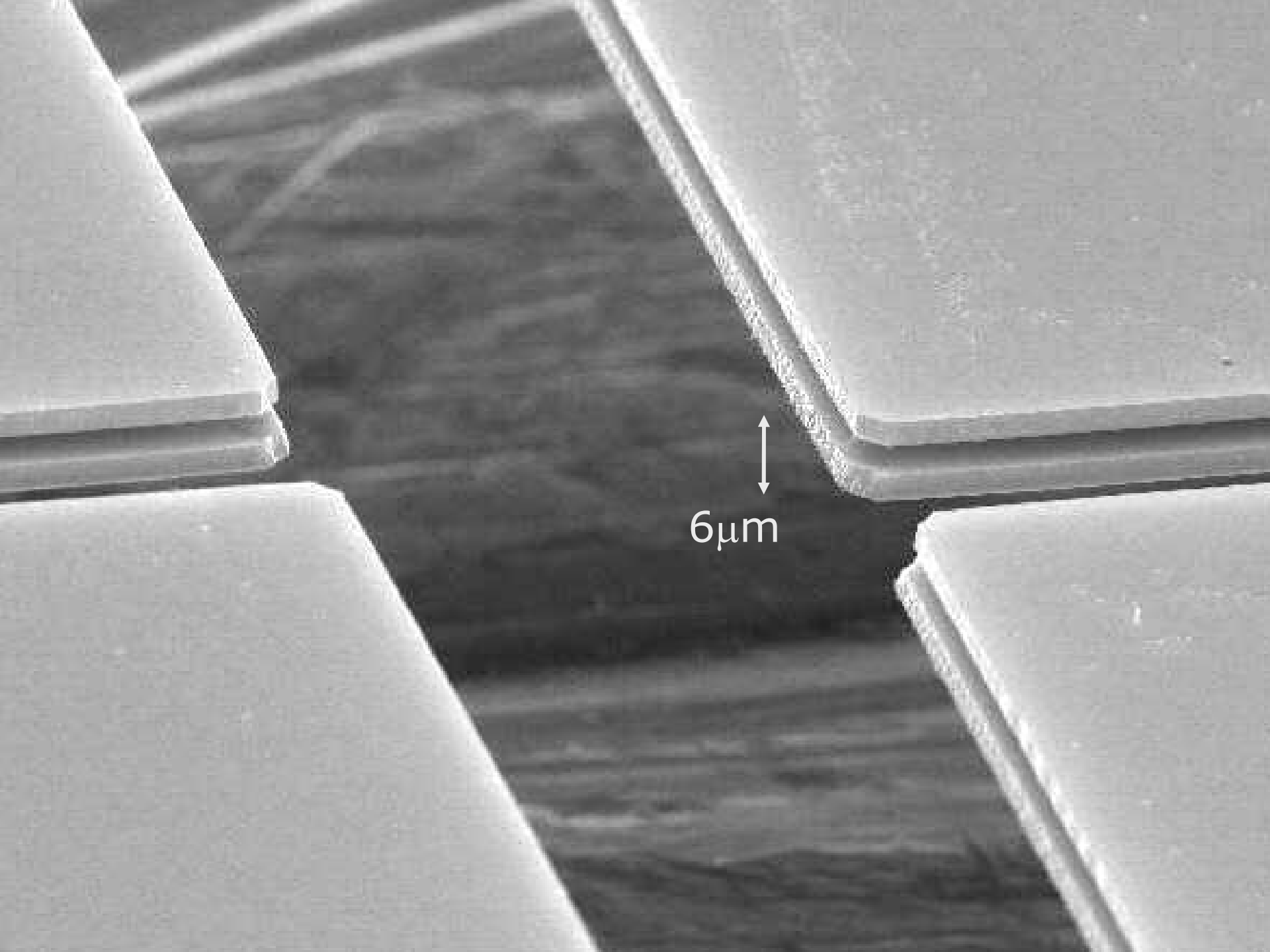
SEI

30.0kV

X450

10 μ m

WD 29.2mm



6 μm

Using a photon as the data bus

“DiVincenzo 6,7”

cavity-QED (trap the photon too)

CalTech

ENS-Paris

MPQ-Garching

...

