

Energy Levels of Germanium, Ge I through Ge XXXII

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Atomic energy levels of germanium have been compiled for all stages of ionization for which experimental data are available. No data have yet been published for Ge VIII through Ge XIII and Ge XXXII. Very accurate calculated values are compiled for Ge XXXI and XXXII. Experimental *g*-factors and leading percentages from calculated eigenvectors of levels are given. A value for the ionization energy, either experimental when available or theoretical, is included for the neutral atom and each ion.

Key words: atomic; energy levels; germanium; ions; spectra.

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1. Introduction

In 1952 C. E. Moore published a compilation of energy levels of germanium containing the results of extensive analyses of Ge I through Ge V. Today, we have experimental energy levels for most stages of ionization of Ge, predicted level values for Ge VIII (Mn-like), Ge XII (Sc-like), and Ge XIII (Ca-like), and very accurate calculated levels for Ge XXXI (He-like) and Ge XXXII (H-like). Only for the ions Ge IX through Ge XI are there no experimental or theoretical data. New results have been published on Ge I, Ge II, Ge IV, and Ge V, the spectra included in the 1952 compilation.

The present critical compilation of the atomic energy levels of germanium in all stages of ionization is part of

an ongoing program of the National Institute of Standards and Technology (formerly the National Bureau of Standards) Atomic Energy Levels Data Center to compile similar data for all the elements. These publications include helium by Martin [1973, 1987], sodium, magnesium, aluminum, and silicon by Martin and Zalubas [1981, 1980, 1979, 1983], and phosphorus and sulfur by Martin, Zalubas, and Musgrove [1985, 1990], potassium through nickel by Sugar and Corliss [1985], copper, krypton, and molybdenum by Sugar and Musgrove [1990, 1991, 1988], and lanthanum through lutetium by Martin, Zalubas, and Hagan [1978].

Companion works containing wavelengths for the higher stages of ionization have been prepared in collaboration with the Japanese Atomic Energy Research Institute in Tokai-Mura, Japan. These include titanium by Mori *et al.* [1986] and vanadium, iron, cobalt, nickel, copper, and molybdenum by Shirai *et al.* [1992b, 1990, 1992a, 1987a, 1991, 1987b]. In addition, wavelength compilations including data for all stages of ionization have

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been published for Sc by Kaufman and Sugar [1988] and for Mg and Al by Kaufman and Martin [1991a, 1991b].

The strong lines of Ge I through Ge V are contained in "Line Spectra of the Elements" edited by Reader and Corliss [1991], which appears in the *CRC Handbook of Chemistry and Physics* [1991]. These data also appear in the NIST Standard Reference Database 38 [1992]. A compilation published by Kelly [1987] gives wavelengths of germanium spectra below 2000 Å and their energy level classifications.

In the present work all energy levels are given in units of cm^{-1} . An estimate of the uncertainty of the energy level values or wavelengths determining them is given in each case. Ionization energies are also given in eV with the conversion factor $8065.5410(24) \text{ cm}^{-1}/\text{eV}$ published by Cohen and Taylor [1987].

We use without comment notations for various coupling schemes as appropriate. Martin, Zalubas, and Hagan [1978] give a complete summary of the coupling notations used here and tables of the allowed terms for equivalent electrons.

We have included under the heading "Leading percentages" the results of calculations that express the eigenvector percentage composition of levels (rounded to the nearest percent) in terms of the basis states of a single configuration, or more than one configuration where configuration interaction has been included. We give first the percentage of the basis state corresponding to the level's name; next the second largest percentage together with the related basis state. Generally, when the leading percentage is less than 40%, no name is given. However, when two different parent states give rise to the same final term type and the sum of their percentages is $\geq 40\%$, the level is designated by the higher percentage term. For an unnamed level, the term symbol for the leading percentage follows the percentage. The user should of course bear in mind that the percentages are model dependent, so that the results of different calculations may yield notably different percentages.

For configurations of equivalent d -electrons, several terms of the same LS type may occur. These are theoretically distinguished by their seniority number. In our compilations they are designated in the notation of Nielson and Koster [1963]. For example, in the $3d^5$ configuration there are three 2D terms with seniorities of 1, 3, and 5. These terms are denoted as 2D1 , 2D2 , and 2D3 respectively, by Nielson and Koster.

In cases where the ionization energies cannot be determined from the experimental data, we have calculated them with Cowan's [1981] Hartree Fock code (HFR) with relativistic and correlation corrections included. The uncertainty in these determinations varies from $\pm 0.1\% - 1.0\%$.

The text for each spectrum does not include a complete review of the literature but is intended to credit the major contributions. In assembling the data for each spectrum, we referred to the following bibliographies:

- i. Papers cited by Moore (1952)
- ii. C. E. Moore (1969)
- iii. L. Hagan and W. C. Martin (1972)
- iv. L. Hagan (1977)
- v. R. Zalubas and A. Albright (1980)
- vi. A. Musgrove and R. Zalubas (1985)
- vii. Bibliographic file of publications since December 1983 maintained by the NIST Atomic Energy Levels Data Center

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3. Tables of Energy Levels

Ge I

Z = 32

Ge I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 4s^2 4p^2 \ ^3P_0$ Ionization energy $63\,713.24 \pm 0.10 \text{ cm}^{-1}$ ($7.899\,438 \pm 0.000\,012 \text{ eV}$)

Early observations of this spectrum combined to produce a list of 184 lines in the range of 1639 Å to 11 144 Å. These include major contributions from Gartlein [1928], Rao [1929], and Kiess [1940]. Gartlein classified 73 lines in the range of 1873 Å to 4685 Å and established all the levels of the ground configuration and 29 levels of $4pnl$ for $nl = 5s - 7s, 4d, 5d$. Rao extended the observations to 1630 Å and classified most of the 50 new lines, modifying some of Gartlein's levels and adding to the known excited configurations. Kiess extended the measurements to the region from 5265 Å to 11 144 Å, and classified 25 of 44 lines as transitions between the excited configurations $4p5s - 4p5p$ and $6p$.

The $4s4p^3 \ ^5S_2^o$ term was discovered by Andrew and Meissner [1957], who recognized the two transitions to $4s^2 4p^2 \ ^3P_{2,1}$.

The spectrum was completely remeasured in the range of 1546–11 253 Å by Andrew and Meissner [1959] who used a low pressure arc in helium. From 4685 Å upward they measured more than 200 lines, most of which were new. From 4685 Å downward they observed 230 lines, 50 of which were new. Using these data they determined values and designations for 68 even and 88 odd levels. These include $4pns$ ($n = 5 - 10$), $4pnp$ ($n = 4 - 7$), $4pnd$ ($n = 4 - 8$), and $4pnf$ ($n = 4 - 7$), as well as $4s4p^3$. They obtained the value for the ionization energy from these series of $63\,715 \pm 10 \text{ cm}^{-1}$. Their wavelength uncertainty was $\pm 0.03 \text{ Å}$ for the spectrum observed in air above 2000 Å, and $\pm 0.02 \text{ Å}$ for vacuum ultraviolet observations below 2000 Å. Some interferometric measurements from 4600 Å to 11 125 Å were made with an uncertainty of $\pm 0.005 \text{ Å}$.

Many interferometric studies of this spectrum have been carried out, beginning with Deverall *et al.* [1954]. Using an atomic beam crossed with an electron source they measured eight lines from 2592–3269 Å with an uncertainty of $\pm 0.000\,04 \text{ Å}$. They were unable to resolve isotopic splitting of the lines but determined that it is less than 0.005 cm^{-1} . With these measurements they redetermined the values of the $4p^2 \ ^3P$ and 1D levels and the $4p5s \ ^1^3P^o$ with an uncertainty of $\pm 0.0005 \text{ cm}^{-1}$.

Van Veld and Meissner [1956] measured 43 lines in the region of 2019–4685 Å using a Fabry-Perot interferometer and a hollow cathode light source. They estimate their wavelength uncertainty to be $\pm 0.0004 \text{ Å}$. Using these measurements they were able to improve the values for the levels $4p^2 \ ^1S$, $4p4d \ ^1^3D^o$, $^1^3F^o$, $^1^3P^o$, and seven addi-

tional higher-lying odd levels. With these values they predicted wavelengths of 29 lines from 1691–1998 Å with an uncertainty of $\pm 0.0004 \text{ Å}$. These were proposed as secondary wavelength standards.

By making long wavelength interferometric measurements, Andrew and Meissner [1958, 1959] extended the number of precisely determined high even levels. These levels, combined with accurate wavelength measurements in the visible and infrared, were used to improve the values of many high odd levels. Transition wavelengths from these levels to the four even levels of the ground configuration are then accurately predicted from 1630 Å to 1997 Å. Using these as wavelength standards, Meissner *et al.* [1958] measured 30 additional Ge I lines from 1633 Å to 1944 Å with an uncertainty of $\pm 0.002 \text{ Å}$.

A new systematic program of interferometric measurements was undertaken by Kaufman and Andrew [1962] in order to improve and increase the vacuum ultraviolet Ritz standards. Their observations extended from 4380 to 12 069 Å and were combined with those of Humphreys and Paul [1961] in the region of 11 255 to 19 284 Å, and with the earlier measurements from 2019 to 4685 Å by Van Veld and Meissner [1956] and Deverall *et al.* [1954] to improve the energy level values. They give 100 Ritz standards below 2000 Å and 95 interferometrically determined levels. The wavelength uncertainties do not exceed $\pm 0.0009 \text{ Å}$, and 68 lines have uncertainties of $\pm 0.0003 \text{ Å}$.

The number of measured lines was increased to 100 in the range of 11 252 to 23 922 Å by Humphreys and Andrew [1964]. This work was intended for establishing the levels of the $4pnf$ configurations that had been tentatively identified earlier by single lines. One $4p5d$ and two $4p6p$ levels were also found. Percentage compositions of the $4p4f$ levels were given by Cowan and Andrew [1965] in jl -coupling. Those levels still determined by a single line are denoted here by question marks.

A high resolution absorption spectrum from 1500 Å to 1900 Å was reported by Brown *et al.* [1977]. They observed 989 lines and determined 549 levels in long series approaching the two $4s^2 4p^2 P^o$ limits. Their series include $4pns$, $4pnd$, and $4png$, most of whose levels are given for the first time. They determined the value for the series limit quoted here. A calculation of these series through $n = 14$ with interaction among the series and $4s4p^3$ was carried out by Dembczynski [1986]. He gave eigenvectors for each of the levels in LS -coupling. These show that be-

yond $n=5$ LS designations are generally meaningless. We give LS names for odd levels of $n=4$ and for most levels of $n=5$, and then group the remaining odd levels in jj -coupling terms based on the quantum defects determined by Brown *et al.* in their Lu-Fano treatment of these series. Dembczynski found that the $^3D^\circ$ and $^3S^\circ$ terms of $4s4p^3$ reported by Andrew and Meissner contain only minor percentages of this configuration. The $4s4p^3$ $^3D^\circ$ component is absorbed into the $4pnd$ series. The $4s4p^3$ $^3S_1^\circ$ level of Andrew and Meissner was changed to $4s^24p6d$ ($^3/2, ^3/2$) $^\circ$ with $J=2$ by Brown *et al.* We were unable to designate the following four levels as series members: 60 429.908 cm^{-1} with $J=3$, 60 516.308 cm^{-1} with $J=2$, 61 101.37 cm^{-1} with $J=1$, and 63 601.65 cm^{-1} with $J=2$.

We quote the energy level values from the following sources. The levels of $4p^2$ and $4p5s$ are from the interferometric atomic beam measurements of Deverall *et al.* The levels of $4p4d$ are from the interferometric hollow cathode measurements of Van Veld and Meissner based on the $4p^2$ levels of Deverall. Kaufman and Andrew extended the interferometric measurements from the visible to 12 000 \AA by means of the spectrum emitted by electrodeless lamps. These data were combined with the measurements of Humphreys and Paul from 11 000–19 000 \AA to extend the known levels to include $4p5p$ to $4p7p$, $4p5d$ to $4p9d$, and $4p6s$ to $4p9s$. Their results are quoted. The $4pnf$ levels are from Humphreys and Andrew and the earlier work by Andrew and Meissner [1959]. The long Rydberg series $4pns$, $4pnd$,

and $4png$ are from the absorption spectra observed by Brown *et al.*

By atomic-beam magnetic resonance Childs and Goodman [1964] measured g -factors with an uncertainty in the fifth decimal place for the levels $4p^2$ 3P_1 , 3P_2 , and 1D_2 . The rest of the observed values were obtained by Andrew *et al.* [1967] with an electrodeless lamp in a uniform magnetic field. They reported uncertainties in the second or third decimal place.

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Ge I

Configuration	Term	J	Level (cm^{-1})	g	Leading percentages	
$4s^24p^2$	3P	0	0.0000			
		1	557.1341	1.50111		
		2	1 409.9609	1.49458		
$4s^24p^2$	1D	2	7 125.2989	1.00639		
$4s^24p^2$	1S	0	16 367.3332			
$4s^24p5s$	$^3P^\circ$	0	37 451.6893		100	
		1	37 702.3054	1.435	82	
		2	39 117.9021	1.500	100	
$4s^24p5s$	$^1P^\circ$	1	40 020.5604	1.068	86	
$4s4p^3$	$^5S^\circ$	2	41 926.726	2.011	100	
		3D	1	45 985.592		
			2	46 834.3798	1.182	
3	48 104.1003		1.332			
$4s^24p5p$	1P	1	46 765.2705	0.980		
$4s^24p5p$	3P	0	47 502.6257			
		1	48 088.3504	1.362		
		2	48 726.1143	1.427		

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages	
4s ² 4p4d	¹ D°	2	48 480.048	0.981	74	14 4s ² 4p4d ³ F°
4s ² 4p4d	³ D°	2	48 882.263	1.008	57	23 4s ² 4p4d ³ F°
		1	48 962.783	0.556	77	14 4s4p ³ ³ D°
		3	49 144.397	1.240	56	28 4s ² 4p4d ³ F°
4s ² 4p5p	³ S	1	49 075.8917	1.862		
4s ² 4p5p	¹ D	2	49 649.5807	1.062		
4s ² 4p4d	³ F°	2	50 068.954	0.867	59	18 4s ² 4p4d ³ D°
		3	50 323.465	1.194	66	24 "
		4	50 786.79		99	
4s ² 4p5p	¹ S	0	51 011.4392			
4s ² 4p4d	³ P°	2	51 437.802	1.460	88	5 4s ² 4p4d ³ D°
		1	51 705.020	1.391	82	9 4s ² 4p4d ¹ P°
		0	51 978.15		89	7 4p6s ³ P°
4s ² 4p6s	(¹ / ₂ , ¹ / ₂)°	1	52 148.726	1.307		
		0	52 170.504			
4s ² 4p4d	¹ F°	3	52 592.224	1.010	90	4 4s ² 4p4d ³ F°
4s ² 4p4d	¹ P°	1	52 847.215	1.140	65	23 4p6s ³ P°
4s ² 4p6s	(³ / ₂ , ¹ / ₂)°	2	53 911.624			
		1	54 174.928	1.103		
4s ² 4p6p	³ D	1	54 935.8479	0.690		
		2	55 266.0902	1.177		
		3	56 793.460	1.297		
4s ² 4p6p	¹ P	1	55 235.8344	1.383		
4s ² 4p5d	³ D°	2	55 372.610		47	15 4p5d ³ P°
		1	55 474.668	0.600	58	10 4p4d ³ D°
		3	55 718.604		42	29 4p5d ³ F°
4s ² 4p6p	³ P	0	55 503.2026			
		1	56 687.1584	1.106		
		2	56 947.7693	1.357		
4s ² 4p5d	³ F°	2	55 686.70		52	37 4s ² 4p5d ¹ D°
		3	56 828.439		51	16 4s ² 4p5d ³ D°
		4	57 556.206		98	
4s ² 4p(² P _{1/2})4f	² [⁷ / ₂]	3	56 765.748		77	23 4s ² 4p(² P _{1/2})4f ² [⁵ / ₂]
		4	56 771.466		100	
4s ² 4p(² P _{1/2})4f	² [⁵ / ₂]	3	56 768.663	1.200	77	23 4s ² 4p(² P _{1/2})4f ² [⁷ / ₂]
		2	56 772.750	0.859	99	
4s ² 4p5d		2	56 921.345		31	4s ² 4p5d ³ F° 24 4s ² 4p5d ¹ D°
4s ² 4p6p	³ S	1	57 083.2021	1.804		

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages	
4s ² 4p7s	(1/2, 1/2) ^o	0	57 168.39			
		1	57 180.198			
4s ² 4p6p	¹ D	2	57 250.9427	1.128		
4s ² 4p5d	³ P ^o	1	57 398.852	1.207	52	21 4s ² 4p5d ³ D ^o
		2	57 430.924		57	28 4s ² 4p5d ¹ D ^o
		0	57 828.98		95	4 4p7s ³ P ^o
4s ² 4p6p	¹ S	0	57 670.326			
4s ² 4p5d	¹ P ^o	1	58 058.060	1.19	40	37 4s ² 4p5d ³ P ^o
4s ² 4p5d	¹ F ^o	3	58 093.352		63	16 4s ² 4p5d ³ F ^o
4s ² 4p7p	³ D	1	58 414.660	1.175		
		2	58 586.8822			
		3	60 218.379			
4s ² 4p(2P _{3/2})4f	² [7/2]	3	58 458.025		100	
		4	58 460.090		100	
4s ² 4p(2P _{3/2} ^o)4f	² [5/2]	3	58 519.672		100	
		2	58 520.737	0.88	99	
4s ² 4p5d		2	58 551.407		21	4s ² 4p5d ³ D ^o 6 4s4p ³ ³ D ^o
4s ² 4p7p	¹ P	1	58 560.827			
4s ² 4p(2P _{3/2} ^o)4f	² [9/2]	5	58 575.02		100	
		4	58 578.867		100	
4s ² 4p(2P _{3/2} ^o)4f	² [3/2]	1	58 642.22		100	
		2	58 645.062	1.10	100	
4s ² 4p7p	³ P	0	58 733.515			
		1	60 187.58			
		2	60 320.616			
4s ² 4p6d	(1/2, 3/2) ^o	1	58 741.040			
		2	58 747.644			
4s ² 4p7s	(3/2, 1/2) ^o	2	58 931.503			
		1	59 114.742			
4s ² 4p5d		3	58 943.436	1.17		
4s ² 4p(2P _{1/2})5f	² [7/2]	3	59 275.91			
		4	59 278.21			
4s ² 4p(2P _{1/2})5f	² [5/2]	3	59 282.900			
		2	59 283.715			
4s ² 4p8s	(1/2, 1/2) ^o	0	59 494.515			
		1	59 524.276			

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p6d	(1/2, 5/2) ^o	3	59 658.391	1.15	
		2	59 690.604		
		1	59 727.511		
4s ² 4p6d	(3/2, 3/2) ^o	2	60 270.231		
		1	60 549.958		
		0	60 607.272		
		3	60 886.24		
4s ² 4p8p	¹ P	1	60 273.738		
4s ² 4p7p	³ S	1	60 354.624		
4s ² 4p7d	(1/2, 3/2) ^o	2	60 403.365		
		1	60 857.12		
4s ² 4p7d	(1/2, 5/2) ^o	3	60 429.908		
		2	60 516.308		
4s ² 4p7p	¹ D	2	60 456.466		
4s ² 4p6d	(3/2, 5/2) ^o	4	60 552.366		
		2	61 091.48		
		1	61 152.37		
		3	61 268.39		
4s ² 4p(2P _{1/2} ^o)6f	² [7/2]	3	60 633.157		
		4	60 634.698		
4s ² 4p(2P _{1/2} ^o)6f	² [5/2]	3	60 634.467		
		2	60 635.152		
4s ² 4p6g	(1/2, 7/2) ^o	3	60 658.13		
4s ² 4p7p	¹ S	0	60 673.064		
4s ² 4p9s	(1/2, 1/2) ^o	1	60 749.318		
		0	60 769.238		
4s ² 4p(2P _{3/2} ^o)5f	² [7/2]	3	61 002.709		
		4	61 004.356		
4s ² 4p(2P _{3/2} ^o)5f	² [5/2]	3	61 035.830		
		2	61 036.635		
4s ² 4p(2P _{3/2} ^o)5f	² [9/2]	5	61 060.41		
		4	61 063.611		
4s ² 4p(2P _{3/2} ^o)5f	² [3/2]	1	61 094.34?		
		2	61 095.40		
4s ² 4p8s	(3/2, 1/2) ^o	2	61 254.02		
		1	61 343.172		
4s ² 4p8d	(1/2, 3/2) ^o	2	61 269.082		
		1	61 571.41		
4s ² 4p(2P _{1/2} ^o)7f	² [7/2]	3	61 454.32		
		4	61 457.24?		

