

## Basics of Two-Fluid Plasma Physics—a Short Summary

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The elementary building blocks for a multi-fluid are the canonical momentum  $\mathbf{P}_\alpha = m_\alpha \mathbf{u}_\alpha + q_\alpha \mathbf{A}/c$ , and the generalized vorticity  $\mathbf{\Omega}_\alpha = \nabla \times \mathbf{P}_\alpha$  (or  $\alpha$ -vorticity) where  $m_\alpha$ ,  $\mathbf{u}_\alpha$ ,  $q_\alpha$  are the species mass, flow velocity and charge, and  $\alpha = i, e$  denotes the species, and  $\mathbf{A}$  is the vector potential. The quadratic invariant of a species is the self helicity, or “ $\alpha$ -helicity,” the “density” of which is  $\mathbf{P}_\alpha \cdot \mathbf{\Omega}_\alpha$ . These are generalizations of helicities that appear in a simple fluids and MHD. For zero electron mass the electron helicity reduces to the familiar magnetic helicity, an invariant in ideal MHD. The evolution of the  $\alpha$ -helicities is governed by the helicity transport equation, derived from Maxwell’s equations and the equations of motion.

Each helicity transport equation has the form,  $n_\alpha D_\alpha (\mathbf{P}_\alpha \cdot \mathbf{\Omega}_\alpha / n_\alpha) / Dt = \nabla \cdot [(\dots) \mathbf{\Omega}_\alpha] + \text{friction}$ , where  $n_\alpha$  is the density. The generalized vorticity appearing in the divergence term implies the existence of a “local”  $\alpha$ -helicity associated with these lines,  $K_\alpha = (c^2/8\pi q_\alpha^2) \int_C d\tau \mathbf{P}_\alpha \cdot \mathbf{\Omega}_\alpha$ , where  $C$  is the volume occupied by a bundle of  $\alpha$ -vortex lines. The constant factor gives  $K_\alpha$  the convenient units of energy-length. The total derivative  $D_\alpha / Dt$  implies that the local  $\alpha$ -helicity convects with its own species. If an  $\alpha$ -vortex line does not intersect the system boundary, then in the strictly ideal (frictionless) case, the associated  $\alpha$ -helicity is constant. There is a circulation theorem,  $\Gamma_\alpha = \int_C \mathbf{P}_\alpha \cdot d\mathbf{x} = \text{const}$ , where  $C$  is an  $\alpha$ -vortex line, and  $d\mathbf{x}$  is a differential length vector along that line. Each species has its own set of  $\alpha$ -vortex lines, its own local  $\alpha$ -helicities, and its own circulation theorem.

In the realistic case with friction, visco-resistive instabilities drive reconnections that break individual  $\alpha$ -vortex lines and destroy their identity. This is a case of non-uniform convergence because even a minute amount of friction is enough to compromise the local  $\alpha$ -helicities. The only quantities immune to these topology altering events are the global  $\alpha$ -helicities,  $K_\alpha = (c^2/8\pi q_\alpha^2) \int_V d\tau \mathbf{P}_\alpha \cdot \mathbf{\Omega}_\alpha$ , where  $V$  is the system volume. Even global invariants may not be *rugged* in the sense that they are *more* “invariant” than the organized energy form, *i.e.* the magnetofluid energy  $W_{mf} = \int_V d\tau (\sum m_\alpha n_\alpha u_\alpha^2 + B^2/8\pi)$ , composed of the flow energy and the magnetic energy (the sum is over species). The ruggedness of the global  $\alpha$ -helicities has been supported by three arguments. (1) *Selective decay*:  $W_{mf}$  decays more rapidly than  $K_\alpha$  in thin reconnection layers. Properly applied, this argument must account for limits on viscous friction coefficients for sharp gradients. (2) *Inverse cascade*: the fluctuation spectrum of  $\tilde{W}_{mf}(k)$  and  $\tilde{K}_\alpha(k)$  satisfy the necessary conditions for a cascade toward larger scale objects ( $k$  is the wave number of the disturbance). (3) *Stability to resistive modes*:  $K_\alpha$  is less affected than  $W_{mf}$  by resistive modes. Each of these is the generalization of arguments previously applied to verify the ruggedness of the magnetic helicity in weakly-dissipative MHD.

A minimum energy state is found formally by minimizing  $W_{mf}$  subject to invariant  $\alpha$ -helicities, and (given axisymmetric system boundary) the global angular momentum,  $L_\theta = \int d\tau r \sum m_\alpha n_\alpha u_{\alpha\theta}$ . The variation with respect to  $\delta \mathbf{u}_\alpha$  leads to the flow equations:  $n_\alpha (\mathbf{u}_\alpha - \Omega r \hat{\theta}) = (\lambda_\alpha / \ell_c^2) \mathbf{\Omega}_\alpha$  where  $\lambda_\alpha$ ,  $\Omega$  are the Lagrange multipliers associated with

invariant  $\alpha$ -helicities and angular momentum, and  $\ell_c = c/\omega_{pi} = (m_i c^2 / 4\pi e^2)^2$  is the length scale. An entropy maximization procedure subject to invariant  $K_\alpha$ ,  $L_\theta$ , and total energy ( $W_{mf}$  + thermal) leads to the same equation. In addition, a global Bernoulli equation links the pressure to the flow by a relation that applies throughout the system volume. Note that an important feature of a two-fluid minimum energy state is the length scale  $\ell_c$ . A two-fluid may or may not relax to the minimum energy state depending on whether the fast mechanisms have been stabilized.

Reference: L.C. Steinhauer and A. Ishida, Phys. Plasmas **5**, 2609 (1998)