Ensemble spin-squeezing

Copenhagen

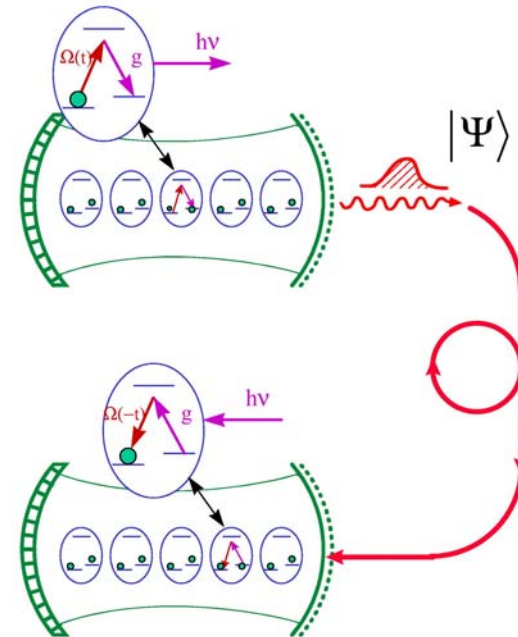
Rochester

Harvard

...

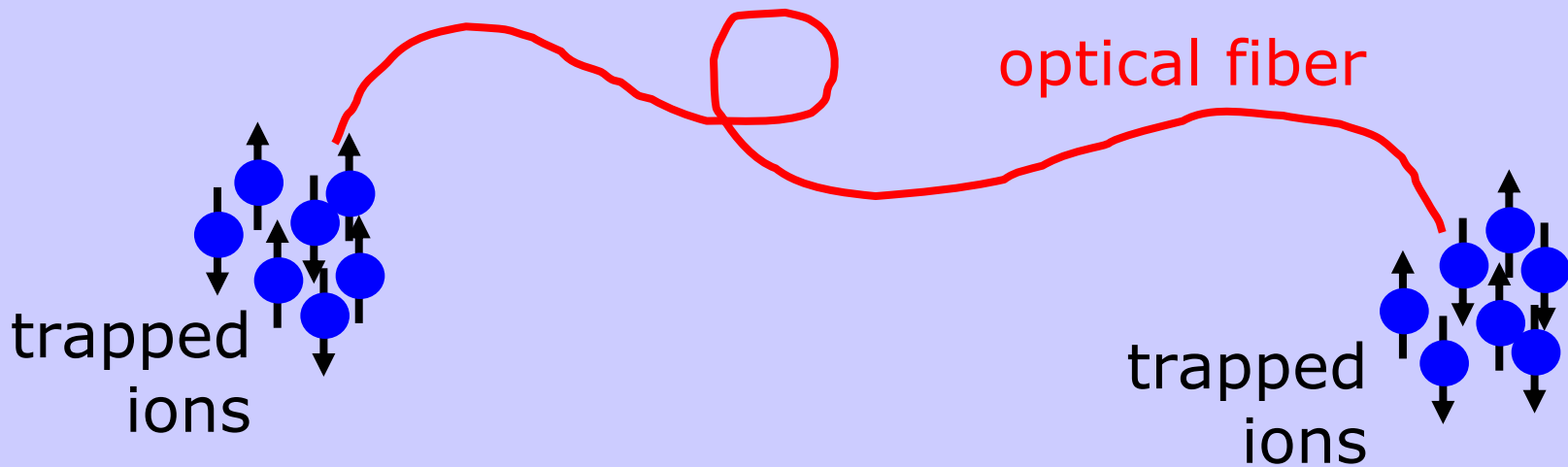
entanglement of atom & photon:

must have excellent control of **both** atom and photon

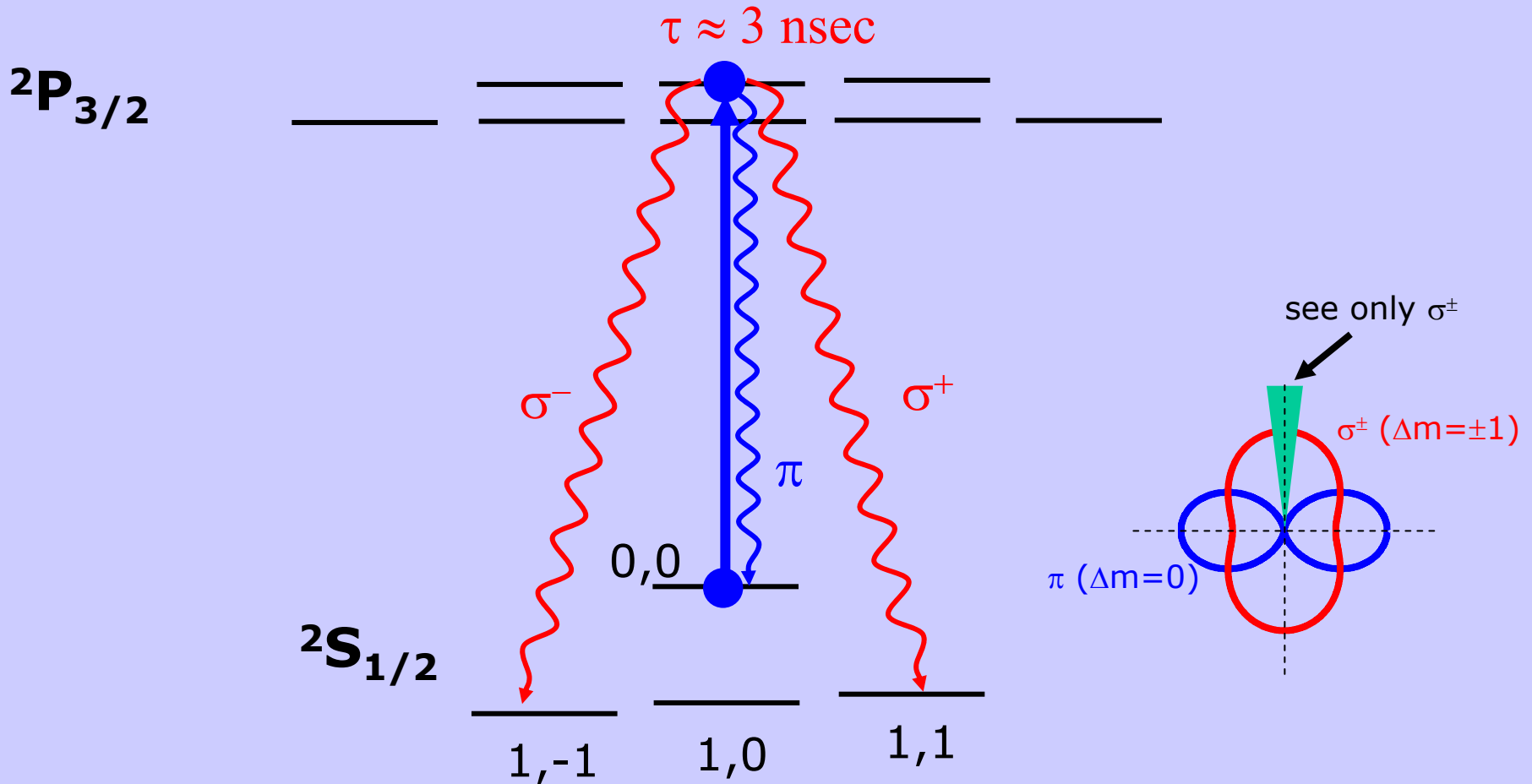


Ideal quantum memory