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p(² P _{1/2})7f	² [⁵ / ₂]	3	61 457.20		
		2	61 457.61?		
4s ² 4p8d	(1/2, ⁵ / ₂)°	3	61 522.73		
		2	61 542.72		
4s ² 4p10s	(1/2, ¹ / ₂)°	0	61 539.12		
		1	61 546.23		
4s ² 4p9d	(1/2, ³ / ₂)°	2	61 849.22		
		1	61 997.14		
4s ² 4p7d	(3/2, ³ / ₂)°	2	61 922.47		
		0	62 124.89		
		1	62 169.26		
		3	62 370.79		
4s ² 4p9d	(1/2, ⁵ / ₂)°	3	61 930.04		
		2	62 054.745		
4s ² 4p8g	(1/2, ⁷ / ₂)°	3	61 995.64		
4s ² 4p11s	(1/2, ¹ / ₂)°	0	62 041.45		
		1	62 044.93		
4s ² 4p12s	(1/2, ¹ / ₂)°	0	62 390.37		
		1	62 398.94		
4s ² 4p13s	(1/2, ¹ / ₂)°	0	62 639.15		
		1	62 647.61		
4s ² 4p14s	(1/2, ¹ / ₂)°	0	62 823.67		
		1	62 826.79		
4s ² 4p15s	(1/2, ¹ / ₂)°	0	62 964.21		
		1	62 965.59		
4s ² 4p16s	(1/2, ¹ / ₂)°	1	63 073.11		
		0	63 075.26		
4s ² 4p17s	(1/2, ¹ / ₂)°	0	63 162.09		
		1	63 165.54		
4s ² 4p18s	(1/2, ¹ / ₂)°	0	63 232.55		
		1	63 234.16		
4s ² 4p19s	(1/2, ¹ / ₂)°	0	63 290.65		
		1	63 293.33		
4s ² 4p20s	(1/2, ¹ / ₂)°	0	63 338		
		1	63 342.57		
4s ² 4p21s	(1/2, ¹ / ₂)°	0	63 378.68		
		1	63 380.02		
4s ² 4p22s	(1/2, ¹ / ₂)°	0	63 412.81		
		1	63 413.65		
4s ² 4p23s	(1/2, ¹ / ₂)°	0	63 441.91		
		1	63 442.57		
4s ² 4p24s	(1/2, ¹ / ₂)°	0	63 466.98		
		1	63 467.49		

Ge I - Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p25s	(1/2, 1/2) ^o	0	63 488.83		
		1	63 489.20		
4s ² 4p26s	(1/2, 1/2) ^o	0	63 507.77		
		1	63 508.13		
4s ² 4p27s	(1/2, 1/2) ^o	0	63 524.54		
		1	63 524.64		
4s ² 4p28s	(1/2, 1/2) ^o	0	63 539.41		
		1	63 539.50		
4s ² 4p29s	(1/2, 1/2) ^o	0	63 552.31		
		1	63 552.44		
4s ² 4p30s	(1/2, 1/2) ^o	0	63 563.99		
		1	63 564.13		
4s ² 4p31s	(1/2, 1/2) ^o	0	63 574.38		
		1	63 574.45		
4s ² 4p32s	(1/2, 1/2) ^o	0	63 583.73		
		1	63 583.80		
4s ² 4p33s	(1/2, 1/2) ^o	0	63 592.27		
		1	63 592.40		
4s ² 4p34s	(1/2, 1/2) ^o	0	63 599.89		
		1	63 600.03		
4s ² 4p35s	(1/2, 1/2) ^o	1	63 607.02		
		0	63 607.05		
4s ² 4p36s	(1/2, 1/2) ^o	0	63 613.13		
4s ² 4p37s	(1/2, 1/2) ^o	0	63 618.78		
4s ² 4p38s	(1/2, 1/2) ^o	0	63 624.34		
4s ² 4p39s	(1/2, 1/2) ^o	0	63 629.31		
4s ² 4p40s	(1/2, 1/2) ^o	0	63 633.64		
4s ² 4p41s	(1/2, 1/2) ^o	0	63 637.76		
4s ² 4p42s	(1/2, 1/2) ^o	0	63 641.56		
4s ² 4p43s	(1/2, 1/2) ^o	0	63 645.07		
4s ² 4p7d	(3/2, 5/2) ^o	4	62 125.233		
		2	62 454.86		
		1	62 467.79		
		3	62 522.98		
4s ² 4p10d	(1/2, 5/2) ^o	3	62 217.75		
		2	62 232.87		
4s ² 4p10d	(1/2, 3/2) ^o	2	62 264.43		
		1	62 355.15		
4s ² 4p9g	(1/2, 7/2) ^o	3	62 355.61		
4s ² 4p(2P _{3/2})6f	2[7/2]	3	62 379.95?		
		4	62 381.86?		

Ge I - Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p(2P _{3/2})6f	2 ¹ [⁹ / ₂]	5	62 412.46		
		4	62 414.70?		
4s ² 4p6g	(³ / ₂ , ⁷ / ₂)°	3	62 437.63		
4s ² 4p9s	(3/2,1/2) ^o	2	62 531.56		
		1	62 576.74		
4s ² 4p11d	(1/2,3/2) ^o	2	62 545.56		
		1	62 639.17		
4s ² 4p12d	(1/2,3/2) ^o	2	62 751.67		
		1	62 814.88		
4s ² 4p13d	(1/2,3/2) ^o	2	62 906.37		
		1	62 948.76		
4s ² 4p14d	(1/2,3/2) ^o	2	63 046.08		
4s ² 4p15d	(1/2,3/2) ^o	2	63 128.81		
		1	63 156.45		
4s ² 4p16d	(1/2,3/2) ^o	2	63 204.74		
		1	63 218.03		
4s ² 4p17d	(1/2,3/2) ^o	2	63 267.51		
		1	63 290.86		
4s ² 4p18d	(1/2,3/2) ^o	2	63 319.49		
		1	63 338.16		
4s ² 4p19d	(1/2,3/2) ^o	2	63 362.52		
		1	63 377.61		
4s ² 4p20d	(1/2,3/2) ^o	2	63 398.92		
		1	63 411.40		
4s ² 4p21d	(1/2,3/2) ^o	2	63 429.94		
		1	63 440.39		
4s ² 4p22d	(1/2,3/2) ^o	2	63 456.73		
		1	63 465.46		
4s ² 4p23d	(1/2,3/2) ^o	2	63 479.74		
		1	63 487.23		
4s ² 4p24d	(1/2,3/2) ^o	2	63 499.78		
		1	63 506.23		
4s ² 4p25d	(1/2,3/2) ^o	2	63 517.32		
		1	63 523.00		
4s ² 4p26d	(1/2,3/2) ^o	2	63 533.09		
		1	63 537.76		
4s ² 4p27d	(1/2,3/2) ^o	2	63 546.3		
		1	63 550.78		
4s ² 4p28d	(1/2,3/2) ^o	1	63 562.46		
4s ² 4p29d	(1/2,3/2) ^o	1	63 572.90		
4s ² 4p30d	(1/2,3/2) ^o	1	63 582.23		
4s ² 4p31d	(1/2,3/2) ^o	1	63 590.62		
4s ² 4p32d	(1/2,3/2) ^o	1	63 598.15		

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p33d	(¹ / ₂ , ³ / ₂)°	1	63 605.06		
4s ² 4p34d	(¹ / ₂ , ³ / ₂)°	1	63 611.24		
4s ² 4p36d	(¹ / ₂ , ³ / ₂)°	1	63 622.05		
4s ² 4p37d	(¹ / ₂ , ³ / ₂)°	1	63 626.76		
4s ² 4p38d	(¹ / ₂ , ³ / ₂)°	1	63 631.12		
4s ² 4p39d	(¹ / ₂ , ³ / ₂)°	1	63 635.17		
4s ² 4p40d	(¹ / ₂ , ³ / ₂)°	1	63 638.98		
4s ² 4p41d	(¹ / ₂ , ³ / ₂)°	1	63 642.53		
4s ² 4p42d	(¹ / ₂ , ³ / ₂)°	1	63 645.72		
4s ² 4p43d	(¹ / ₂ , ³ / ₂)°	1	63 648.99		
4s ² 4p44d	(¹ / ₂ , ³ / ₂)°	1	63 651.91		
4s ² 4p45d	(¹ / ₂ , ³ / ₂)°	1	63 654.63		
4s ² 4p46d	(¹ / ₂ , ³ / ₂)°	1	63 657.22		
4s ² 4p47d	(¹ / ₂ , ³ / ₂)°	1	63 659.69		
4s ² 4p48d	(¹ / ₂ , ³ / ₂)°	1	63 661.94		
4s ² 4p49d	(¹ / ₂ , ³ / ₂)°	1	63 664.10		
4s ² 4p50d	(¹ / ₂ , ³ / ₂)°	1	63 666.10		
4s ² 4p51d	(¹ / ₂ , ³ / ₂)°	1	63 667.95		
4s ² 4p52d	(¹ / ₂ , ³ / ₂)°	1	63 669.81		
4s ² 4p53d	(¹ / ₂ , ³ / ₂)°	1	63 671.46		
4s ² 4p54d	(¹ / ₂ , ³ / ₂)°	1	63 673.04		
4s ² 4p55d	(¹ / ₂ , ³ / ₂)°	1	63 674.47		
4s ² 4p56d	(¹ / ₂ , ³ / ₂)°	1	63 675.93		
4s ² 4p57d	(¹ / ₂ , ³ / ₂)°	1	63 677.24		
4s ² 4p58d	(¹ / ₂ , ³ / ₂)°	1	63 678.56		
4s ² 4p59d	(¹ / ₂ , ³ / ₂)°	1	63 679.74		
4s ² 4p60d	(¹ / ₂ , ³ / ₂)°	1	63 680.84		
4s ² 4p61d	(¹ / ₂ , ³ / ₂)°	1	63 682.03		
4s ² 4p62d	(¹ / ₂ , ³ / ₂)°	1	63 683.1		
4s ² 4p63d	(¹ / ₂ , ³ / ₂)°	1	63 684.03		
4s ² 4p64d	(¹ / ₂ , ³ / ₂)°	1	63 684.87		
4s ² 4p65d	(¹ / ₂ , ³ / ₂)°	1	63 685.70		
4s ² 4p66d	(¹ / ₂ , ³ / ₂)°	1	63 686.70		
4s ² 4p67d	(¹ / ₂ , ³ / ₂)°	1	63 687.40		
4s ² 4p10g	(¹ / ₂ , ⁷ / ₂)°	3	62 614.61		
4s ² 4p11d	(¹ / ₂ , ⁵ / ₂)°	3	62 629.23		
		2	62 635.22		

Ge I - Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p 12d	(1/2, 5/2) ^o	3	62 791.97		
		2	62 793.56		
4s ² 4p 13d	(1/2, 5/2) ^o	2	62 926.05		
		3	62 930.26		
4s ² 4p 14d	(1/2, 5/2) ^o	2	63 033.89		
		3	63 037.26		
4s ² 4p 15d	(1/2, 5/2) ^o	3	63 113.76		
		2	63 117.23		
4s ² 4p 16d	(1/2, 5/2) ^o	2	63 186.44		
		3	63 223.49		
4s ² 4p 17d	(1/2, 5/2) ^o	2	63 294.10		
		3	63 302.32		
4s ² 4p 18d	(1/2, 5/2) ^o	3	63 337.03		
		2	63 337.89		
4s ² 4p 19d	(1/2, 5/2) ^o	3	63 374.34		
		2	63 375.62		
4s ² 4p 20d	(1/2, 5/2) ^o	3	63 407.82		
		2	63 408.58		
4s ² 4p 21d	(1/2, 5/2) ^o	3	63 436.95		
		2	63 437.27		
4s ² 4p 22d	(1/2, 5/2) ^o	2,3	63 462.29		
4s ² 4p 23d	(1/2, 5/2) ^o	2	63 484.12		
		3	63 484.37		
4s ² 4p 24d	(1/2, 5/2) ^o	2	63 503.29		
		3	63 503.72		
4s ² 4p 25d	(1/2, 5/2) ^o	2	63 520.23		
		3	63 520.69		
4s ² 4p 26d	(1/2, 5/2) ^o	2	63 535.22		
		3	63 535.68		
4s ² 4p 27d	(1/2, 5/2) ^o	2	63 548.46		
		3	63 549.00		
4s ² 4p 28d	(1/2, 5/2) ^o	2	63 560.43		
		3	63 560.90		
4s ² 4p 29d	(1/2, 5/2) ^o	2	63 571.11		
		3	63 571.58		
4s ² 4p 30d	(1/2, 5/2) ^o	2	63 580.74		
		3	63 581.16		
4s ² 4p 31d	(1/2, 5/2) ^o	2	63 589.46		
		3	63 589.77		
4s ² 4p 32d	(1/2, 5/2) ^o	2	63 597.38		
		3	63 597.62		
4s ² 4p 33d	(1/2, 5/2) ^o	2	63 604.41		
4s ² 4p 34d	(1/2, 5/2) ^o	2	63 611.21		
4s ² 4p 35d	(1/2, 5/2) ^o	2	63 617.0		
4s ² 4p 36d	(1/2, 5/2) ^o	2	63 622.57		

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p 37d	(¹ / ₂ , ⁵ / ₂)°	2	63 627.36		
4s ² 4p 38d	(¹ / ₂ , ⁵ / ₂)°	2	63 632.2		
4s ² 4p 39d	(¹ / ₂ , ⁵ / ₂)°	2	63 635.90		
4s ² 4p 40d	(¹ / ₂ , ⁵ / ₂)°	2	63 639.91		
4s ² 4p 41d	(¹ / ₂ , ⁵ / ₂)°	2	63 643.50		
4s ² 4p 43d	(¹ / ₂ , ⁵ / ₂)°	2	63 650.03		
4s ² 4p 44d	(¹ / ₂ , ⁵ / ₂)°	2	63 652.94		
4s ² 4p 45d	(¹ / ₂ , ⁵ / ₂)°	2	63 655.62		
4s ² 4p 46d	(¹ / ₂ , ⁵ / ₂)°	2	63 658.17		
4s ² 4p 47d	(¹ / ₂ , ⁵ / ₂)°	2	63 660.52		
4s ² 4p 48d	(¹ / ₂ , ⁵ / ₂)°	2	63 662.83		
4s ² 4p 49d	(¹ / ₂ , ⁵ / ₂)°	2	63 664.90		
4s ² 4p 50d	(¹ / ₂ , ⁵ / ₂)°	2	63 666.78		
4s ² 4p 51d	(¹ / ₂ , ⁵ / ₂)°	2	63 668.68		
4s ² 4p 52d	(¹ / ₂ , ⁵ / ₂)°	2	63 670.44		
4s ² 4p 53d	(¹ / ₂ , ⁵ / ₂)°	2	63 672.12		
4s ² 4p 54d	(¹ / ₂ , ⁵ / ₂)°	2	63 673.57		
4s ² 4p 55d	(¹ / ₂ , ⁵ / ₂)°	2	63 675.04		
4s ² 4p 56d	(¹ / ₂ , ⁵ / ₂)°	2	63 676.45		
4s ² 4p 57d	(¹ / ₂ , ⁵ / ₂)°	2	63 677.76		
4s ² 4p 58d	(¹ / ₂ , ⁵ / ₂)°	2	63 679.02		
4s ² 4p 59d	(¹ / ₂ , ⁵ / ₂)°	2	63 680.17		
4s ² 4p 60d	(¹ / ₂ , ⁵ / ₂)°	2	63 681.27		
4s ² 4p 61d	(¹ / ₂ , ⁵ / ₂)°	2	63 682.28		
4s ² 4p 62d	(¹ / ₂ , ⁵ / ₂)°	2	63 683.36		
4s ² 4p 63d	(¹ / ₂ , ⁵ / ₂)°	2	63 684.39		
4s ² 4p 64d	(¹ / ₂ , ⁵ / ₂)°	2	63 685.26		
4s ² 4p 65d	(¹ / ₂ , ⁵ / ₂)°	2	63 686.17		
4s ² 4p 66d	(¹ / ₂ , ⁵ / ₂)°	2	63 686.96		
4s ² 4p 67d	(¹ / ₂ , ⁵ / ₂)°	2	63 687.71		
4s ² 4p 68d	(¹ / ₂ , ⁵ / ₂)°	2	63 688.52		
4s ² 4p 69d	(¹ / ₂ , ⁵ / ₂)°	2	63 689.26		
4s ² 4p 70d	(¹ / ₂ , ⁵ / ₂)°	2	63 689.94		
4s ² 4p 11g	(¹ / ₂ , ⁷ / ₂)°	3	62 805.32		
4s ² 4p 12g	(¹ / ₂ , ⁷ / ₂)°	3	62 950.30		
4s ² 4p 13g	(¹ / ₂ , ⁷ / ₂)°	3	63 063.31		
4s ² 4p 15g	(¹ / ₂ , ⁷ / ₂)°	3	63 224.64		

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p8d	(3/2, 3/2) ^o	2	62 995.59		
		0	63 032.97		
		1	63 089.10		
		3	63 168.53		
4s ² 4p(2P _{3/2} ^o)7f	2[9/2]	5	63 227.73		
4s ² 4p8d	(3/2, 5/2) ^o	2	63 244.22		
		1	63 251.20		
		3	63 271.31		
4s ² 4p10s	(3/2, 1/2) ^o	2	63 305.19		
		1	63 330.45		
4s ² 4p9d	(3/2, 3/2) ^o	2	63 606.08		
		0	63 625.69		
		3	63 715		
Ge II (2P _{1/2} ^o)	Limit		63 713.24		
4s ² 4p9d	(3/2, 5/2) ^o	1	63 763.1		
		3	63 790		
4s ² 4p11s	(3/2, 1/2) ^o	2	63 809.51		
		1	63 827.3		
4s ² 4p12s	(3/2, 1/2) ^o	2	64 156.72		
		1	64 168.9		
4s ² 4p13s	(3/2, 1/2) ^o	2	64 405.88		
		1	64 415		
4s ² 4p14s	(3/2, 1/2) ^o	2	64 590.78		
		1	64 597.3		
4s ² 4p15s	(3/2, 1/2) ^o	2	64 731.78		
		1	64 737.0		
4s ² 4p16s	(3/2, 1/2) ^o	2	64 841.72		
		1	64 845.8		
4s ² 4p17s	(3/2, 1/2) ^o	2	64 929.13		
		1	64 932.3		
4s ² 4p18s	(3/2, 1/2) ^o	2	64 999.66		
		1	65 002.4		
4s ² 4p19s	(3/2, 1/2) ^o	2	65 057.56		
		1	65 059.8		
4s ² 4p20s	(3/2, 1/2) ^o	2	65 105.54		
		1	65 107.4		
4s ² 4p21s	(3/2, 1/2) ^o	2	65 145.83		
		1	65 147.4		
4s ² 4p22s	(3/2, 1/2) ^o	2	65 179.99		
		1	65 181.30		
4s ² 4p23s	(3/2, 1/2) ^o	2	65 209.14		
		1	65 210.2		