trapped atomic ion

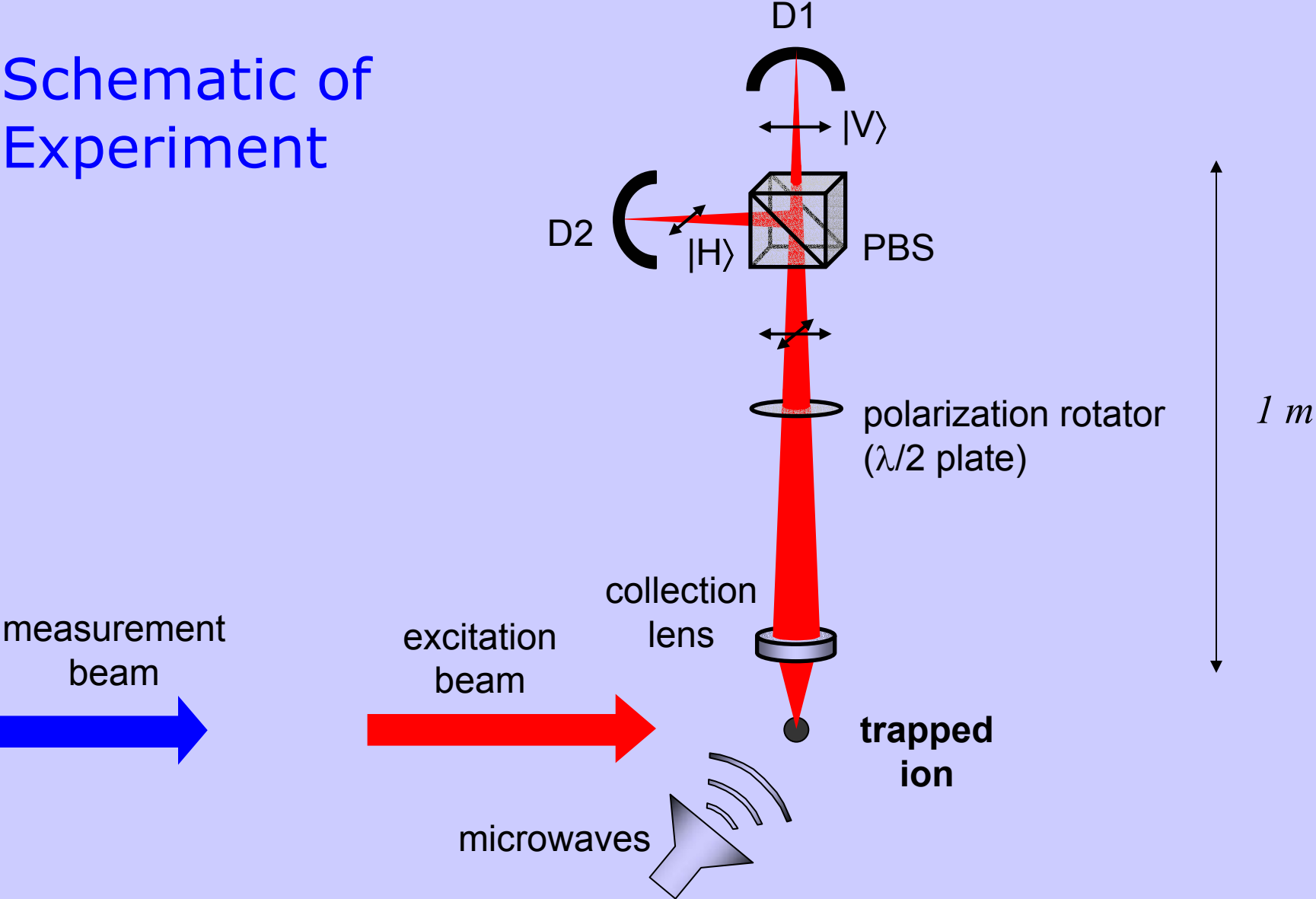
Quantum communication channel
photon: "flying qubit"



Probabilistic entanglement between a single atom and single photon

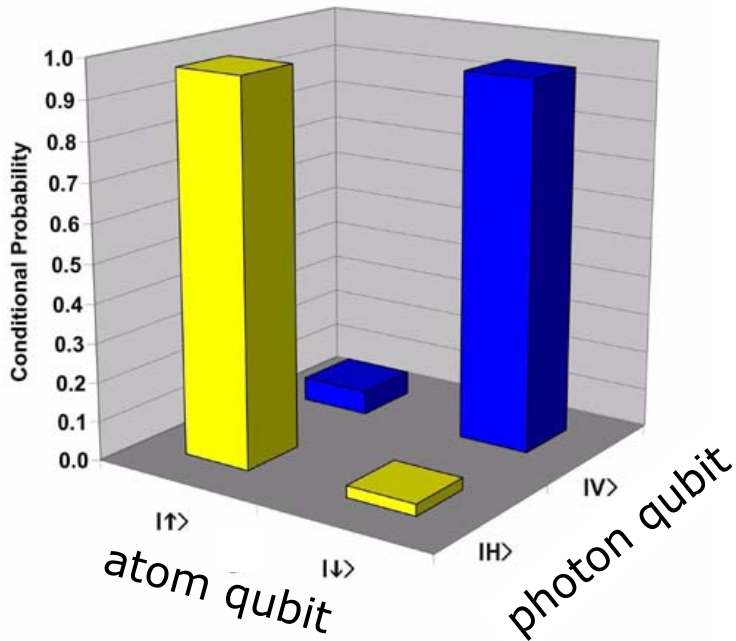


Schematic of Experiment

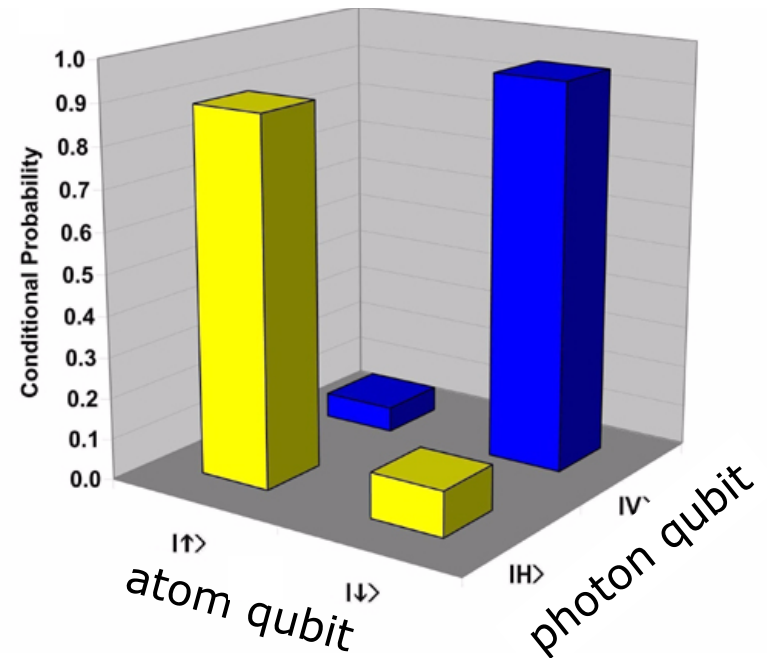


$$|\Psi\rangle_{\text{ideal}} = |H\rangle|\uparrow\rangle + |V\rangle|\downarrow\rangle$$

Measured correlation
between atomic and
photonic qubits...

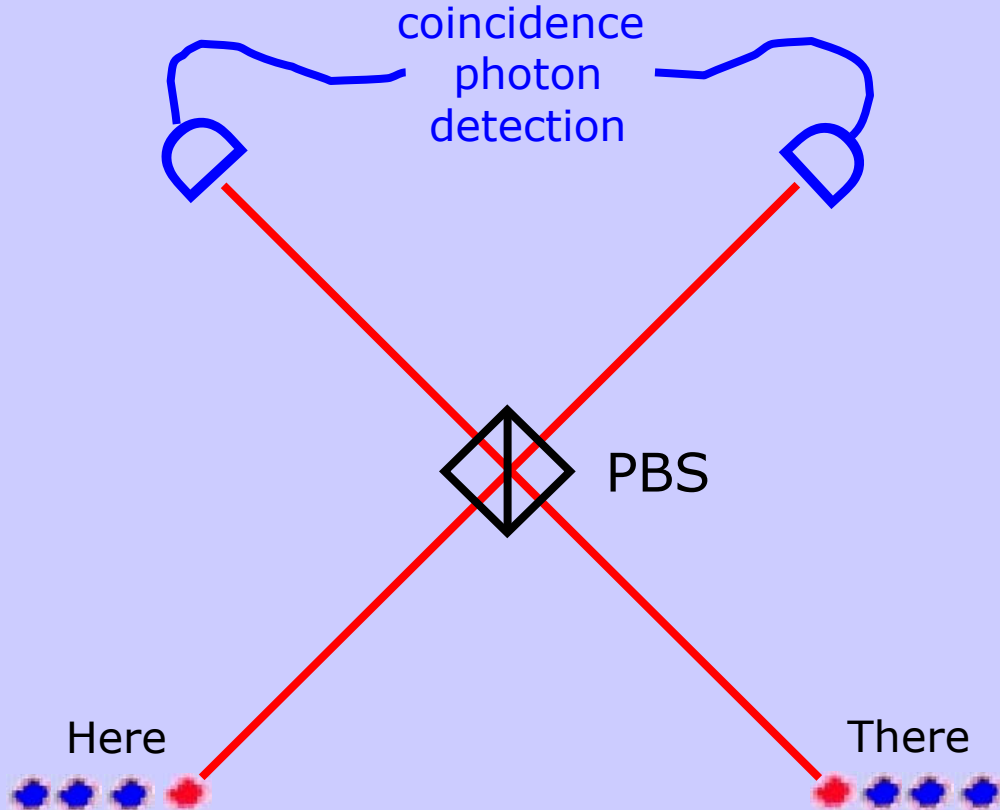


... and after rotation
of each qubit by $\pi/2$
before measurement



Entanglement Fidelity: $F = \langle \Psi_{\text{ideal}} | \rho | \Psi_{\text{ideal}} \rangle > 87\%$

Can use this technique to seed
remote ion-ion entanglement...