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p24s	(3/2, 1/2) ^o	2	65 234.22		
		1	65 235.12		
4s ² 4p25s	(3/2, 1/2) ^o	2	65 256.09		
		1	65 256.9		
4s ² 4p26s	(3/2, 1/2) ^o	2	65 275.0		
		1	65 275.79		
4s ² 4p27s	(3/2, 1/2) ^o	1	65 292.2		
4s ² 4p28s	(3/2, 1/2) ^o	1	65 307.0		
4s ² 4p29s	(3/2, 1/2) ^o	1	65 319.6		
4s ² 4p30s	(3/2, 1/2) ^o	1	65 331.6		
4s ² 4p31s	(3/2, 1/2) ^o	1	65 342.0		
4s ² 4p32s	(3/2, 1/2) ^o	1	65 351.3		
4s ² 4p33s	(3/2, 1/2) ^o	1	65 359.9		
4s ² 4p10d	(3/2, 3/2) ^o	1	64 040		
		3	64 084		
4s ² 4p10d	(3/2, 5/2) ^o	1	64 127		
		3	64 144		
4s ² 4p11d	(3/2, 3/2) ^o	0	64 298.98		
		1	64 320		
		3	64 357		
4s ² 4p12d	(3/2, 3/2) ^o	1	64 505		
		3	64 556		
4s ² 4p13d	(3/2, 3/2) ^o	0	64 671.62		
		1	64 680		
		3	64 702		
4s ² 4p14d	(3/2, 3/2) ^o	0	64 794.06		
		1	64 800		
		3	64 818		
4s ² 4p15d	(3/2, 3/2) ^o	0	64 890.64		
		1	64 902		
		3	64 910		
4s ² 4p16d	(3/2, 3/2) ^o	0	64 968.5		
		1	64 975		
		3	64 980		
4s ² 4p17d	(3/2, 3/2) ^o	0	65 031.66		
		1	65 036		
		3	65 045		
4s ² 4p18d	(3/2, 3/2) ^o	0	65 083.90		
		1	65 089		
		3	65 095		
4s ² 4p19d	(3/2, 3/2) ^o	1	65 132		
		3	65 138.0		
4s ² 4p20d	(3/2, 3/2) ^o	0	65 164.47		
		1	65 167		
		3	65 173		

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p21d	(3/2, 3/2) ^o	0	65 195.0		
		1	65 199		
		3	65 203		
4s ² 4p22d	(3/2, 3/2) ^o	0	65 222.66		
		1	65 225		
		3	65 229		
4s ² 4p23d	(3/2, 3/2) ^o	1	65 249		
		3	65 251.2		
4s ² 4p24d	(3/2, 3/2) ^o	0	65 266.37		
		1	65 268.8		
		3	65 270.5		
4s ² 4p25d	(3/2, 3/2) ^o	1	65 286		
		3	65 287.99		
4s ² 4p26d	(3/2, 3/2) ^o	1	65 301.0		
		3	65 303.17		
4s ² 4p27d	(3/2, 3/2) ^o	1	65 315		
		3	65 316.51		
4s ² 4p28d	(3/2, 3/2) ^o	1	65 327		
		3	65 328.64		
4s ² 4p29d	(3/2, 3/2) ^o	1	65 337.75		
		3	65 339.41		
4s ² 4p30d	(3/2, 3/2) ^o	1	65 347.9		
		3	58 348.99		
4s ² 4p31d	(3/2, 3/2) ^o	1	65 356.50		
		3	65 357.65		
4s ² 4p32d	(3/2, 3/2) ^o	1	65 364.38		
		3	65 365.48		
4s ² 4p33d	(3/2, 3/2) ^o	1	65 371.74		
		3	65 372.54		
4s ² 4p34d	(3/2, 3/2) ^o	1	65 378.40		
		3	65 378.91		
4s ² 4p35d	(3/2, 3/2) ^o	1	65 384.13		
		3	65 384.94		
4s ² 4p36d	(3/2, 3/2) ^o	1	65 389.68		
		3	65 390.24		
4s ² 4p37d	(3/2, 3/2) ^o	1	65 394.62		
4s ² 4p38d	(3/2, 3/2) ^o	1	65 399.90		
4s ² 4p39d	(3/2, 3/2) ^o	1	65 403.5		
4s ² 4p40d	(3/2, 3/2) ^o	1	65 407.3		
4s ² 4p41d	(3/2, 3/2) ^o	1	65 411.0		
4s ² 4p11d	(3/2, 5/2) ^o	1	64 382.4		
		3	64 396		
4s ² 4p12d	(3/2, 5/2) ^o	1	64 572.90		
		3	64 581.6		

Ge I - Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p 13d	(3/2, 5/2) ^o	1	64 717.9		
		3	64 724.4		
4s ² 4p 14d	(3/2, 5/2) ^o	1	64 830.30		
		3	64 835.7		
		2	64 837		
4s ² 4p 15d	(3/2, 5/2) ^o	1	64 920.14		
		2	64 922		
		3	64 924.1		
4s ² 4p 16d	(3/2, 5/2) ^o	1	64 992.35		
		2	64 994		
		3	64 995.6		
4s ² 4p 17d	(3/2, 5/2) ^o	1	65 051.60		
		2	65 052		
		3	65 054.1		
4s ² 4p 18d	(3/2, 5/2) ^o	1	65 100.56		
		2	65 102		
		3	65 102.7		
4s ² 4p 19d	(3/2, 5/2) ^o	1	65 141.72		
		2	65 142		
		3	65 143.2		
4s ² 4p 20d	(3/2, 5/2) ^o	1	65 176.32		
		2	65 177		
		3	65 177.72		
4s ² 4p 21d	(3/2, 5/2) ^o	1	65 206.13		
		2	65 207		
		3	65 207.2		
4s ² 4p 22d	(3/2, 5/2) ^o	1	65 231.75		
		2	65 232		
		3	65 232.5		
4s ² 4p 23d	(3/2, 5/2) ^o	1	65 253.80		
		2	65 254		
		3	65 254.5		
4s ² 4p 24d	(3/2, 5/2) ^o	1	65 273.18		
		3	65 273.93		
4s ² 4p 25d	(3/2, 5/2) ^o	1	65 290.15		
		3	65 290.5		
4s ² 4p 26d	(3/2, 5/2) ^o	1	65 305.08		
		3	65 305.43		
4s ² 4p 27d	(3/2, 5/2) ^o	1	65 318.34		
		3	65 318.6		
4s ² 4p 28d	(3/2, 5/2) ^o	1	65 330.17		
		3	65 330.3		
4s ² 4p 29d	(3/2, 5/2) ^o	1	65 340.78		
		3	65 341.1		
4s ² 4p 30d	(3/2, 5/2) ^o	1	65 350.27		
		3	65 350.5		
4s ² 4p 31d	(3/2, 5/2) ^o	3	65 358.8		
		1	65 358.93		

Ge I — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p32d	(3/2, 5/2) ^o	1	65 366.69		
		3	65 366.7		
4s ² 4p33d	(3/2, 5/2) ^o	1	65 373.66		
		3	65 373.8		
4s ² 4p34d	(3/2, 5/2) ^o	1	65 380.10		
		3	65 380.1		
4s ² 4p35d	(3/2, 5/2) ^o	1	65 385.93		
		3	65 386.1		
4s ² 4p36d	(3/2, 5/2) ^o	3	65 391.2		
		1	65 391.31		
4s ² 4p37d	(3/2, 5/2) ^o	3	65 396		
		1	65 396.14		
4s ² 4p38d	(3/2, 5/2) ^o	3	65 400.6		
		1	65 400.67		
4s ² 4p39d	(3/2, 5/2) ^o	1	65 404.78		
		3	65 404.8		
4s ² 4p40d	(3/2, 5/2) ^o	1	65 408.59		
		3	65 408.6		
4s ² 4p41d	(3/2, 5/2) ^o	1	65 412.07		
		3	65 412.2		
4s ² 4p42d	(3/2, 5/2) ^o	1	65 415.38		
		3	65 415.5		
4s ² 4p43d	(3/2, 5/2) ^o	3	65 418.5		
		1	65 418.59		
4s ² 4p44d	(3/2, 5/2) ^o	1	65 421.32		
		3	65 421.4		
4s ² 4p45d	(3/2, 5/2) ^o	1	65 424.01		
		3	65 424.2		
4s ² 4p46d	(3/2, 5/2) ^o	1	65 426.49		
		3	65 426.5		
4s ² 4p47d	(3/2, 5/2) ^o	1	65 428.86		
		3	65 428.9		
4s ² 4p48d	(3/2, 5/2) ^o	1	65 431.07		
		3	65 431.1		
4s ² 4p49d	(3/2, 5/2) ^o	1	65 433.02		
		3	65 433.1		
4s ² 4p50d	(3/2, 5/2) ^o	3	65 435.1		
		1	65 435.15		
4s ² 4p51d	(3/2, 5/2) ^o	1	65 436.82		
		3	65 436.9		
4s ² 4p52d	(3/2, 5/2) ^o	1	65 438.51		
		3	65 438.6		
4s ² 4p53d	(3/2, 5/2) ^o	1	65 440.13		
		3	65 440.24		
4s ² 4p54d	(3/2, 5/2) ^o	1	65 441.62		
		3	65 441.79		

Ge I -- Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s ² 4p55d	(³ / ₂ , ⁵ / ₂)°	1	65 443.1		
4s ² 4p56d	(³ / ₂ , ⁵ / ₂)°	1	65 444.42		
4s ² 4p57d	(³ / ₂ , ⁵ / ₂)°	1 3	65 445.69 65 445.85		
4s ² 4p58d	(³ / ₂ , ⁵ / ₂)°	1 3	65 446.93 65 446.99		
4s ² 4p59d	(³ / ₂ , ⁵ / ₂)°	1 3	65 448.12 65 448.19		
4s ² 4p60d	(³ / ₂ , ⁵ / ₂)°	1,3	65 449.16		
4s ² 4p61d	(³ / ₂ , ⁵ / ₂)°	1,3	65 450.29		
4s ² 4p62d	(³ / ₂ , ⁵ / ₂)°	1 3	65 451.12 65 451.28		
4s ² 4p63d	(³ / ₂ , ⁵ / ₂)°	1 3	65 452.16 65 452.28		
4s ² 4p64d	(³ / ₂ , ⁵ / ₂)°	1	65 453.02		
Ge II (² P _{3/2})	<i>Limit</i>		65 480.60		

Ge II

Z = 32

Ga I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 4s^2 4p^2 P_{1/2}^{\circ}$ Ionization energy $128\,521.3 \pm 0.2 \text{ cm}^{-1}$ ($15.93462 \pm 0.00002 \text{ eV}$)

Lang [1928, 1929] observed this spectrum with a spark in air, in hydrogen and in vacuum and classified 33 lines in the range of 999–7147 Å. Gartlein [1950] extended the observations from 813 Å to 6484 Å and classified 131 lines. His unpublished results are quoted by Moore [1952].

Accurate measurements of 22 lines from 4741 to 7145 Å, twelve of them measured interferometrically with an uncertainty of ± 0.005 Å, enabled Andrew and Meissner [1959] to determine the fine structure of the $4s^2 nf$ terms with $n = 4-7$. They discovered that all of these terms are inverted.

New interferometric measurements of 23 lines with an uncertainty of ± 0.001 Å were made by Kaufman and Andrew [1962] in the range of 2831–9475 Å. These data enabled them to improve energy level values and to calculate 12 Ritz standards from 999 to 1966 Å with an uncertainty of ± 0.0009 Å. Their improved level values for the lower members of the nf and ng series were used to calculate the value for the ionization energy quoted here.

The spectrum was reobserved from 600 to 10 000 Å by Shenstone [1963] with a hollow cathode source and with a low pressure spark in helium. In addition to improving the energy level values and resolving some of the doublets, he established many levels above the ionization threshold and classified many diffuse lines as transitions from these levels.

New measurements of five lines from 1649 to 2007 Å with an uncertainty of ± 0.001 Å were made by Kaufman *et al.* [1966]. Four of the lines are the spin-forbidden transitions from $4s4p^2 \text{ } ^4P$ to the $4s^2 4p^2 \text{ } ^2P^{\circ}$ ground term and one line is from $5s \text{ } ^2S$.

Kaufman and Ward [1966] measured 14 new lines in the range of 999–3224 Å with an uncertainty of ± 0.001 Å or better, depending on the order of diffraction in which they were measured. They have revised the values of the energy levels and give 25 calculated lines from 843 to 1649 Å with an uncertainty of ± 0.0006 Å.

We have derived level values from the following sources: those of $4s^2 4p$ and $4s4p^2 \text{ } ^4P$ are from the wavelengths of Kaufman, Radziemski, and Andrew [1966], those of $4s^2 5s$, $6d$, and $7d$ are derived from the wavelengths of Kaufman and Ward [1966]. The levels of $4s4p^2 \text{ } ^2D$, $4s^2 5p$, $6p$, $4d$, $5d$, $4f$, $6f$, $5g$, $6g$, and $7g$ are derived from the interferometric measurements by Kaufman and Andrew [1963]. The remaining levels are taken from the extensive analysis by Shenstone [1963].

The observed g -values were reported by Andrew *et al.* [1967].

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Ge II

Configuration	Term	J	Level (cm^{-1})	g
4s ² 4p	2P°	1/2	0.000	
		3/2	1 767.357	
4s4p ²	4P	1/2	51 575.885	
		3/2	52 290.942	
		5/2	53 366.738	
4s ² 5s	2S	1/2	62 403.028	2.003
4s4p ²	2D	3/2	65 015.672	0.803
		5/2	65 184.738	1.197
4s ² 5p	2P°	1/2	79 006.853	0.668
		3/2	79 366.494	1.330
4s ² 4d	2D	3/2	80 836.880	0.800
		5/2	81 012.598	1.200
4s4p ²	2S	1/2	85 890.55	
4s4p ²	2P	1/2	91 016.2	
		3/2	92 122.7	
4s ² 6s	2S	1/2	94 784.381	1.996
4s ² 5d	2D	3/2	100 089.970	0.797
		5/2	100 130.812	1.201
4s ² 4f	2F°	7/2	100 317.283	1.143
		5/2	100 317.976	0.857
4s ² 6p	2P°	1/2	101 098.446	
		3/2	101 244.635	
4s ² 7s	2S	1/2	107 936.16	
4s ² 6d	2D	3/2	110 378.07	
		5/2	110 397.63	
4s ² 5f	2F°	7/2	110 504.58	
		5/2	110 505.22	
4s ² 5g	2G	7/2	110 868.453	
		9/2	110 868.486	
4s ² 7p	2P°	1/2	111 016.30	
		3/2	111 091.55	
4s ² 8s	2S	1/2	114 638.47	
4s ² 7d	2D	3/2	115 977.80	
		5/2	115 988.77	
4s ² 6f	2F°	7/2	116 041.231	
		5/2	116 041.81	
4s ² 6g	2G	7/2, 9/2	116 266.75	

Ge II - Continued

Configuration	Term	J	Level (cm ⁻¹)	g
4s ² 8p	2P°	1/2	116 361.78	
		3/2	116 406.61	
4s ² 9s	2S	1/2	118 523.24	
4s ² 8d	2D	3/2	119 339.11	
		5/2	119 345.75	
4s ² 7f	2F°	7/2	119 373.92	
		5/2	119 374.34	
4s ² 7g	2G	7/2	119 521.835	
		9/2	119 521.837	
4s ² 9p	2P°	1/2	119 577.60	
		3/2	119 607.77	
4s ² 10s	2S	1/2	120 977.05	
4s ² 9d	2D	3/2	121 511.36	
		5/2	121 515.82	
4s ² 8f	2F	7/2	121 532.26	
		5/2	121 532.56	
4s ² 8g	2G	7/2, 9/2	121 633.87	
4s ² 10p	2P°	1/2	121 662.87	
		3/2	121 685.74	
4s4p(3P°)5s	4P°	1/2	122 694.3	
		3/2	123 401.2	
		5/2	124 732.4	
4s ² 10d	2D	3/2	122 995.45	
		5/2	122 998.43	
4s ² 9f	2F°	7/2	123 008.73	
		5/2	123 008.96	
4s ² 9g	2G	7/2, 9/2	123 081.34	
4s ² 11p	2P°	3/2	123 102.60	
4s ² 10f	2F°	5/2, 7/2	124 062.8	
4s ² 10g	2G	7/2, 9/2	124 116.37	
4s ² 11f	2F°	5/2, 7/2	124 841.2	
4s ² 11g	2G	7/2, 9/2	124 881.91	
4s4p(3P°)4d	2D°	3/2	126 922.8	
		5/2	127 271.4	
4s4p(3P°)5s	2P°	1/2	127 081.8	
		3/2	127 916.3	

Ge II — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>
Ge III (¹ S ₀)	Limit		128 521.3	
4s 4p (³ P°)4d	⁴ F°	³ / ₂	130 412.8	
		⁵ / ₂	130 777.7	
		⁷ / ₂	131 348.9	
		⁹ / ₂	132 150	
4s 4p (³ P°)4d	⁴ D°	¹ / ₂	136 203.4	
		³ / ₂	136 235.7	
		⁵ / ₂	136 351.1	
		⁷ / ₂	137 064.7	
4s 4p (³ P°)4d	⁴ P°	⁵ / ₂	138 187.5	
		³ / ₂	138 567.0	
		¹ / ₂	138 761.8	
4s 4p (³ P°)5p	⁴ D	¹ / ₂	138 720.5?	
		³ / ₂	139 170?	
		⁵ / ₂	139 913?	
4p ³	⁴ S°	³ / ₂	139 760	
4s 4p (³ P°)5p	² P	¹ / ₂	140 016?	
		³ / ₂	140 752?	
4s 4p (³ P°)4d	² P°	¹ / ₂	140 432?	
		³ / ₂	140 983	
4s 4p (³ P°)5p	⁴ D	⁷ / ₂	141 179.70	
4s 4p (³ P°)4d	² F°	⁵ / ₂	141 279.2	
		⁷ / ₂	142 701.1	
4s 4p (³ P°)5p	⁴ P	¹ / ₂	141 452	
		³ / ₂	141 787	
		⁵ / ₂	142 443.9	
4p ³	² D°	³ / ₂	149 238.2	
		⁵ / ₂	149 393.1	
4p ³	² P°	³ / ₂	156 013.0?	
		¹ / ₂	156 083.8?	
4s 4p (³ P°)6s	⁴ P°	¹ / ₂	156 183.4	
		³ / ₂	156 942?	
		⁵ / ₂	158 353.4	
4s 4p (³ P°)5d	² D°	³ / ₂	160 611.5	
		⁵ / ₂	160 948.2	
4s 4p (³ P°)4f	² F°	⁵ / ₂	161 858	
4s 4p (³ P°)4f	⁴ F°	⁷ / ₂	162 324?	
		⁹ / ₂	162 454?	
4s 4p (³ P°)4f	⁴ G	⁵ / ₂	163 220?	
		⁷ / ₂	163 573?	
		⁹ / ₂	164 013?	
		¹¹ / ₂	164 535	

Ge II — Continued

Configuration	Term	J	Level (cm ⁻¹)	g
4s4p(³ P°)4f	² G	⁷ / ₂	164 157	
		⁹ / ₂	164 735	
4s4p(³ P°)4f	⁴ D°	⁷ / ₂	164 258	
		⁵ / ₂	164 606?	
		³ / ₂	164 768?	
		¹ / ₂	164 860?	
Ge III 4s4p(³ P°)	<i>Limit</i>		190 250	

Ge III

Z = 32

Zn I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 4s^2 \ ^1S_0$ Ionization energy $274\,693 \pm 10 \text{ cm}^{-1}$ ($34.0576 \pm 0.0010 \text{ eV}$)

The analysis was first carried out by Lang (1929) who classified 86 lines in the range of 540–5300 Å measured with an uncertainty of about ± 0.05 Å. He determined values for levels of the configurations $4s4p$, $4p^2$, $4s4d$, $4s5s$, $4s5p$, $4p4d$, $4s4f$, $4s5d$, $4s6s$, and $4d^2$.

The spectrum was reobserved by Ryabstev *et al.* [1993] in the range of 300–7000 Å with an uncertainty of ± 0.005 Å below 6000 Å and ± 0.02 Å above 6000 Å. All level values were redetermined with improved accuracy. The previous levels $4p^2 \ ^1D$, $4s5s \ ^1S$ and all $4p4d$ and $4d^2$ levels were discarded. The $4s4d \ ^1D_2$ was reassigned to $4p^2 \ ^1D_2$. All of these levels except for the $^1F^\circ$ of the $4p4d$ configuration and the $^1P^\circ$ of the $4p5s$ were established. In addition, the $4sns$, np , nd , nf , and ng series were extended. A tentative assignment of a line at 377.031 Å to the inner-shell excitation $3d^{10} 4s^2 \ ^1S_0 - 3d^9 4s^2 4p \ ^1P_1^\circ$ was given.

We have adopted the energy level values given by Ryabstev *et al.* as well as their calculated percentage composition of the levels. We quote their value for the ionization energy derived from the $4sns \ ^3S$ series.

The g -values were derived by Moore [1952] from observed Zeeman patterns given by van den Bosch and Klinkenberg [1941].

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Ge III

Configuration	Term	J	Level (cm^{-1})	g	Leading percentages
$4s^2$	1S	0	0.00		97 3 $4p^2 \ ^1S$
$4s4p$	$^3P^\circ$	0	61 733.82		100
		1	62 499.95		100
		2	64 143.85		100
$4s4p$	$^1P^\circ$	1	91 872.68		97
$4p^2$	1D	2	144 974.86		64 32 $4s4d \ ^1D$
$4p^2$	3P	0	147 685.45		99
		1	148 640.09		100
		2	150 371.74		97
$4s5s$	3S	1	158 565.02	1.99	100
$4s4d$	3D	1	162 840.25		100
		2	162 910.34		100
		3	163 016.95		100
$4s5s$	1S	0	163 532.53		98
$4p^2$	1S	0	174 113.91		94
$4s4d$	1D	2	178 140.35		65 30 $4p^2 \ ^1D$

Ge III — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	<i>g</i>	Leading percentages
4s5p	³ P°	0	181 859.17		100
		1	182 027.84	1.48	98
		2	182 487.04	1.49	100
4s5p	¹ P°	1	184 298.06		95 3 4p5s ¹ P°
4s4f	³ F°	2	210 445.99		96 4 4p4d ³ F°
		3	210 461.50		72 24 4s4f ¹ F°
		4	210 523.91		96 4 4p4d ³ F°
4s4f	¹ F°	3	210 536.58		73 24 4s4f ³ F°
4s6s	³ S	1	211 140.07	1.94	100
4s6s	¹ S	0	212 662.0		99
4s5d	³ D	1	213 123.16	0.50	100
		2	213 152.45	1.16	100
		3	213 197.64	1.32	100
4s5d	¹ D	2	215 576.96		95 3 4p ² ¹ D
4s7s	³ S	1	234 432.3		100
4s5g	³ G	3	234 916.12		100
		4	234 916.38		51 48 4s5g ¹ G
		5	234 920.10		100
4s5g	¹ G	4	234 920.04		51 48 4s5g ³ G
4s7s	¹ S	0	235 099.8		100
4s6d	³ D	1	235 436.5		100
		2	235 452.2		100
		3	235 479.5		100
4s6d	¹ D	2	236 248.5		98
4p4d	¹ D°	2	237 851.3		56 21 4p4d ³ F°
4s8s	³ S	1	246 875.0		100
4s6g	³ G	3	247 083.70		100
		4	247 084.01		72 28 4s6g ¹ G
		5	247 087.25		100
4s6g	¹ G	4	247 087.28		72 28 4s6g ³ G
4p4d	¹ P°	1	262 715.2		94 3 3p5s ¹ P°
Ge IV (² S _{1/2})	<i>Limit</i>		274 693		

Ge IV

Z = 32

Cu I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 4s^2 S_{1/2}$ Ionization energy $368\,720 \pm 10 \text{ cm}^{-1}$ ($45.7155 \pm 0.0010 \text{ eV}$)

The spectrum was first observed by Carroll [1925] who classified nine lines, from which he derived the $4s$, $4p$, $4d$, and $5s$ term values. Improved wavelengths were given by Lang [1929] who extended the term system by locating $3d^{10}6p$, $3d^{10}5g$, and $3d^9 4s^2$ terms. He identified lines in the range of 440 – 3676 \AA .

The spectrum was reobserved by Ryabtsev and Wyart [1987]. They classified 56 additional lines in the range of 301 – 1229 \AA with a wavelength uncertainty that we estimate to be $\pm 0.005 \text{ \AA}$ for most lines. Several lines originating from autoionizing levels of $3d^9 4p^2$ are very broad and probably have an uncertainty of $\pm 0.05 \text{ \AA}$. With these new data they found 16 new levels, extending the nl series to $9s$, $8p$, $6d$, and $6f$. They also replaced the $6s$ and $6p$ levels of Lang, which they found to be incorrect, and redetermined values for all the levels except $4f^2 F^\circ$, $5g^2 G$, and $3d^9 4s^2 {}^2D$. With the ns series they determined the value for the ionization energy.

In addition Ryabtsev and Wyart identified the $3d^{10} 4p - 3d^9 4p^2$ array in the range of 301 – 338 \AA . Sixteen levels of $3d^9 4p^2$ were found, all of which are above the first ionization threshold. The percentage compositions for these levels are given, derived with configuration interaction with $3d^9 4s 4d$.

The transitions $3d^{10} 4s - 3d^9 4s 4p$ were identified by Wyart *et al.* [1984] in the range of 314 – 354 \AA . They

classified eight lines from $3d^9 4s 4p$ levels with $J = 1/2$ and $3/2$, and gave the first component of the percentage composition for these levels. Ryabtsev and Wyart corrected the transition $3d^{10} 4s {}^2S_{1/2} - 3d^9 ({}^2D) 4s 4p ({}^3P^\circ) {}^4P_{3/2}$ to 354.968 \AA and added the transition $3d^{10} 4s {}^2S_{1/2} - 3d^9 ({}^2D) 4s 4p ({}^3P^\circ) {}^4P_{1/2}$ at 352.102 \AA .

We quote the results of Ryabtsev and Wyart for all the $3d^{10} nl$ levels, with the exception of $4f^2 F^\circ$ and $5g^2 G$ taken from Lang's work, and those of $3d^9 4p^2$. The $3d^9 4s 4p$ levels are from Wyart *et al.* with the additions from Ryabtsev and Wyart. The $3d^9 4s^2 {}^2D$ levels are taken from Lang's paper.

The g -values were derived by Moore [1952] from the Zeeman patterns observed by van den Bosch and Klinkenberg [1941].

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Ge IV

Configuration	Term	J	Level (cm ⁻¹)	g	Leading percentages
$3d^{10} 4s$	2S	$1/2$	0		
$3d^{10} 4p$	${}^2P^\circ$	$1/2$	81 311.4		
		$3/2$	84 102.3		
$3d^{10} 4d$	2D	$3/2$	190 601.5	0.82	
		$5/2$	190 852.5	1.26	
$3d^{10} 5s$	2S	$1/2$	199 263.7	2.04	
$3d^{10} 5p$	${}^2P^\circ$	$1/2$	226 456.7	0.68	
		$3/2$	227 393.9	1.34	
$3d^{10} 4f$	${}^2F^\circ$	$7/2$	257 492		
		$5/2$	257 496		

Ge IV — Continued

Configuration	Term	J	Level (cm^{-1})	g	Leading percentages
$3d^9(^2D)4s^2$	2D	$5/2$	259 938		
		$3/2$	264 444		
$3d^{10}5d$	2D	$3/2$	267 455.5	0.80	
		$5/2$	267 575.6	1.19	
$3d^{10}6s$	2S	$1/2$	270 603.4		
$3d^9(^2D)4s4p(^3P^\circ)$	$^4P^\circ$	$3/2$	281 715		94
		$1/2$	284 009		97
$3d^{10}6p$	$^2P^\circ$	$1/2$	283 027.8		
		$3/2$	283 461.4		
$3d^9(^2D)4s4p(^3P^\circ)$	$^4D^\circ$	$3/2$	291 652		43
		$1/2$	293 876		75
$3d^9(^2D)4s4p(^3P^\circ)$	$^2P^\circ$	$1/2$	292 755		73
		$3/2$	293 306		83
$3d^9(^2D)4s4p(^3P^\circ)$	$^2D^\circ$	$3/2$	295 255		58
$3d^{10}5f$	$^2F^\circ$	$5/2$	297 563.6		
		$7/2$	297 564.2		
$3d^{10}5g$	2G	$7/2, 9/2$	298 373		
$3d^{10}6d$	2D	$3/2$	303 105.8		
		$5/2$	303 168.5		
$3d^{10}7s$	2S	$1/2$	304 667.6		
$3d^{10}7p$	$^2P^\circ$	$1/2$	311 342		
		$3/2$	311 482		
$3d^9(^2D)4s4p(^1P^\circ)$	$^2P^\circ$	$3/2$	318 124		97
		$1/2$	322 139		99
$3d^{10}6f$	$^2F^\circ$	$7/2$	319 277.1		
		$5/2$	319 277.9		
$3d^{10}8s$	2S	$1/2$	323 605.5		
$3d^{10}9s$	2S	$1/2$	335 228		
.....					
Ge V (1S_0)	Limit		368 720		
$3d^9(^2D)4p^2(^1D)$	2S	$1/2$	379 514		58
$3d^9(^2D)4p^2(^1D)$	2P	$3/2$	380 814		59
		$1/2$	384 514		41
$3d^9(^2D)4p^2(^3P)$	4D	$3/2$	381 702		82
		$1/2$	383 260		75
$3d^9(^2D)4p^2(^1D)$	2D	$3/2$	387 565		39

Ge IV — Continued

Configuration	Term	J	Level (cm^{-1})	g	Leading percentages
$3d^9(^2D)4p^2(^3P)$	4P	$5/2$	390 374		86
		$3/2$	393 376		81
		$1/2$	394 870		76
$3d^9(^2D)4p^2(^3P)$	2D	$3/2$	392 040		64
		$5/2$	394 170		58
$3d^9(^2D)4p^2(^3P)$		$5/2$	393 460		
$3d^9(^2D)4p^2(^3P)$	2P	$1/2$	396 620		76
		$3/2$	397 366		83
$3d^9(^2D)4p^2(^1S)$	2D	$5/2$	415 660?		92

Ge v

Z = 32

Ni I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^{10} 1S_0$ Ionization energy $729\,930 \pm 150 \text{ cm}^{-1}$ ($94.500 \pm 0.02 \text{ eV}$)

Sixteen levels comprising all of $3d^9 4s$ and $3d^9 4p$ were determined by Mack *et al.* [1928] from observations by Carroll [1925]. Their values relative to the ground state were found by Kruger and Shoupp [1934] who identified three resonance lines. The wavelengths in the range of 942–1222 Å were measured with improved accuracy by Joshi and van Kleef [1974]. They redetermined all the energy level values and added the $3d^9 4p \ ^3P_1^\circ$ level.

The spectrum was reobserved by Ryabtsev [1975] and Ryabtsev and Churilov [1991] in the range of 178–1222 Å with an uncertainty of $\pm 0.005 \text{ Å}$. They found all the levels of $3d^9 4d$, $4f$, $5s$ to $7s$, $5p$ and $6p$, and part of $5d$ and $6d$, as well as improved values for the pre-

viously known levels. Their results are quoted here along with their value for the ionization energy.

We have calculated the percentage composition of all the levels and designated them accordingly.

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Ge v

Configuration	Term	J	Level (cm ⁻¹)	Leading percentages	
$3d^{10}$	$1S$	0	0		
$3d^9 4s$	$3D$	3	234 219.3	100	
		2	235 967.0	83	17 $3d^9 4s \ ^1D$
		1	238 764.9	100	
$3d^9 4s$	$1D$	2	241 935.2	83	17 $3d^9 4s \ ^3D$
$3d^9 4p$	$3P^\circ$	2	323 748.9	96	4 $3d^9 4p \ ^3D^\circ$
		1	327 753.0	96	3 "
		0	330 332.5	100	
$3d^9 4p$	$3F^\circ$	3	327 891.0	68	27 $3d^9 4p \ ^1F^\circ$
		4	329 847.8	100	
		2	330 790.6	95	4 $3d^9 4p \ ^3D^\circ$
$3d^9 4p$	$1F^\circ$	3	335 161.3	60	27 $3d^9 4p \ ^3D^\circ$
$3d^9 4p$	$1D^\circ$	2	335 560.3	64	32 $3d^9 4p \ ^3D^\circ$
		3	337 168.1	68	19 $3d^9 4p \ ^3F^\circ$
		1	339 540.2	72	24 $3d^9 4p \ ^1P^\circ$
$3d^9 4p$	$3D^\circ$	2	340 295.7	60	36 $3d^9 4p \ ^1D^\circ$
		1	338 273.5	75	25 $3d^9 4p \ ^3D^\circ$
		3	456 051.8	93	5 $3d^9 4d \ ^3P$
$3d^9 4d$	$3S$	1	456 051.8	93	
$3d^9 4d$	$3G$	5	461 417.5	100	
		4	461 642.6	64	36 $3d^9 4d \ ^1G$
		3	464 077.0	68	19 $3d^9 4d \ ^1F$

Ge v - Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages			
3d ⁹ 4d	¹ P	1	461 814.9	52	28	3d ⁹ 4d ³ P	
3d ⁹ 4d	³ P	2	461 828.5	76	24	3d ⁹ 4d ³ D	
		0	463 957.9	94	6	3d ⁹ 4d ¹ S	
		1	464 852.7	57	36	3d ⁹ 4d ¹ P	
3d ⁹ 4d	³ D	3	463 361.6	89	9	3d ⁹ 4d ³ F	
		2	464 652.2	50	24	3d ⁹ 4d ¹ D	
		1	467 029.1	80	10	3d ⁹ 4d ¹ P	
3d ⁹ 4d		4	464 705.7	45	3d ⁹ 4d ³ F	39	3d ⁹ 4d ¹ G
3d ⁹ 4d	¹ F	3	466 780.2	42	31	3d ⁹ 4d ³ G	
3d ⁹ 4d	³ F	4	467 386.7	54	25	3d ⁹ 4d ¹ G	
		3	469 686.4	53	38	3d ⁹ 4d ¹ F	
		2	469 889.6	64	35	3d ⁹ 4d ¹ D	
3d ⁹ 4d	¹ D	2	468 695.8	41	27	3d ⁹ 4d ³ D	
3d ⁹ 5s	^(5/2, 1/2)	3	490 741.7	100	or	100	3d ⁹ 5s ³ D
		2	491 443.3	98	or	55	"
3d ⁹ 4d	¹ S	0	493 765.5?	94	6	3d ⁹ 4d ³ P	
3d ⁹ 5s	^(3/2, 1/2)	1	495 288.4	100	or	100	3d ⁹ 5s ³ D
		2	495 907.7	98	or	55	3d ⁹ 5s ¹ D
3d ⁹ 5p	³ P°	2	522 959	93	7	3d ⁹ 5p ³ D°	
		1	525 343	72	22	3d ⁹ 5p ¹ P°	
		0	528 156	100			
3d ⁹ 5p	³ F°	3	524 232	57	36	3d ⁹ 5p ¹ F°	
		4	524 950	100			
		2	528 625	71	26	3d ⁹ 5p ¹ D°	
3d ⁹ 5p	¹ D°	2	526 373	42	29	3d ⁹ 5p ³ D°	
3d ⁹ 5p	³ D°	3	526 902	80	20	3d ⁹ 5p ¹ F°	
		1	530 928	94	5	3d ⁹ 5p ³ P°	
		2	531 321	61	32	3d ⁹ 5p ¹ D°	
3d ⁹ 5p	¹ P°	1	529 476	77	22	3d ⁹ 5p ³ P°	
3d ⁹ 5p	¹ F°	3	530 113	44	42	3d ⁹ 5p ³ F°	
3d ⁹ 4f	^(5/2, 5/2) °	0	[553 660]	100			
		1	553 910.9	75	24	3d ⁹ 4f ^(5/2, 7/2) °	
		5	554 690.2	95	5	"	
		2	555 298.8	77	23	"	
		4	555 860.0	86	12	"	
		3	555 912.7	72	28	"	
3d ⁹ 4f	^(5/2, 7/2) °	2	554 370.9	75	22	3d ⁹ 4f ^(5/2, 5/2) °	
		6	554 658.0	100			
		3	555 337.1	72	28	"	
		1	555 730.2	72	25	"	
		4	555 798.4	87	13	"	
		5	555 852.3	93	5	"	

Ge v - Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages	
3d ⁹ 4f	(3/2, 7/2) ^o	2	558 877.1	89	8 3d ⁹ 4f (3/2, 5/2) ^o
		5	559 467.6	98	2 3d ⁹ 4f (5/2, 7/2) ^o
		3	560 096.6	96	4 3d ⁹ 4f (3/2, 5/2) ^o
		4	560 547.2	95	5 "
3d ⁹ 4f	(3/2, 5/2) ^o	4	559 463.1	93	5 3d ⁹ 4f (3/2, 7/2) ^o
		2	560 034.6	91	8 "
		1	560 149.3	96	4 3d ⁹ 4f (5/2, 7/2) ^o
		3	560 587.7	96	4 3d ⁹ 4f (3/2, 7/2) ^o
3d ⁹ 5d	³ S	1	574 389	79	15 3d ⁹ 5d ³ P
3d ⁹ 5d	³ G	5	576 067	100	
		4	576 226	55	44 3d ⁹ 5d ¹ G
		3	580 695	78	14 3d ⁹ 5d ¹ F
3d ⁹ 5d	³ D	3	576 963	84	13 3d ⁹ 5d ³ F
3d ⁹ 5d	¹ F	3	577 459	48	29 3d ⁹ 5d ³ F
3d ⁹ 5d	³ F	4	577 553	80	14 3d ⁹ 5d ¹ G
		3	582 225	50	36 3d ⁹ 5d ¹ F
		2	582 308	68	30 3d ⁹ 5d ¹ D
3d ⁹ 5d	¹ G	4	581 114	41	40 3d ⁹ 5d ³ G
3d ⁹ 6s	(5/2, 1/2)	3	588 094		
		2	588 403		
3d ⁹ 6s	(3/2, 1/2)	1	592 644		
		2	592 872		
3d ⁹ 6p	(5/2, 1/2) ^o	2	603 388	73	26 3d ⁹ 6p (5/2, 3/2) ^o
		3	603 921	95	5 "
3d ⁹ 6p	(5/2, 3/2) ^o	4	604 256	100	
		1	604 656	97	3 3d ⁹ 6p (3/2, 3/2) ^o
		2	604 965	73	26 3d ⁹ 6p (5/2, 1/2) ^o
		3	605 176	95	4 "
3d ⁹ 6p	(3/2, 3/2) ^o	0	608 172	100	
		3	609 029	100	
		1	609 384	57	42 3d ⁹ 6p (3/2, 1/2) ^o
		2	609 581	93	7 3d ⁹ 6p (1/2, 1/2) ^o
3d ⁹ 6p	(3/2, 1/2) ^o	2	608 262	93	7 3d ⁹ 6p (3/2, 3/2) ^o
		1	608 530	56	40 "
3d ⁹ 7s	(5/2, 1/2)	3	635 973		
		2	636 143		
		1	640 509		
3d ⁹ 7s	(3/2, 1/2)	2	640 625		
Ge VI (² D _{5/2})	Limit		729 930		

Ge vi

Z = 32

Co I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^9 \ ^2D_{5/2}$ Ionization energy $934\,800 \pm 1\,000 \text{ cm}^{-1}$ ($115.9 \pm 0.1 \text{ eV}$)

The $3d^9-3d^8 4p$ array was observed in the range of 222–242 Å by Kononov [1967], who measured the spectrum with a wavelength uncertainty of ± 0.03 Å. The ground state 2D splitting and 21 levels of $3d^8 4p$ were determined. Aksenov and Ryabtsev [1975] extended the range of observation to include 200–1100 Å and classified 110 lines in the arrays $3d^9-3d^8 4p$ and $3d^8 4s-3d^8 4p$. Their measurement uncertainty is ± 0.005 Å. An improved value for the 2D ground state splitting and all levels of $3d^8 4s$ and $3d^8 4p$ were found except for those based on the $3d^8 \ ^1S$ parent. Ramonas and Ryabtsev [1979] classified an additional 20 lines and found the missing levels of the above configurations. They also calculated the percentage composition of the levels. All energy level values were redetermined. Their results are quoted here.