Speed of gate:

$$R = \Gamma(P_{\text{succ}})^2$$

$$= 0.01/\text{sec} \text{ now}$$

$$= 10^6/\text{sec} \text{ possible}$$

... and form the basis for scalable QC

New Frontiers in Quantum Information With Atoms and Ions

Both the precision control of trapped-ion systems and very large samples of cold neutral atoms are opening important new possibilities for quantum computation and simulation.

J. Ignacio Cirac and Peter Zoller

The success story of quantum optics during the past 10 years is largely based on progress in gaining control of systems at the single-quantum level while suppressing unwanted interactions with the environment, which cause decoherence. Those achievements, illustrated by storage and laser cooling of single trapped ions and atoms and by the manipulation of single photons in cavity quantum electrodynamics, have opened a new field: the engineering of interesting and useful quantum states. In the meantime, the frontier has moved toward building larger composite systems of a few atoms and photons while still maintaining complete quantum control of the individual particles. The new physics to be studied in these systems is based on entangled states and ranges from a fundamental point of testing quantum mechanics for larger and larger systems to possible new applications such as quantum information processing and precision measurements.^{1,2}

The past few years have seen extraordinary progress in experimental atomic, molecular, and optical (AMO) physics. Two highlights of those developments are laser-cooled trapped ions³⁻⁶ and cold atoms in optical lattices.⁹⁻¹¹ These two examples also illustrate the different perspectives and strengths of AMO systems. Systems of a few trapped ions have demonstrated quantum-entanglement engineering with high fidelity (that is, low error rate) in the laboratory, and these systems are well on their way toward scalable quantum computing (see box 1), with no fundamental obstacles in sight—at least from our current understanding. Neutral atoms can be loaded from a Bose-Einstein condensate (BEC) into an optical lattice via a quantum phase transition and can provide a huge number of qubits that can be entangled in massively parallel operations. Such a system holds the promise of a quantum simulator (see box 2) that may offer insight into other fields of physics, such as condensed matter physics.

Although we focus on these two AMO systems in this

during recent years. Those systems include single photons, nuclear spins of donor atoms in doped silicon, superconducting Josephson junctions in both the charge- and flux-quantization regimes, semiconductor quantum dots, nuclear magnetic resonance samples, and electrons floating on liquid helium. Some of the ideas we re-

view here will likely apply to these systems if they ultimately succeed as quantum computers.

Cold trapped ions

Right after Peter Shor's discovery in 1994 of a factoring algorithm for quantum computers¹ (see PHYSICS TODAY, October 1995, page 24), trapped ions interacting with laser light were identified as one of the most promising candidates to build a small-scale quantum computer.³ The reason is that, for many years, the technology to control and manipulate single (or few) ions had been very strongly developed for ultrahigh-precision spectroscopy and atomic clocks.¹² In particular, ions can be trapped and cooled in such a way that they remain practically frozen in a specific region of space; their internal states can be precisely manipulated using lasers and can be measured with practically 100% efficiency; and they interact with each other very strongly due to the Coulomb repulsion, yet they can, at the same time, be decoupled from the environment very efficiently.

Ions stored and laser-cooled in an electromagnetic trap (see figure 1) can be described in terms of a set of external and internal degrees of freedom. The external degrees of freedom are closely related to the center-of-mass motion of each ion; the internal, related to the motion of electrons within each ion and to the presence of electronic and nuclear spins, are responsible for the existence of a discrete energy-level structure in each ion. Each qubit can be stored in two of the internal levels, typically denoted by $|0\rangle$ and $|1\rangle$. These levels have to be very long-lived and suffer no decoherence, so that they are not disturbed during the computation. That condition can be achieved, for example, by choosing them as ground hyperfine or metastable Zeeman levels, where spontaneous emission is practically absent.

To start a computation, one can prepare all the qubits

Physics Today
March, 2004

University of Michigan

Trapped Ion Quantum Computing

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Chitra Rangan (postdoc/theory)

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<http://iontrap.physics.lsa.umich.edu/>