New observations in the range of 128–148 Å by Ryabtsev *et al.* [1985] with an uncertainty of ± 0.005 Å were interpreted as the transitions $3d^9-3d^8 4f$ and

$3p^6 3d^9-3p^5 3d^{10}$. Percentage compositions were calculated for the mixed configurations $3d^8 4f + 3p^5 3d^{10}$. We quote these results.

We determined the value for the ionization energy from the center of gravity of the $3d^8 4f$ configuration and an interpolated value for its term energy, and the center of gravity of the $3d^8$ configuration calculated with the Cowan [1981] HFR program.

References

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Ge vi

Configuration	Term	J	Level (cm ⁻¹)	Leading percentages	
$3p^6 3d^9$	2D	$5/2$	0		
		$3/2$	4 560		
$3p^6 3d^8(^3F)4s$	4F	$9/2$	303 696.2	100	
		$7/2$	306 243.4	92	7 $3p^6 3d^8(^3F)4s \ ^2F$
		$5/2$	308 657.2	98	
		$3/2$	310 199.1	98	
$3p^6 3d^8(^3F)4s$	2F	$7/2$	313 024.6	92	7 $3p^6 3d^8(^3F)4s \ ^4F$
		$5/2$	316 936.8	94	4 $3p^6 3d^8(^1D)4s \ ^2D$
$3p^6 3d^8(^3P)4s$	4P	$5/2$	327 537.9	55	42 $3p^6 3d^8(^1D)4s \ ^2D$
		$3/2$	332 376.6	79	19 "
		$1/2$	332 497.0	100	
$3p^6 3d^8(^1D)4s$	2D	$3/2$	329 073.6	70	10 $3p^6 3d^8(^3P)4s \ ^4P$
		$5/2$	333 034.6	53	45 "
$3p^6 3d^8(^3P)4s$	2P	$3/2$	339 335.3	92	7 $3p^6 3d^8(^1D)4s \ ^2D$
		$1/2$	340 525.3	100	
$3p^6 3d^8(^1G)4s$	2G	$9/2$	343 624.4	100	
		$7/2$	343 674.0	100	
$3p^6 3d^8(^1S)4s$	2S	$1/2$	391 183.2	100	

Ge VI — Continued

Configuration	Term	J	Level (cm ⁻¹)	Leading percentages	
3p ⁶ 3d ⁸ (³ F)4p	4D°	7/2	404 892.3	88	6 3p ⁶ 3d ⁸ (³ F)4p 4F°
		5/2	408 981.1	86	7 "
		3/2	412 037.7	86	8 3p ⁶ 3d ⁸ (³ P)4p 4D°
		1/2	413 872.6	88	9 "
3p ⁶ 3d ⁸ (³ F)4p	4G°	9/2	409 188.4	53	27 3p ⁶ 3d ⁸ (³ F)4p 2G°
		11/2	411 591.6	100	
		7/2	411 885.9	74	14 3p ⁶ 3d ⁸ (³ F)4p 4F°
		5/2	413 728.0	86	8 "
3p ⁶ 3d ⁸ (³ F)4p	4F°	9/2	415 142.7	77	16 3p ⁶ 3d ⁸ (³ F)4p 2G°
		7/2	416 709.7	56	10 3p ⁶ 3d ⁸ (³ F)4p 2F°
		3/2	417 792.2	62	21 3p ⁶ 3d ⁸ (³ F)4p 2D°
		5/2	417 941.9	61	20 "
3p ⁶ 3d ⁸ (³ F)4p	2G°	9/2	417 421.2	56	40 3p ⁶ 3d ⁸ (³ F)4p 4G°
		7/2	421 310.0	62	14 3p ⁶ 3d ⁸ (³ F)4p 2F°
3p ⁶ 3d ⁸ (³ F)4p	2D°	5/2	420 117.8	58	15 3p ⁶ 3d ⁸ (³ F)4p 2F°
		3/2	423 030.0	53	31 3p ⁶ 3d ⁸ (³ F)4p 4F°
3p ⁶ 3d ⁸ (³ F)4p	2F°	7/2	420 542.0	55	24 3p ⁶ 3d ⁸ (³ F)4p 2G°
		5/2	424 506.4	72	12 3p ⁶ 3d ⁸ (³ F)4p 2D°
3p ⁶ 3d ⁸ (³ P)4p	4P°	3/2	429 997.3	72	9 3p ⁶ 3d ⁸ (¹ D)4p 2P°
		5/2	430 736.0	66	12 3p ⁶ 3d ⁸ (¹ D)4p 2D°
		1/2	431 041.9	90	5 3p ⁶ 3d ⁸ (¹ D)4p 2P°
3p ⁶ 3d ⁸ (¹ D)4p	2F°	5/2	433 507.0	69	19 3p ⁶ 3d ⁸ (³ P)4p 4P°
		7/2	436 173.3	70	14 3p ⁶ 3d ⁸ (¹ G)4p 2F°
3p ⁶ 3d ⁸ (¹ D)4p	2D°	3/2	436 023.6	40	19 3p ⁶ 3d ⁸ (³ P)4p 4P°
		5/2	437 801.0	74	8 "
3p ⁶ 3d ⁸ (¹ D)4p	2P°	1/2	436 727.0	58	29 3p ⁶ 3d ⁸ (³ P)4p 2P°
		3/2	438 871.4	46	34 3p ⁶ 3d ⁸ (¹ D)4p 2D°
3p ⁶ 3d ⁸ (³ P)4p	4D°	1/2	441 221.0	86	8 3p ⁶ 3d ⁸ (³ F)4p 4D°
		3/2	441 278.0	74	7 3p ⁶ 3d ⁸ (¹ D)4p 2P°
		5/2	441 386.0	59	19 3p ⁶ 3d ⁸ (³ P)4p 2D°
		7/2	442 790.4	76	17 3p ⁶ 3d ⁸ (¹ G)4p 2F°
3p ⁶ 3d ⁸ (¹ G)4p	2H°	9/2	444 442.2	98	
		11/2	448 258.9	100	
3p ⁶ 3d ⁸ (³ P)4p	2P°	3/2	445 439.0	72	13 3p ⁶ 3d ⁸ (¹ D)4p 2P°
		1/2	449 093.0	58	22 "
3p ⁶ 3d ⁸ (³ P)4p	2D°	5/2	445 670.0	66	27 3p ⁶ 3d ⁸ (³ P)4p 4D°
		3/2	447 373.0	83	10 "
3p ⁶ 3d ⁸ (¹ G)4p	2F°	7/2	447 780.0	64	22 3p ⁶ 3d ⁸ (¹ D)4p 2D°
		5/2	450 153.0	85	8 3p ⁶ 3d ⁸ (¹ D)4p 2F°
3p ⁶ 3d ⁸ (³ P)4p	2S°	1/2	451 362.0	83	9 3p ⁶ 3d ⁸ (¹ D)4p 2P°
3p ⁶ 3d ⁸ (³ P)4p	4S°	3/2	451 501.0	96	

Ge VI — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages	
$3p^6 3d^8(^1G)4p$	$^2G^\circ$	$7/2$	458 378.0	96	
		$9/2$	459 014.5	98	
$3p^6 3d^8(^1S)4p$	$^2P^\circ$	$1/2$	496 179.5	98	
		$3/2$	500 575.3	98	
$3p^6 3d^8(^3F)4f$	$^2S^\circ$	$1/2$	682 720	45	23 $3p^6 3d^8(^3F)4f^2P^\circ$
$3p^6 3d^8(^3F)4f$	$^4D^\circ$	$7/2$	682 980	61	35 $3p^6 3d^8(^3F)4f^4F^\circ$
$3p^6 3d^8(^3F)4f$	$^2P^\circ$	$3/2$	683 460	45	20 $3p^6 3d^8(^3F)4f^4D^\circ$
		$1/2$	689 590	48	22 $3p^6 3d^8(^3F)4f^2G^\circ$
$3p^6 3d^8 4f$		$7/2$	684 530	34	$3p^6 3d^8(^3F)4f^2F^\circ$ 21 $3p^6 3d^8(^3F)4f^4G^\circ$
$3p^6 3d^8 4f$		$5/2$	684 730	30	$3p^6 3d^8(^3F)4f^4F^\circ$ 26 $3p^6 3d^8(^3F)4f^2D^\circ$
$3p^6 3d^8(^3F)4f$	$^4P^\circ$	$1/2$	686 810	58	27 $3p^6 3d^8(^3F)4f^2P^\circ$
$3p^6 3d^8 4f$		$3/2$	687 320	30	$3p^6 3d^8(^3F)4f^2P^\circ$ 26 $3p^6 3d^8(^3F)4f^4D^\circ$
$3p^6 3d^8 4f$		$5/2$	687 690	34	$3p^6 3d^8(^3F)4f^4D^\circ$ 27 $3p^6 3d^8(^3F)4f^4P^\circ$
$3p^6 3d^8 4f$		$3/2$	687 750	32	$3p^6 3d^8(^3F)4f^4S^\circ$ 30 $3p^6 3d^8(^3F)4f^4F^\circ$
$3p^6 3d^8 4f$		$7/2$	688 330	33	$3p^6 3d^8(^3F)4f^4G^\circ$ 32 $3p^6 3d^8(^3F)4f^2F^\circ$
$3p^6 3d^8(^3F)4f$	$^2D^\circ$	$5/2$	688 460	41	31 $3p^6 3d^8(^3F)4f^4F^\circ$
		$3/2$	691 020	67	14 $3p^6 3d^8(^3F)4f^2P^\circ$
$3p^6 3d^8(^3F)4f$	$^2G^\circ$	$7/2$	690 680	45	31 $3p^6 3d^8(^3F)4f^2F^\circ$
$3p^6 3d^8(^3F)4f$	$^2F^\circ$	$5/2$	691 420	53	31 $3p^6 3d^8(^3F)4f^2D^\circ$
$3p^6 3d^8(^1D)4f$	$^2F^\circ$	$5/2$	705 420	57	16 $3p^6 3d^8(^3P)4f^2F^\circ$
		$7/2$	706 410	45	28 $3p^6 3d^8(^1D)4f^2G^\circ$
$3p^6 3d^8(^1D)4f$	$^2D^\circ$	$3/2$	706 300	50	23 $3p^6 3d^8(^3P)4f^4F^\circ$
		$5/2$	706 650	64	15 "
$3p^6 3d^8(^1D)4f$	$^2P^\circ$	$1/2$	706 490	73	24 $3p^6 3d^8(^3P)4f^4D^\circ$
		$3/2$	706 850	63	17 "
$3p^6 3d^8(^3P)4f$	$^4F^\circ$	$5/2$	710 830	67	17 $3p^6 3d^8(^1D)4f^2D^\circ$
		$3/2$	710 940	72	21 "
$3p^6 3d^8(^3P)4f$	$^2F^\circ$	$7/2$	711 920	67	19 $3p^6 3d^8(^1D)4f^2G^\circ$
$3p^6 3d^8(^3P)4f$	$^4G^\circ$	$7/2$	712 940	81	12 $3p^6 3d^8(^3P)4f^2G^\circ$
$3p^6 3d^8 4f$		$5/2$	713 230	32	$3p^6 3d^8(^3P)4f^2F^\circ$ 25 $3p^6 3d^8(^3P)4f^4D^\circ$
$3p^6 3d^8(^3P)4f$	$^4D^\circ$	$3/2$	713 250	66	14 $3p^6 3d^8(^1D)4f^2P^\circ$
		$5/2$	713 650	47	27 $3p^6 3d^8(^3P)4f^2F^\circ$
		$7/2$	714 030	71	12 $3p^6 3d^8(^3P)4f^2G^\circ$
$3p^6 3d^8(^3P)4f$	$^2G^\circ$	$7/2$	713 518	67	20 $3p^6 3d^8(^3P)4f^4D^\circ$

Ge VI — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages	
$3p^6 3d^8(^3P)4f$	$^2D^\circ$	$3/2$	714 100	67	16 $3p^6 3d^8(^1D)4f^2 D^\circ$
		$5/2$	714 210	79	14 $3p^6 3d^8(^3P)4f^4 D^\circ$
$3p^6 3d^8(^1G)4f$	$^2P^\circ$	$3/2$	715 650	79	8 $3p^6 3d^{10} ^2P^\circ$
		$1/2$	717 595	95	3 "
$3p^6 3d^8(^1G)4f$	$^2D^\circ$	$5/2$	720 370	98	
		$3/2$	720 550	98	
$3p^6 3d^8(^1G)4f$	$^2F^\circ$	$7/2$	721 400	99	
		$5/2$	721 470	98	
$3p^5 3d^{10}$	$^2P^\circ$	$3/2$	753 930	87	11 $3p^6 3d^8(^1G)4f^2 P^\circ$
		$1/2$	783 600	95	3 "
$3p^6 3d^8(^1S)4f$	$^2F^\circ$	$7/2$	768 990	99	
Ge VII (3F_4)	<i>Limit</i>		934 800		

Ge VII

Z = 32

Fe I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^8 \ ^3F_4$ Ionization energy $1\ 169\ 000 \pm 12\ 000\ \text{cm}^{-1}$ ($144.9 \pm 1.5\ \text{eV}$)

The spectrum was obtained with a spark discharge and recorded in the range of 160–210 Å with an uncertainty of $\pm 0.004\ \text{Å}$ by Podobedova *et al.* [1980]. They classified 220 lines in the array $3d^8 - 3d^7 4p$. All the levels of the ground configuration and most of $3d^7 4p$ were determined with an uncertainty of $\pm 10\ \text{cm}^{-1}$.

New observations obtained by Podobedova and Ryabtsev [1986] enabled them to classify lines of the $3p^6 3d^8 - 3p^5 3d^9$ array in the range of 123 to 142 Å and the $3d^7 4s - 3d^7 4p$ in the range of 739–987 Å as well as improving the earlier work by Podobedova *et al.* [1980]. They also give the percentage composition of all the levels.

We obtained the levels and percentages for the $3d^8$ configuration from the earlier Podobedova *et al.* [1980]

paper and the remaining levels from the later Podobedova and Ryabtsev [1986] paper.

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

References

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 Podobedova, L. I., and Ryabtsev, A. N., [1986] *Opt. Spectrosc. (USSR)* **60**, 694.

Ge VII

Configuration	Term	J	Level (cm^{-1})	Leading percentages	
$3p^6 3d^8$	3F	4	0	100	
		3	4 055	100	
		2	6 468	96	
$3p^6 3d^8$	1D	2	22 658	67	30 $3p^6 3d^8 \ ^3P$
$3p^6 3d^8$	3P	2	28 827	70	29 $3p^6 3d^8 \ ^1D$
		1	29 212	100	
		0	29 769	100	
$3p^6 3d^8$	1G	4	37 923	100	
$3p^6 3d^8$	1S	0	86 203	100	
$3p^6 3d^7(^4F)4s$	5F	5	393 964	99	
		4	396 659	98	
		3	398 813	98	
		2	400 338	99	
		1	401 310	99	
$3p^6 3d^7(^4F)4s$	3F	4	405 833	97	
		3	409 481	98	
		2	411 818	98	
$3p^6 3d^7(^4F)4s$	5P	3	421 833	99	
		2	422 209	88	12 $3p^6 3d^7(^2P)4s \ ^3P$
		1	423 790	98	

Ge VII -- Continued

Configuration	Term	J	Level (cm ⁻¹)	Leading percentages	
3p ⁶ 3d ⁷ (² G)4s	³ G	5	427 035	94	
		4	428 431	86	
		3	430 455	99	
3p ⁶ 3d ⁷ (⁴ P)4s	³ P	2	432 065	47	38 3p ⁶ 3d ⁷ (² P)4s ³ P
		1	432 097	44	24 3p ⁶ 3d ⁷ (² P)4s ¹ P
		0	438 647?	55	45 3p ⁶ 3d ⁷ (² P)4s ³ P
3p ⁶ 3d ⁷ (² P)4s	³ P	0	433 275?	55	45 3p ⁶ 3d ⁷ (⁴ P)4s ³ P
		2	434 645	51	35 3p ⁶ 3d ⁷ (² P)4s ³ P
		1	435 977	52	35 "
3p ⁶ 3d ⁷ (² G)4s	¹ G	4	434 261	80	11 3p ⁶ 3d ⁷ (² G)4s ³ G
3p ⁶ 3d ⁷ (² H)4s	³ H	6	437 438	100	
		5	438 931	90	
		4	441 167	87	11 3p ⁶ 3d ⁷ (² G)4s ¹ G
3p ⁶ 3d ⁷ (² D2)4s	³ D	3	438 698	76	23 3p ⁶ 3d ⁷ (² D1)4s ³ D
		2	440 920	57	15 "
		1	445 585	43	37 3p ⁶ 3d ⁷ (² P)4s ¹ P
3p ⁶ 3d ⁷ 4s		1	439 551	38	3p ⁶ 3d ⁷ (² P)4s ¹ P 26 3p ⁶ 3d ⁷ (² D1)4s ³ D
3p ⁶ 3d ⁷ (² H)4s	¹ H	5	444 819	91	
3p ⁶ 3d ⁷ (² D2)4s	¹ D	2	447 235	61	15 3p ⁶ 3d ⁷ (² D1)4s ¹ D
3p ⁶ 3d ⁷ (² F)4s	³ F	2	462 416	99	
		3	462 884	98	
		4	463 729	99	
3p ⁶ 3d ⁷ (² F)4s	¹ F	3	469 133	98	
3p ⁶ 3d ⁷ (² D1)4s	³ D	1	498 640	81	18 3p ⁶ 3d ⁷ (² D2)4s ³ D
		2	499 206	78	19 "
		3	500 364	76	23 "
3p ⁶ 3d ⁷ (² D1)4s	¹ D	2	505 648	76	21 3p ⁶ 3d ⁷ (² D2)4s ¹ D
3p ⁶ 3d ⁷ (⁴ F)4p	⁵ D°	4	506 743	47	41 3p ⁶ 3d ⁷ (⁴ F)4p ⁵ F°
		3	514 478	42	34 3p ⁶ 3d ⁷ (⁴ F)4p ⁵ G°
		1	517 882	75	16 3p ⁶ 3d ⁷ (⁴ P)4p ⁵ D°
		0	518 296?	81	18 "
3p ⁶ 3d ⁷ (⁴ F)4p	⁵ F°	5	507 521	80	13 3p ⁶ 3d ⁷ (⁴ F)4p ⁵ G°
		3	509 634	60	30 3p ⁶ 3d ⁷ (⁴ F)4p ⁵ D°
		2	511 950	74	17 "
		1	513 576	87	
3p ⁶ 3d ⁷ 4p		4	511 843	35	3p ⁶ 3d ⁷ (⁴ F)4p ⁵ D° 33 3p ⁶ 3d ⁷ (⁴ F)4p ⁵ F°
3p ⁶ 3d ⁷ (⁴ F)4p	⁵ G°	0	514 281	99	
		5	514 626	56	23 3p ⁶ 3d ⁷ (⁴ F)4p ³ G°
		4	516 182	56	21 3p ⁶ 3d ⁷ (⁴ F)4p ⁵ F°
		2	516 277	56	32 3p ⁶ 3d ⁷ (⁴ F)4p ⁵ D°
		3	517 123	51	23 3p ⁶ 3d ⁷ (⁴ F)4p ⁵ F°
3p ⁶ 3d ⁷ 4p		2	517 526	35	3p ⁶ 3d ⁷ (⁴ F)4p ⁵ D° 33 3p ⁶ 3d ⁷ (⁴ F)4p ⁵ G°

Ge VII — Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages	
$3p^6 3d^7(^4F)4p$	$^3G^\circ$	5	520 237	69	30 $3p^6 3d^7(^4F)4p ^5G^\circ$
		4	524 060	83	14 "
		3	526 696	90	
$3p^6 3d^7(^4F)4p$	$^3F^\circ$	4	520 617	81	
		3	523 344	63	21 $3p^6 3d^7(^4F)4p ^3D^\circ$
		2	525 786	60	23 $3p^6 3d^7(^4P)4p ^5S^\circ$
$3p^6 3d^7(^4F)4p$	3D	3	525 123?	72	16 $3p^6 3d^7(^4F)4p ^3F$
		2	527 177	73	
		1	528 653	83	
$3p^6 3d^7(^4P)4p$	$^5S^\circ$	2	525 200	67	20 $3p^6 3d^7(^4F)4p ^3F^\circ$
$3p^6 3d^7 4p$		1	537 900	33	$3p^6 3d^7(^2P)4p ^3P^\circ$ 20 $3p^6 3d^7(^4P)4p ^3S^\circ$
$3p^6 3d^7(^4P)4p$	$^5D^\circ$	2	539 685	50	11 $3p^6 3d^7(^4F)4p ^5D^\circ$
		3	540 316	63	13 $3p^6 3d^7(^2P)4p ^3P^\circ$
		1	541 541	58	13 $3p^6 3d^7(^4P)4p ^3S^\circ$
		4	542 862	86	
$3p^6 3d^7(^2G)4p$	$^3H^\circ$	5	540 170	56	20 $3p^6 3d^7(^2G)4p ^1H^\circ$
		4	542 376	79	
		6	544 249	89	
$3p^6 3d^7(^2G)4p$	$^3F^\circ$	4	540 217	53	13 $3p^6 3d^7(^2G)4p ^3G^\circ$
		3	544 670	53	28 "
		2	549 137	90	
$3p^6 3d^7 4p$		0	540 556?	36	$3p^6 3d^7(^4P)4p ^5D^\circ$ 16 $3p^6 3d^7(^4F)4p ^5D^\circ$
$3p^6 3d^7(^2P)4p$	$^3P^\circ$	0	543 430?	47	27 $3p^6 3d^7(^4P)4p ^5D^\circ$
		2	543 913	50	10 "
$3p^6 3d^7 4p$		1	544 951	31	$3p^6 3d^7(^2P)4p ^3P^\circ$ 24 $3p^6 3d^7(^4P)4p ^5P^\circ$
$3p^6 3d^7 4p$		2	545 592	26	$3p^6 3d^7(^4P)4p ^5P^\circ$ 21 $3p^6 3d^7(^2P)4p ^1D^\circ$
$3p^6 3d^7 4p$		4	546 178	39	$3p^6 3d^7(^2G)4p ^1G^\circ$ 29 $3p^6 3d^7(^2G)4p ^3F^\circ$
$3p^6 3d^7(^4P)4p$	$^5P^\circ$	3	546 636	42	18 $3p^6 3d^7(^4P)4p ^3D^\circ$
		2	547 728	44	13 "
$3p^6 3d^7(^2G)4p$	$^3G^\circ$	5	547 345	76	13 $3p^6 3d^7(^2G)4p ^1H^\circ$
		4	549 426	68	14 $3p^6 3d^7(^2G)4p ^3H^\circ$
		3	549 742	44	18 $3p^6 3d^7(^2G)4p ^3F^\circ$
$3p^6 3d^7 4p$		1	548 195	24	$3p^6 3d^7(^4P)4p ^5P^\circ$ 23 $3p^6 3d^7(^4P)4p ^3S^\circ$
$3p^6 3d^7(^2H)4p$	$^3G^\circ$	5	548 225	84	
		4	551 889	81	
		3	554 987	63	
$3p^6 3d^7(^4P)4p$	$^3D^\circ$	3	548 808	44	21 $3p^6 3d^7(^4P)4p ^5P^\circ$
$3p^6 3d^7 4p$		1	549 121	33	$3p^6 3d^7(^4P)4p ^5P^\circ$ 24 $3p^6 3d^7(^2P^\circ)4p ^3D^\circ$
$3p^6 3d^7(^2P)4p$	$^1S^\circ$	0	549 503?	65	23 $3p^6 3d^7(^4P)4p ^3P^\circ$

Ge VII - Continued

Configuration	Term	J	Level (cm ⁻¹)	Leading percentages	
3p ⁶ 3d ⁷ (² H)4p	³ I°	6	549 592	58	35 3p ⁶ 3d ⁷ (² H)4p ¹ I°
		5	552 176	83	10 3p ⁶ 3d ⁷ (² G)4p ¹ H°
		7	553 185	100	
3p ⁶ 3d ⁷ (² G)4p	¹ H°	5	549 675	50	37 3p ⁶ 3d ⁷ (² G)4p ³ H°
3p ⁶ 3d ⁷ 4p		2	551 188	36	3p ⁶ 3d ⁷ (⁴ P)4p ³ P° 25 3p ⁶ 3d ⁷ (⁴ P)4p ³ D°
3p ⁶ 3d ⁷ (² D2)4p	³ D°	3	551 831	47	15 3p ⁶ 3d ⁷ (² D1)4p ³ D°
3p ⁶ 3d ⁷ 4p		1	552 495	35	3p ⁶ 3d ⁷ (⁴ P)4p ³ D° 23 3p ⁶ 3d ⁷ (⁴ P)4p ³ P°
3p ⁶ 3d ⁷ 4p		3	552 823	35	3p ⁶ 3d ⁷ (² P)4p ³ D° 27 3p ⁶ 3d ⁷ (² D2)4p ³ F°
3p ⁶ 3d ⁷ 4p		2	553 030	32	3p ⁶ 3d ⁷ (⁴ P)4p ³ P° 16 3p ⁶ 3d ⁷ (² D2)4p ³ F°
3p ⁶ 3d ⁷ 4p		2	553 730	29	3p ⁶ 3d ⁷ (² D2)4p ³ D° 19 3p ⁶ 3d ⁷ (⁴ P)4p ³ D°
3p ⁶ 3d ⁷ 4p		1	553 741	37	3p ⁶ 3d ⁷ (² D2)4p ³ D° 23 3p ⁶ 3d ⁷ (² P)4p ¹ P°
3p ⁶ 3d ⁷ (⁴ P)4p	³ P°	1	555 075	47	16 3p ⁶ 3d ⁷ (⁴ P)4p ³ D°
		0	558 187	69	30 3p ⁶ 3d ⁷ (² P)4p ¹ S°
3p ⁶ 3d ⁷ (² P)4p	³ D°	2	555 549	44	18 3p ⁶ 3d ⁷ (² P)4p ¹ D°
3p ⁶ 3d ⁷ (² D2)4p	³ F°	4	557 572	76	19 3p ⁶ 3d ⁷ (² D1)4p ³ F°
3p ⁶ 3d ⁷ (² H)4p	¹ I°	6	557 878	60	33 3p ⁶ 3d ⁷ (² H)4p ³ I°
3p ⁶ 3d ⁷ 4p		3	558 778	34	3p ⁶ 3d ⁷ (² D2)4p ³ F° 15 3p ⁶ 3d ⁷ (² G)4p ¹ F°
3p ⁶ 3d ⁷ 4p		2	559 843	37	3p ⁶ 3d ⁷ (² D2)4p ³ F° 16 3p ⁶ 3d ⁷ (² D2)4p ¹ D°
3p ⁶ 3d ⁷ (² P)4p	³ S°	1	559 953	60	11 3p ⁶ 3d ⁷ (² P)4p ³ P°
3p ⁶ 3d ⁷ 4p		2	560 486	30	3p ⁶ 3d ⁷ (² D2)4p ¹ D° 21 3p ⁶ 3d ⁷ (² D2)4p ³ P°
3p ⁶ 3d ⁷ 4p		1	560 894	39	3p ⁶ 3d ⁷ (² P)4p ¹ P° 11 3p ⁶ 3d ⁷ (² D2)4p ³ D°
3p ⁶ 3d ⁷ (² H)4p	³ H°	6	561 279	94	
		5	562 236	88	
		4	563 599	92	
3p ⁶ 3d ⁷ 4p		2	564 446	32	3p ⁶ 3d ⁷ (² D2)4p ³ P° 19 3p ⁶ 3d ⁷ (² D2)4p ¹ D°
3p ⁶ 3d ⁷ (² H)4p	¹ G°	4	564 577	61	35 3p ⁶ 3d ⁷ (² G)4p ¹ G°
3p ⁶ 3d ⁷ (² D2)4p	¹ F°	3	566 263	52	14 3p ⁶ 3d ⁷ (² G)4p ¹ F°
3p ⁶ 3d ⁷ (² D2)4p	¹ P°	1	567 391	40	20 3p ⁶ 3d ⁷ (² D2)4p ³ P°
3p ⁶ 3d ⁷ 4p		1	568 575	36	3p ⁶ 3d ⁷ (² D2)4p ¹ P° 30 3p ⁶ 3d ⁷ (² D2)4p ³ P°
3p ⁶ 3d ⁷ (² D2)4p	³ P°	0	568 862	68	14 3p ⁶ 3d ⁷ (² P)4p ³ P°
3p ⁶ 3d ⁷ (² H)4p	¹ H°	5	570 244	90	
3p ⁶ 3d ⁷ (² F)4p	¹ D°	2	577 022	47	38 3p ⁶ 3d ⁷ (² F)4p ³ F°

Ge VII - Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages	
$3p^6 3d^7 ({}^2F) 4p$	${}^3G^\circ$	3	578 696	63	19 $3p^6 3d^7 ({}^2F) 4p {}^3F^\circ$
		4	579 547	45	26 "
		5	584 341	92	
$3p^6 3d^7 ({}^2F) 4p$	${}^3D^\circ$	3	580 763	56	21 $3p^6 3d^7 ({}^2F) 4p {}^3F^\circ$
		2	583 647	68	19 $3p^6 3d^7 ({}^2F) 4p {}^1D^\circ$
		1	583 788	87	
$3p^6 3d^7 ({}^2F) 4p$	${}^3F^\circ$	2	582 542	52	27 $3p^6 3d^7 ({}^2F) 4p {}^1D^\circ$
		3	582 977	51	24 $3p^6 3d^7 ({}^2F) 4p {}^3D^\circ$
		4	584 123	42	34 $3p^6 3d^7 ({}^2F) 4p {}^1G^\circ$
$3p^6 3d^7 ({}^2F) 4p$	${}^1G^\circ$	4	584 712	45	48 $3p^6 3d^7 ({}^2F) 4p {}^3F^\circ$
$3p^6 3d^7 ({}^2F) 4p$	${}^1F^\circ$	3	592 619	92	
$3p^6 3d^7 ({}^2D1) 4p$	${}^3P^\circ$	2	609 898	75	20 $3p^6 3d^7 ({}^2D2) 4p {}^3P^\circ$
		1	610 851	76	15 "
		0	612 085	83	15 "
$3p^6 3d^7 ({}^2D1) 4p$	${}^3F^\circ$	2	612 681	74	20 $3p^6 3d^7 ({}^2D2) 4p {}^3F^\circ$
		3	614 904	69	18 "
		4	618 502	73	22 "
$3p^6 3d^7 ({}^2D1) 4p$	${}^1P^\circ$	1	620 557	70	10 $3p^6 3d^7 ({}^2D2) 4p {}^1P^\circ$
$3p^6 3d^7 ({}^2D1) 4p$	${}^1F^\circ$	3	621 422	69	18 $3p^6 3d^7 ({}^2D2) 4p {}^1F^\circ$
$3p^6 3d^7 ({}^2D1) 4p$	${}^3D^\circ$	1	625 836	65	19 $3p^6 3d^7 ({}^2D2) 4p {}^3D^\circ$
		2	626 496	64	20 "
		3	629 033	66	24 "
$3p^5 3d^9$	${}^3F^\circ$	4	728 432	100	
		3	742 467	91	
		2	759 175	53	22 $3p^5 3d^9 {}^3D^\circ$
$3p^5 3d^9$	${}^1D^\circ$	2	733 491	77	18 $3p^5 3d^9 {}^3F^\circ$
$3p^5 3d^9$	${}^3P^\circ$	1	769 022	81	19 $3p^5 3d^9 {}^3D^\circ$
		0	769 854?	100	
		2	794 440	53	43 "
$3p^5 3d^9$	${}^3D^\circ$	3	772 711	90	
		1	784 470	75	16 $3p^5 3d^9 {}^3P^\circ$
$3p^5 3d^9$		2	775 405	34	$3p^5 3d^9 {}^3D^\circ$ 28 $3p^5 3d^9 {}^3F^\circ$
$3p^5 3d^9$	${}^1F^\circ$	3	845 547	97	
$3p^5 3d^9$	${}^1P^\circ$	1	848 346?	91	
Ge VIII (${}^4F_{9/2}$)	Limit		1 169 000		

Ge VIII

Z = 32

Mn I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^7 \ ^4F_{9/2}$ Ionization energy $1\,423\,000 \pm 14\,000 \text{ cm}^{-1}$ ($176.4 \pm 1.7 \text{ eV}$)

No classified lines have been reported for this spectrum. The energy levels of the $3d^7$ configuration have been predicted by Wyart *et al.* [1985] in a parametric study of the isoelectronic sequence with an uncertainty of a few hundred cm^{-1} .

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

References

- Cowan, R. D. [1981], *The Theory of Atomic Structure and Spectra*, (Univ. California Press, Berkeley, CA).
 Wyart, J. -F., Raassen, A. J. J., and Uylings, P. H. M. [1985], *Phys. Scr.* **32**, 169.

Ge IX

Z = 32

Cr I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^6 \ ^5D_4$ Ionization energy $1\,714\,000 \pm 17\,000 \text{ cm}^{-1}$ ($212.7 \pm 2.1 \text{ eV}$)

No classified lines have been reported for this spectrum. We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

Reference

- Cowan, R. D. [1981], *The Theory of Atomic Structure and Spectra*, (Univ. California Press, Berkeley, CA).

Ge X

Z = 32

V I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^5 \ ^6S_{5/2}$ Ionization energy $2\,030\,000 \pm 20\,000 \text{ cm}^{-1}$ ($251.7 \pm 2.5 \text{ eV}$)

No classified lines have been reported for this spectrum. We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

Reference

- Cowan, R. D. [1981], *The Theory of Atomic Structure and Spectra*, (Univ. California Press, Berkeley, CA).

Ge XI

 $Z = 32$

Ti I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^4 \ ^5D_0$ Ionization energy $2\,304\,000 \pm 23\,000 \text{ cm}^{-1}$ ($285.7 \pm 2.8 \text{ eV}$)

No classified lines have been reported for this spectrum. We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

Reference

Cowan, R. D. [1981], *The Theory of Atomic Structure and Spectra*, (Univ. California Press, Berkeley, CA).

Ge XII

 $Z = 32$

Sc I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^3 \ ^4F_{3/2}$ Ionization energy $2\,634\,000 \pm 26\,000 \text{ cm}^{-1}$ ($326.6 \pm 3.2 \text{ eV}$)

No lines have been classified in this spectrum. Wyart *et al.* [1985] have predicted the values for the energy levels of the $3d^3$ ground configuration and their percentage composition in a study of the isoelectronic sequence with an uncertainty of a few hundred cm^{-1} .

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

References

Cowan, R. D. [1981], *The Theory of Atomic Structure and Spectra*, (Univ. California Press, Berkeley, CA).
Wyart, J. -F., Raassen, A. J. J., and Uylings, P. H. M. [1985], *Phys. Scr.* **32**, 169.

Ge XIII

 $Z = 32$

Ca I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^2 \ ^3F_2$ Ionization energy $2\,960\,000 \pm 30\,000 \text{ cm}^{-1}$ ($368 \pm 4 \text{ eV}$)

No observations have been reported for this spectrum. Energy levels values for the $3d^2$ ground configuration were predicted by Wyart *et al.* [1985] with an uncertainty of a few hundred cm^{-1} in an isoelectronic study.

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

References

Cowan, R. D. [1981], *The Theory of Atomic Structure and Spectra*, (Univ. California Press, Berkeley, CA).
Wyart, J. F., Raassen, A. J. J., and Uylings, P. H. M. [1985], *Phys. Scr.* **32**, 169.

Ge XIV

Z = 32

K I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 3d^2 D_{3/2}$ Ionization energy $3\,300\,000 \pm 16\,000 \text{ cm}^{-1}$ ($409 \pm 2 \text{ eV}$)

Three transitions of the array $3p^6 3d - 3p^5 3d^2$ were observed by Fawcett and Hayes [1975] with an uncertainty of $\pm 0.03 \text{ \AA}$ in a laser-produced plasma. Kaufman *et al.* [1989] observed this spectrum in a tokamak plasma with an uncertainty of $\pm 0.005 \text{ \AA}$ and classified three additional lines of this array. They derived the energy levels from the observed transitions and obtained an interpolated value for the 2D ground term splitting of $7714 \pm 5 \text{ cm}^{-1}$. The $3p^5(^2P^\circ)3d^2(^3P) ^2P_{1/2}$ level value was predicted from an isoelectronic interpolation.

The value for the ionization energy was calculated by Finkelnburg and Humbach [1955] by extrapolation of the effective charge on the residual ion.

References

- Fawcett, B. C., and Hayes, R. W. [1975], *J. Opt. Soc. Am.* **65**, 623.
 Finkelnburg, W., and Humbach, W. [1955], *Naturwiss.* **42**, 35.
 Kaufman, V., Sugar, J., and Rowan, W. L. [1989], *J. Opt. Soc. Am. B* **6**, 142.

Ge XIV

Configuration	Term	J	Level (cm ⁻¹)
$3s^2 3p^6 3d$	2D	$3/2$	0
		$5/2$	7 714
$3s^2 3p^5(^2P^\circ)3d^2(^3F)$	$^2F^\circ$	$5/2$	800 830
$3s^2 3p^5(^2P^\circ)3d^2(^1G)$	$^2F^\circ$	$7/2$	822 250
$3s^2 3p^5(^2P^\circ)3d^2(^3P)$	$^2P^\circ$	$1/2$	[877 550]
		$3/2$	887 740
$3s^2 3p^5(^2P^\circ)3d^2(^3F)$	$^2D^\circ$	$3/2$	885 420
		$5/2$	885 710
.....			
Ge xv ($1S_0$)	Limit		3 300 000

Ge xv

 $Z = 32$

Ar I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^6 \ ^1S_0$ Ionization energy $4\,290\,000 \pm 8000 \text{ cm}^{-1}$ ($532 \pm 1 \text{ eV}$)

The resonance transition $3p^6 \ ^1S_0 - 3p^5 3d \ ^1P_1^o$ was first observed by Fawcett and Hayes [1975] at $117.25 \pm 0.03 \text{ \AA}$ by means of a laser-produced plasma. It was later measured by Sugar *et al.* [1987] in the spectrum of a tokamak plasma at $117.222 \pm 0.01 \text{ \AA}$. They also reported a value of $145.829 \pm 0.01 \text{ \AA}$ for the transition $3p^6 \ ^1S_0 - 3p^5 3d \ ^3D_1^o$.

The value for the ionization energy was calculated by Finkelnburg and Humbach [1955] by extrapolation of the effective charge on the residual ion.

References

- Fawcett, B. C., and Hayes, R. W. [1975], *J. Opt. Soc. Am.* **65**, 623.
 Finkelnburg, W., and Humbach, W. [1955], *Naturwiss.* **42**, 35.
 Sugar, J., Kaufman, V., and Rowan, W. L. [1987], *J. Opt. Soc. Am. B* **4**, 1927.

Ge xv

Configuration	Term	J	Level (cm^{-1})
$3s^2 3p^6$	1S	0	0
$3s^2 3p^5 3d$	$^3D^o$	1	685 735
$3s^2 3p^5 3d$	$^1P^o$	1	853 082
.....			
Ge xvi ($^2P_{3/2}$)	<i>Limit</i>		4 290 000

Ge XVI

 $Z = 32$

Cl I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^5 {}^2P_{3/2}^{\circ}$ Ionization energy $4\,570\,000 \pm 16\,000 \text{ cm}^{-1}$ ($567 \pm 2 \text{ eV}$)

The interval of the $3p^5 {}^2P^{\circ}$ ground term is determined by the ${}^2P_{1/2}^{\circ} - {}^2P_{3/2}^{\circ}$ M1 transition observed by Denne *et al.* [1983] with a tokamak plasma at $2085.1 \pm 0.1 \text{ \AA}$.

Kaufman *et al.* [1989] observed five lines of the $3s^2 3p^5 - 3s^2 3p^4 3d$ array with a tokamak plasma. They fall in the range of $114 - 126 \text{ \AA}$ and have a measurement uncertainty of $\pm 0.005 \text{ \AA}$. Corrected values for the sequence $3s^2 3p^5 {}^2P_{1/2}^{\circ} - 3s^2 3p^4 ({}^3P) 3d {}^2P_{1/2}$ were given by Kaufman *et al.* [1990]. We quote the energy levels from the two papers of Kaufman *et al.* The level value 838 414 is given in brackets because it was obtained by isoelectronic interpolation of the corresponding wavelength. We estimate the wavelength uncertainty to be $\pm 0.010 \text{ \AA}$.

The value for the ionization energy was calculated by Finkelnburg and Humbach [1955] by extrapolation of the effective charge on the residual ion. This was given in their paper as 367 eV but it is a misprint, as shown by our calculation with the Cowan HFR code giving a value of 568 eV.

References

- Cowan, R. D. [1981], *The Theory of Atomic Structure and Spectra*, (Univ. California Press, Berkeley, CA).
- Denne, B., Hinnov, E., Suckewer, S., and Cohen, S. [1983], *Phys. Rev. A* **28**, 206.
- Finkelnburg, W., and Humbach, W. [1955], *Naturwiss.* **42**, 35.
- Kaufman, V., Sugar, J., and Rowan, W. L. [1989], *J. Opt. Soc. Am. B* **6**, 1444.
- Kaufman, V., Sugar, J., and Rowan, W. L. [1990], *J. Opt. Soc. Am. B* **7**, 1169.

Ge XVI

Configuration	Term	J	Level (cm^{-1})
$3s^2 3p^5$	${}^2P^{\circ}$	$3/2$	0
		$1/2$	47 944
$3s^2 3p^4 ({}^1D) 3d$	2S	$1/2$	790 233
$3s^2 3p^4 ({}^3P) 3d$	2P	$3/2$	820 499
		$1/2$	[838 414]
$3s^2 3p^4 ({}^3P) 3d$	2D	$5/2$	827 801
		$3/2$	869 839
.....
Ge XVII (3P_2)	Limit		4 570 000

Ge xvii

Z = 32

S I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^4 \ ^3P_2$ Ionization energy $4\,920\,000 \pm 8000 \text{ cm}^{-1}$ ($610 \pm 1 \text{ eV}$)

Denne [1983] observed three M1 lines in a tokamak plasma due to transitions among the $3s^2 3p^4$ levels. Their measurement uncertainty is $\pm 0.3 \text{ \AA}$. These lines determine all levels of $3s^2 3p^4$ except for 3P_0 , which was derived by Kaufman *et al.* [1990] by isoelectronic fitting at $33\,290 \pm 30 \text{ cm}^{-1}$. Level values were derived from these lines with an uncertainty of ± 5 to $\pm 30 \text{ cm}^{-1}$ for the 1S_0 level.

Kaufman *et al.* [1990] observed ten lines of the transition array $3s^2 3p^4 - 3s^2 3p^3 3d$ in the range of $125 - 143 \text{ \AA}$ and three lines of $3s^2 3p^4 - 3s 3p^5$ at $218 - 231 \text{ \AA}$. Their measurement uncertainty was $\pm 0.007 \text{ \AA}$. Their calculated percentage compositions of the levels of $3s^2 3p^4$ are

given here. We calculated the compositions of the mixed $3s 3p^5$ and $3s^2 3p^3 3d$ configurations.

The value for the ionization energy was calculated by Finkelnburg and Humbach [1955] by extrapolation of the effective charge on the residual ion.

References

- Denne, B., Hinnov, E., Suckewer, S., and Cohen, S. [1983], *Phys. Rev. A* **28**, 206.
 Finkelnburg, W., and Humbach, W. [1955], *Naturwiss.* **42**, 35.
 Kaufman, V., Sugar, J., and Rowan, W. L. [1990], *J. Opt. Soc. Am. B* **7**, 1169.

Ge xvii

Configuration	Term	J	Level (cm ⁻¹)	Leading percentages	
$3s^2 3p^4$	3P	2	0	88	12 $3s^2 3p^4 \ ^1D$
		0	[33 290]	78	22 $3s^2 3p^4 \ ^1S$
		1	41 535	100	
$3s^2 3p^4$	1D	2	73 461	88	12 $3s^2 3p^4 \ ^3P$
$3s^2 3p^4$	1S	0	146 478	78	22 $3s^2 3p^4 \ ^3P$
$3s 3p^5$	$^3P^\circ$	2	431 810	82	12 $3s^2 3p^3(^2D^\circ)3d \ ^3P^\circ$
		1	457 340	76	12 "
$3s^2 3p^3(^2P^\circ)3d$	$^3P^\circ$	2	743 200	71	9 $3s^2 3p^3(^2P^\circ)3d \ ^1D^\circ$
$3s^2 3p^3(^2D^\circ)3d$	$^3P^\circ$	2	773 690	78	12 $3s 3p^5 \ ^3P^\circ$
$3s^2 3p^3(^2P^\circ)3d$	$^1F^\circ$	3	782 600	46	31 $3s^2 3p^3(^2D^\circ)3d \ ^1F^\circ$
$3s^2 3p^3(^4S^\circ)3d$	$^3D^\circ$	3	799 240	45	20 $3s^2 3p^3(^2D^\circ)3d \ ^1F^\circ$
		2	817 680	28	22 $3s^2 3p^3(^2P^\circ)3d \ ^3D^\circ$
$3s^2 3p^3(^2D^\circ)3d$	$^1D^\circ$	2	851 540	50	17 $3s^2 3p^3(^2P^\circ)3d \ ^1D^\circ$
$3s^2 3p^3(^2D^\circ)3d$	$^1F^\circ$	3	867 340	58	32 $3s^2 3p^3(^2P^\circ)3d \ ^1F^\circ$
$3s^2 3p^3(^2P^\circ)3d$	$^1P^\circ$	1	915 030	80	7 $3s^2 3p^3(^4S^\circ)3d \ ^3D^\circ$
Ge xviii ($^4S_{3/2}$)	Limit		4 920 000		

Ge xviii

Z = 32

P I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^3 \ ^4S_{3/2}$ Ionization energy $5\,390\,000 \pm 32\,000 \text{ cm}^{-1}$ ($668 \pm 4 \text{ eV}$)

All levels of the ground configuration are determined with the M1 transitions observed by Denne *et al.* [1984] in a tokamak discharge. Except for the $^4S_{3/2} - ^2P_{3/2}$ line at 703.6 Å, all the M1 lines are tentatively assigned. Their measurement uncertainty is $\pm 0.2 \text{ Å}$, giving a level uncertainty of $\pm 40 \text{ cm}^{-1}$.

The transition array $3p^3 - 3p^2 3d$ was observed in a tokamak plasma by Sugar *et al.* [1991], who classified seven lines in the range of 131–137 Å measured with an uncertainty of $\pm 0.005 \text{ Å}$. We quote their levels. They also calculated the percentage composition of the $3s^2 3p^3$ levels. We have calculated the percentages for the interacting $3s 3p^4$ and $3s^2 3p^2 3d$ configurations. The two level

values given in parentheses were obtained by Sugar *et al.* by interpolation of observed minus calculated energies with an uncertainty of $\pm 0.01 \text{ Å}$.

The value for the ionization energy was calculated by Finkelnburg and Humbach [1955] by extrapolation of the effective charge on the residual ion.

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Ge xviii

Configuration	Term	J	Level (cm ⁻¹)	Leading percentages	
$3s^2 3p^3$	$^4S^{\circ}$	$3/2$	0	86	11 $3s^2 3p^3 \ ^2P^{\circ}$
$3s^2 3p^3$	$^2D^{\circ}$	$3/2$	56 117	75	16 $3s^2 3p^3 \ ^2P^{\circ}$
		$5/2$	73 646	100	
$3s^2 3p^3$	$^2P^{\circ}$	$1/2$	112 334	100	22 $3s^2 3p^3 \ ^2D^{\circ}$
		$3/2$	142 132	73	
$3s^2 3p^2(^3P)3d$	4P	$5/2$	731 502	80	7 $3s 3p^4 \ ^4P$
		$3/2$	740 598	56	15 $3s^2 3p^2(^1D)3d \ ^2P$
$3s^2 3p^2(^1S)3d$	2D	$5/2$	805 800	51	8 $3s^2 3p^2(^1D)3d \ ^2F$
$3s^2 3p^2(^3P)3d$	2F	$7/2$	834 392	40	34 $3s^2 3p^2(^1D)3d \ ^2F$
$3s^2 3p^2(^1D)3d$	2P	$3/2$	[840 300]	46	27 $3s^2 3p^2(^3P)3d \ ^2P$
$3s^2 3p^2(^3P)3d$	2D	$5/2$	876 810	39	23 $3s^2 3p^2(^1S)3d \ ^2D$
Ge XIX (3P_0)	Limit		5 390 000		

Ge XIX

 $Z = 32$

Si I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^2 \ ^3P_0$ Ionization energy $5\,700\,000 \pm 8000 \text{ cm}^{-1}$ ($707 \pm 1 \text{ eV}$)

Measurements of M1 radiation in tokamak discharges provided sufficient information with which to determine all the energy levels of the $3p^2$ configuration. The interval $^3P_0 - ^3P_1$ and $^3P_1 - ^3P_2$ were observed by Denne *et al.* [1983] at $2933.7 \pm 0.2 \text{ \AA}$ and $5170.3 \pm 0.3 \text{ \AA}$, respectively. Datla *et al.* [1989] have given the $^3P_1 - ^1D_2$ and $^3P_2 - ^1D_2$ M1 transitions at $1341.2 \pm 0.5 \text{ \AA}$ and $1809.9 \pm 0.5 \text{ \AA}$. Hinnov [1985] reported the $^3P_1 - ^1S_0$ line at $746.9 \pm 0.3 \text{ \AA}$.

The transition $3p^2 \ ^3P_2 - 3p3d \ ^3D_3$ was reported by Fawcett and Hayes [1975]. With spectra of a laser-produced plasma and a tokamak discharge, Sugar *et al.* [1990] classified nine additional lines of this array and three lines of $3s^2 3p^2 - 3s3p^3$. Their wavelength uncertainty was $\pm 0.005 \text{ \AA}$. They also included the percentage composition of the levels.

New observations of this spectrum with a high energy laser-produced plasma by Ekberg *et al.* [1992] confirm all but the $3s3p^3 \ ^3P_2$ level in the results of Sugar *et al.* They

report three additional levels: the $3s3p^3 \ ^3P_1$, 3D_2 , and 3D_1 . We include their values for these levels and their calculated percentage compositions.

The value for the ionization energy was calculated by Finkelnburg and Humbach [1955] by extrapolation of the effective charge on the residual ion.

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Ge XIX

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages	
3s ² 3p ²	³ P	0	0	90	10 3s ² 3p ² ¹ S
		1	34 076.7	100	
		2	53 412.5	71	29 3s ² 3p ² ¹ D
3s ² 3p ²	¹ D	2	108 650	71	29 3s ² 3p ² ³ P
3s ² 3p ²	¹ S	0	167 970	90	10 3s ² 3p ² ³ P
3s 3p ³	³ D°	1	444 690	77	11 3s ² 3p 3d ³ D°
		2	446 159	75	10 "
3s 3p ³	³ P°	1	515 931	78	10 3s ² 3p 3d ³ P°
		2	521 554	59	14 3s 3p ³ ³ D°
3s 3p ³	³ S°	1	617 645	67	28 3s 3p ³ ¹ P°
3s 3p ³	¹ P°	1	674 263	55	29 3s 3p ³ ³ S°
3s ² 3p 3d	³ P°	2	704 381	50	16 3s ² 3p 3d ¹ D°
		1	752 767	60	30 3s ² 3p 3d ³ D°
3s ² 3p 3d	³ D°	1	716 266	57	29 3s ² 3p 3d ³ P°
		3	752 792	86	8 "
		2	759 430	43	37 "
3s ² 3p 3d	¹ D°	2	741 021	38	31 3s ² 3p 3d ³ D°
3s ² 3p 3d	¹ F°	3	817 036	95	4 3s ² 3p 3d ³ D°
3s ² 3p 3d	¹ P°	1	839 250	82	12 3s 3p ³ ¹ P°
Ge xx (² P _{1/2})	<i>Limit</i>		5 700 000		

Ge xx

 $Z = 32$

Al I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 3p^2 P_{1/2}^o$ Ionization energy $6\,030\,000 \pm 24\,000 \text{ cm}^{-1}$ ($748 \pm 3 \text{ eV}$)

The first observations were reported by Fawcett and Hayes [1975] who identified the $3s^2 3p - 3s^2 3d$ $^2P_{1/2}^o - ^2D_{3/2}$ and $^2P_{3/2}^o - ^2D_{5/2}$ lines as well as the transition $3s^2 3p^2 P_{3/2}^o - 3s^2 3p^2 P_{1/2}^o$ with a laser-produced plasma.

The spectrum was observed in a tokamak discharge by Hinnov *et al.* [1986]. They classified nine lines in the $3s^2 3p - 3s^2 3p^2$ and $3s^2 3p - 3s^2 3d$ arrays and gave wavelengths with an uncertainty of $\pm 0.05 \text{ \AA}$. The spectrum was reobserved by Sugar *et al.* [1988] who also used a tokamak discharge but obtained a wavelength uncertainty of $\pm 0.01 \text{ \AA}$ by means of photographic detection and an improved distribution of wavelength standards.

With spectra emitted by a laser-produced plasma in the range of $130 - 235 \text{ \AA}$ Ekberg *et al.* [1991] observed transitions from the $3p^3$ and $3s^2 3p^2 3d$ configurations and determined the positions of 16 of these levels. Their wavelength uncertainty is reported as $\pm 0.01 \text{ \AA}$, giving an average level uncertainty of $\pm 40 \text{ cm}^{-1}$. They also derived

the percentage composition of all the levels. The $3s^2 3p^2 P^o$ splitting is from the M1 line of Hinnov *et al.* Levels of $3s^2 3p^2$ and $3s^2 3d$ are those derived by Sugar *et al.*

The value for the ionization energy was calculated by Finkelnburg and Humbach [1955] by extrapolation of the effective charge on the residual ion.

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Ge xx

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages	
3s ² 3p	2P°	1/2	0	97	
		3/2	54 564	97	
3s3p ²	2D	3/2	453 869	83	12 3s ² 3d 2D
		5/2	466 407	76	11 "
3s3p ²	2P	1/2	530 521	48	47 3s3p ² 2S
		3/2	597 020	93	
3s3p ²	2S	1/2	585 257	48	49 3s3p ² 2P
3s ² 3d	2D	3/2	682 246	84	13 3s3p ² 2D
		5/2	689 533	86	12 "
3p ³	4S°	3/2	875 850	75	13 3p ³ 2D°
3s3p(3P°)3d	4P°	5/2	996 200	60	28 3s3p(3P°)3d 4D°
		1/2	1 035 630	90	10 "
3s3p(3P°)3d	4D°	3/2	1 000 790	62	32 3s3p(3P°)3d 4P°
		1/2	1 002 640	88	10 "
		7/2	1 033 390	95	
3s3p(3P°)3d	2D°	3/2	1 050 590	46	23 3s3p(1P°)3d 2D°
3s3p(3P°)3d	2F°	7/2	1 114 490	69	30 3s3p(1P°)3d 2F°
3s3p(3P°)3d	2P°	3/2	1 157 630	63	18 3s3p(1P°)3d 2P°
		1/2	1 184 750	86	10 3p ³ 2P°
3s3p(1P°)3d	2F°	7/2	1 184 190	66	30 3s3p(3P°)3d 2F°
		5/2	1 190 040	66	28 "
3s3p(1P°)3d	2P°	1/2	1 211 960	88	6 3p ³ 2P°
		3/2	1 223 700	73	9 3s3p(1P°)3d 2D°
3s3p(1P°)3d	2D°	3/2	1 217 450	44	18 3s3p(3P°)3d 2P°
		5/2	1 222 780	64	16 3s3p(3P°)3d 2D°
Ge XXI (1S ₀)	Limit		6 030 000		

Ge XXI

 $Z = 32$

Mg I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 \ ^1S_0$ Ionization energy $6\ 730\ 000 \pm 67\ 000\ \text{cm}^{-1}$ ($834 \pm 8\ \text{eV}$)

The first observations of this spectrum were reported by Fawcett and Hayes [1975] who identified the transitions $3s^2 \ ^1S_0 - 3s3p \ ^1P_1^o$, $3s3p \ ^3P_2^o - 3s3d \ ^3D_3$, and $3s3p \ ^1P_1^o - 3s3d \ ^1D_2$ with a laser-produced plasma. The spin-forbidden resonance line $3s^2 \ ^1S_0 - 3s3p \ ^3P_1^o$ and the M1 transition $3s3p \ ^3P_1^o - ^3P_2^o$ were both identified in spectra of tokamak discharges by Finkenthal *et al.* [1982] and by Denne *et al.* [1983], respectively. With a laser-generated plasma, Litzén and Redfors [1987] identified the $3s3d - 3p3d$ transitions $^3D - ^3F^o$ and $^1D - ^1F^o$.

The spectrum was observed more completely by Sugar *et al.* [1989] with both laser and tokamak excitations. They identified 38 lines in the range of 152–266 Å measured with an uncertainty of $\pm 0.005\ \text{Å}$. Their values for the $3s^2 - 3s3p$ resonance lines are 196.573 Å and 293.362 Å. Lines of the arrays $3s3p - 3p^2$, $3s3p - 3s3d$, $3s3d - 3p3d$, $3p^2 - 3p3d$, and $3p3d - 3d^2$ were classified.

With a line-focused high-powered laser plasma Ekberg *et al.* [1989] produced a more complete excitation and improved and extended the analysis. We report their values for the levels obtained with wavelength measurements

having an uncertainty of $\pm 0.01\ \text{Å}$. They also gave the percentage compositions of the levels with configuration interaction among all the even configurations $3s^2$, $3p^2$, $3d^2$, and $3s3d$ and between the odd configurations $3s3p$ and $3p3d$.

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

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Ge XXI

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages	
3s ²	¹ S	0	0	98	2 3p ² ¹ S
3s3p	³ P ^o	0	326 384	100	
		1	340 862	97	2 3s3p ¹ P ^o
		2	383 421	100	
3s3p	¹ P ^o	1	508 690	95	3 3s3p ³ P ^o
3p ²	³ P	0	777 547	91	8 3p ² ¹ S
		1	813 275	100	
		2	864 401	76	16 3p ² ¹ D
3p ²	¹ D	2	808 162	62	23 3p ² ³ P
3p ²	¹ S	0	963 328	88	9 3p ² ³ P
3s3d	³ D	1	965 838	100	
		2	969 563	100	
		3	975 540	100	
3s3d	¹ D	2	1 080 920	78	21 3p ² ¹ D
3p3d	³ F ^o	2	1 315 892	82	17 3p3d ¹ D ^o
		3	1 344 446	95	4 3p3d ³ D ^o
		4	1 380 973	100	
3p3d	¹ D ^o	2	1 361 297	65	15 3p3d ³ P ^o
3p3d	³ D ^o	1	1 392 766	75	21 3p3d ³ P ^o
		2	1 406 809	46	32 "
		3	1 431 589	94	4 3p3d ³ F ^o
3p3d	³ P ^o	0	1 435 471	100	
		1	1 435 827	77	23 3p3d ³ D ^o
		2	1 437 695	52	45 "
3p3d	¹ F ^o	3	1 521 245	97	2 3p3d ³ D ^o
3p3d	¹ P ^o	1	1 538 477	94	3 3p3d ³ D ^o
3d ²	³ F	2	1 952 219	99	1 3d ² ¹ D
		3	1 958 840	100	
		4	1 966 176	99	1 3d ² ¹ G
3d ²	¹ D	2	2 000 311	58	40 3d ² ³ P
3d ²	³ P	1	2 005 861	100	
		2	2 013 067	59	40 3d ² ¹ D
3d ²	¹ G	4	2 011 823	99	1 3d ² ³ F
Ge XXII (² S _{1/2})			Limit		6 730 000

Ge xxii

 $Z = 32$

Na I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 3s^2 S_{1/2}$ Ionization energy $7\,101\,700 \pm 500 \text{ cm}^{-1}$ ($880.50 \pm 0.06 \text{ eV}$)

The $3s - 3p$ and $3p - 3d$ doublets were first observed in a laser-induced plasma by Fawcett and Hayes [1975] with a wavelength uncertainty of 0.03 \AA . Using a similar light source Kononov *et al.* [1979] considerably extended the observed lines to include $3s - np$ ($n = 3 - 5$), $3p - nd$ ($n = 3 - 6$), $3p - 4s$, $3d - nf$ ($n = 4 - 6$), $3d - 4p$, $4p - 5d$, $4d - 5f$, and $4f - 5g$. The $3s^2 S_{1/2} - 3p^2 P_{1/2}$ line was not observed. The wavelength uncertainties are within $\pm 0.01 \text{ \AA}$ for all lines except for $3d - 6d$ and $3d - 7f$. Improved values are derived by Reader *et al.* [1987] for the $3s - 3p$, $3p - 3d$, and $3d - 4f$ transitions with uncertainties of $\pm 0.007 \text{ \AA}$, including the $3s^2 S_{1/2} - 3p^2 P_{1/2}$ line. Their values agree with Kononov *et al.* to within $\pm 0.005 \text{ \AA}$, except for the $3p^2 P_{1/2} - 3d^2 D_{3/2}$ line, which deviates by 0.011 \AA .

We have derived the values for the energy levels combining the data of Reader *et al.* with those of Kononov *et al.* We derived the value for the ionization energy from the nf series ($n = 4 - 6$), assuming a linear quantum defect dependence on the term value.

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Ge xxii

Configuration	Term	J	Level (cm^{-1})
3s	2S	$1/2$	0
3p	$^2P^\circ$	$1/2$	382 403
		$3/2$	441 501
3d	2D	$3/2$	955 830
		$5/2$	966 130
4s	2S	$1/2$	3 305 400
4p	$^2P^\circ$	$1/2$	3 460 300
		$3/2$	3 483 200
4d	2D	$3/2$	3 674 900
		$5/2$	3 679 300
4f	$^2F^\circ$	$5/2$	3 770 000
		$7/2$	3 771 800
5p	$^2P^\circ$	$1/2$	4 816 000
		$3/2$	4 826 700
5d	2D	$5/2$	4 921 000
		$7/2$	4 923 400
5f	$^2F^\circ$	$5/2$	4 968 800
		$7/2$	4 969 500
5g	2G	$7/2$	4 975 200
		$9/2$	4 975 800
6d	2D	$5/2$	5 593 000
6f	$^2F^\circ$	$5/2$	5 620 400
		$7/2$	5 621 400
7f	$^2F^\circ$	$7/2$	6 019 000
Ge xxiii (1S_0)	Limit		7 101 700

Ge xxiii

Z = 32

Ne I isoelectronic sequence

Ground state $1s^2 2s^2 2p^6 \ ^1S_0$ Ionization energy $17\,558\,000 \pm 18\,000 \text{ cm}^{-1}$ ($2176.9 \pm 2.2 \text{ eV}$)

Resonance lines in the range of 7–10 Å originating from $2p^5 nl$ configurations where $nl = 3s, 3d, 4s,$ and $4d$ as well as $2s 2p^6 3p$ were observed by Burkhalter *et al.* [1975] with a laser-generated plasma and a measurement uncertainty of $\pm 0.01 \text{ \AA}$. Gordon *et al.* [1980] extended the observations to $nl = 5d$ and $6d$, added $2s 2p^6 4p$, and improved the wavelength uncertainty to $\pm 0.005 \text{ \AA}$. We quote the results of Gordon *et al.*

By observing axially along a line-focused laser plasma, Lee *et al.* [1987] identified three lines of the array $2p^5 3s - 2p^5 3p$ from which we derive the $2p^5 3p$ levels.

The value for the ionization energy was derived by Hutcheon [1980] from the $2p^5 nd$ series.

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Ge xxiii

Configuration	Term	J	Level (cm^{-1})
$2s^2 2p^6$	1S	0	0
$2s^2 2p^5(^2P_{3/2}^o)3s$	$(^3/2, ^1/2)^o$	1	9 990 000
$2s^2 2p^5(^2P_{1/2}^o)3s$	$(^1/2, ^1/2)^o$	1	10 244 000
$2s^2 2p^5(^2P_{3/2}^o)3p$	$(^3/2, ^3/2)$	2	10 421 000
$2s^2 2p^5(^2P_{1/2}^o)3p$	$(^1/2, ^3/2)$	2	10 667 000
$2s^2 2p^5(^2P_{3/2}^o)3p$	$(^1/2, ^1/2)$	0	10 754 000
$2s^2 2p^5(^2P_{3/2}^o)3d$	$(^3/2, ^3/2)^o$	1	10 846 000
$2s^2 2p^5(^2P_{3/2}^o)3d$	$(^3/2, ^5/2)^o$	1	10 976 000
$2s^2 2p^5(^2P_{1/2}^o)3d$	$(^1/2, ^3/2)^o$	1	11 197 000
$2s 2p^6 3p$	$(^1/2, ^1/2)^o$	1	11 867 000
$2s 2p^6 3p$	$(^1/2, ^3/2)^o$	1	11 943 000
$2s^2 2p^5(^2P_{3/2}^o)4s$	$(^3/2, ^1/2)^o$	1	13 490 000
$2s^2 2p^5(^2P_{1/2}^o)4s$	$(^1/2, ^1/2)^o$	1	13 749 000
$2s^2 2p^5(^2P_{3/2}^o)4d$	$(^3/2, ^5/2)^o$	1	13 879 000
$2s^2 2p^5(^2P_{1/2}^o)4d$	$(^1/2, ^3/2)^o$	1	14 122 000
$2s^2 2p^5(^2P_{3/2}^o)5d$	$(^3/2, ^3/2)^o$	1	15 188 000
$2s 2p^6 4p$	$(^1/2, ^1/2)^o$	1	15 188 000
$2s^2 2p^5(^2P_{3/2}^o)5d$	$(^3/2, ^5/2)^o$	1	15 218 000
$2s 2p^6 4p$	$(^1/2, ^3/2)^o$	1	15 218 000
$2s^2 2p^5(^2P_{1/2}^o)5d$	$(^1/2, ^3/2)^o$	1	15 468 000
$2s^2 2p^5(^2P_{3/2}^o)6d$	$(^3/2, ^5/2)^o$	1	15 944 000
$2s^2 2p^5(^2P_{1/2}^o)6d$	$(^1/2, ^3/2)^o$	1	16 202 000
Ge xxiv ($^2P_{3/2}^o$)	Limit		17 558 000

Ge xxiv

Z = 32

F I isoelectronic sequence

Ground state $1s^2 2s^2 2p^5 \ ^2P_{3/2}^o$ Ionization energy $18\,580\,000 \pm 180\,000 \text{ cm}^{-1}$ ($2303 \pm 22 \text{ eV}$)

The ground state $^2P^o$ term interval was determined with the magnetic dipole transition observed at $379.5 \pm 0.1 \text{ \AA}$ in a tokamak plasma by Denne and Hinnov [1987]. The $2s\,2p^6 \ ^2S_{1/2}$ level was derived from the identification of the two resonance lines to the ground $^2P^o$ term in a laser plasma by Behring *et al.* [1976] and by Kononov *et al.* [1977]. Improved wavelengths with an uncertainty of $\pm 0.01 \text{ \AA}$ were obtained by Feldman *et al.* [1989] by observing along the axis of a line-focused laser plasma.

Burkhalter *et al.* [1977] classified lines in the range of $8.5\text{--}9.5 \text{ \AA}$ observed with a laser plasma. These are identified as transitions from $2p^4 3s$ and $2p^4 3d$ to the $2p^5$ ground term. Their wavelength uncertainty is given as $\pm 0.003 \text{ \AA}$. The same transition arrays were observed by Boiko *et al.* [1979] including nine additional classified lines with the same wavelength uncertainty of Burkhalter *et al.* Gordon *et al.* [1980] repeated the observations with an estimated uncertainty of $\pm 0.005 \text{ \AA}$ but extended the wavelength range to 6.7 \AA to include the upper configurations $2s\,2p^5 3p$, $2p^4 4s$, and $2p^4 4d$. We give the results of Gordon *et al.* and their calculated percentage compositions for the levels. No levels could be derived for the

$2p^4 4d$ configuration because of insufficient resolution of the spectrum.

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

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- Behring, W. E., Cohen, L., Doschek, G. A., and Feldman, U. [1976], *J. Opt. Soc. Am.* **66**, 376.
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Ge xxiv

Configuration	Term	J	Level (cm ⁻¹)	Leading percentages	
$2s^2 2p^5$	$^2P^o$	$3/2$	0		
		$1/2$	263 500		
$2s\,2p^6$	2S	$1/2$	1 517 440		
$2s^2 2p^4(^3P)3s$	4P	$5/2$	10 468 000	83	17 $2s^2 2p^4(^1D)3s \ ^2D$
		$1/2$	10 594 000	51	38 $2s^2 2p^4(^1S)3s \ ^2S$
		$3/2$	10 707 000	80	19 $2s^2 2p^4(^3P)3s \ ^2P$
$2s^2 2p^4(^3P)3s$	2P	$3/2$	10 498 000	60	21 $2s^2 2p^4(^1D)3s \ ^2D$
		$1/2$	10 732 000	73	25 $2s^2 2p^4(^3P)3s \ ^4P$
$2s^2 2p^4(^1D)3s$	2D	$5/2$	10 803 000	83	17 $2s^2 2p^4(^3P)3s \ ^4P$
		$3/2$	10 811 000	78	21 $2s^2 2p^4(^3P)3s \ ^2P$
$2s^2 2p^4(^1S)3s$	2S	$1/2$	11 133 000	60	24 $2s^2 2p^4(^3P)3s \ ^4P$
$2s^2 2p^4(^3P)3d$	4P	$1/2$	11 357 000	47	29 $2s^2 2p^4(^3P)3d \ ^2P$
$2s^2 2p^4 3d$		$3/2$	11 379 000	37	$2s^2 2p^4(^3P)3d \ ^4P$ 29 $2s^2 2p^4(^3P)3d \ ^2D$
$2s^2 2p^4 3d$		$5/2$	11 406 000	28	$2s^2 2p^4(^3P)3d \ ^2D$ 24 $2s^2 2p^4(^3P)3d \ ^2F$

Ge xxiv - Continued

Configuration	Term	<i>J</i>	Level (cm ⁻¹)	Leading percentages	
$2s^2 2p^4(^3P)3d$	⁴ D	¹ / ₂	11 526 000	71	15 $2s^2 2p^4(^3P)3d$ ² P
$2s^2 2p^4 3d$		³ / ₂	11 557 000	38	$2s^2 2p^4(^3P)3d$ ⁴ D 17 $2s^2 2p^4(^3P)3d$ ⁴ P
$2s^2 2p^4(^3P)3d$	² F	⁵ / ₂	11 570 000	45	24 $2s^2 2p^4(^3P)3d$ ² D
$2s^2 2p^4(^3P)3d$	² P	³ / ₂	11 594 000	41	17 $2s^2 2p^4(^1D)3d$ ² P
$2s^2 2p^4(^1D)3d$	² S	¹ / ₂	11 660 000	72	18 $2s^2 2p^4(^3P)3d$ ⁴ P
$2s^2 2p^4(^1D)3d$	² P	³ / ₂	11 680 000	62	13 $2s^2 2p^4(^3P)3d$ ² P
		¹ / ₂	11 746 000	56	35 "
$2s^2 2p^4 3d$		⁵ / ₂	11 682 000	38	$2s^2 2p^4(^1D)3d$ ² F 27 $2s^2 2p^4(^1D)3d$ ² D
$2s^2 2p^4(^1D)3d$	² D	³ / ₂	11 727 000	60	28 $2s^2 2p^4(^1D)3d$ ² D
$2s 2p^5(^3P^o)3p$	⁴ D	⁵ / ₂	12 177 000	48	32 $2s 2p^5(^3P^o)3p$ ² D
		³ / ₂	12 347 000	42	25 $2s 2p^5(^3P^o)3p$ ⁴ P
$2s 2p^5(^3P^o)3p$	⁴ P	⁵ / ₂	12 231 000	52	47 $2s 2p^5(^3P^o)3p$ ² D
$2s 2p^5(^3P^o)3p$	² P	³ / ₂	12 291 000	47	37 $2s 2p^5(^3P^o)3p$ ⁴ D
$2s 2p^5 3p$		¹ / ₂	12 330 000	38	$2s 2p^5(^3P^o)3p$ ² S 35 $2s 2p^5(^3P^o)3p$ ² P
$2s 2p^5 3p$		¹ / ₂	12 483 000	38	$2s 2p^5(^3P^o)3p$ ² S 35 $2s 2p^5(^3P^o)3p$ ² P
$2s 2p^5(^1P^o)3p$	² D	³ / ₂	12 658 000	74	12 $2s 2p^5(^1P^o)3p$ ² P
		⁵ / ₂	12 724 000	85	8 $2s 2p^5(^3P^o)3p$ ⁴ D
$2s 2p^5(^1P^o)3p$	² P	¹ / ₂	12 721 000	78	10 $2s 2p^5(^3P^o)3p$ ² S
		³ / ₂	12 745 000	76	11 $2s 2p^5(^1P^o)3p$ ² D
$2s^2 2p^4(^3P)4s$	² P	³ / ₂	14 198 000	66	10 $2s^2 2p^4(^1D)4s$ ² D
$2s^2 2p^4(^1D)4s$	² D	⁵ / ₂	14 522 000	82	18 $2s^2 2p^4(^3P)4s$ ⁴ P
		³ / ₂	14 529 000	80	17 $2s^2 2p^4(^3P)4s$ ² P
Ge xxv (³ P ₂)	Limit		18 580 000		

Ge xxv

Z = 32

O I isoelectronic sequence

Ground state $1s^2 2s^2 2p^4 \ ^3P_2$ Ionization energy $19\,700\,000 \pm 200\,000 \text{ cm}^{-1}$ ($2442 \pm 25 \text{ eV}$)

Two magnetic dipole lines were identified by Denne and Hinnov [1987] in a tokamak plasma: the $3p^4 \ ^3P_2 - ^3P_1$ and $^3P_1 - ^1D_2$ at $410.7 \pm 0.3 \text{ \AA}$ and $297.5 \pm 0.3 \text{ \AA}$, respectively.

Behring *et al.* [1976] gave classifications for six lines of the $2s^2 2p^4 - 2s 2p^5$ array observed in a laser plasma. This was extended to 14 lines by Feldman *et al.* [1989], who also identified two transitions from the $2p^6 \ ^1S_0$ level. These authors improved the wavelength uncertainty to $\pm 0.01 \text{ \AA}$ by observing axially along a line-focused laser plasma. Their level values are given, rounded off to correspond with their measurement uncertainty. Gordon *et al.* [1980] have given classifications for transitions in the arrays $2p^4 - 2p^3 3d$ and $2p^4 - 2p^3 3s$, but most are unresolved.

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

References

- Behring, W. E., Cohen, L., Doschek, G. A., and Feldman, U. [1976], *J. Opt. Soc. Am.* **66**, 376.
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Ge xxv

Configuration	Term	J	Level (cm ⁻¹)
$2s^2 2p^4$	3P	2	0
		0	127 240
		1	243 540
$2s^2 2p^4$	1D	2	336 200
$2s^2 2p^4$	1S	0	646 650
$2s 2p^5$	$^3P^\circ$	2	1 324 130
		1	1 457 310
		0	1 596 720
$2s 2p^5$	$^1P^\circ$	1	1 842 920
$2p^6$	1S	0	3 021 850
Ge xxvi ($^4S_{3/2}$)	Limit		19 700 000

Ge xxvi

Z = 32

N I isoelectronic sequence

Ground state $1s^2 2s^2 2p^3 \ ^4S_{3/2}^{\circ}$ Ionization energy $21\,500\,000 \pm 220\,000 \text{ cm}^{-1}$ ($2666 \pm 27 \text{ eV}$)

Two magnetic dipole transitions were identified by Denne and Hinnov [1987] in a tokamak plasma: the $2p^3 \ ^4S_{3/2}^{\circ} - ^2D_{3/2}^{\circ}$ and $2p^3 \ ^4S_{3/2}^{\circ} - ^2D_{5/2}^{\circ}$ lines at $427.9 \pm 0.3 \text{ \AA}$ and $319.1 \pm 0.3 \text{ \AA}$, respectively.

Behring *et al.* [1985] observed nine lines of the arrays $2s^2 2p^3 - 2s 2p^4$ and $2s 2p^4 - 2p^5$. Their results were extended by Feldman *et al.* [1989] to 22 lines with a wavelength uncertainty of $\pm 0.01 \text{ \AA}$ by observing a line-focused laser plasma axially. We give the energy level values of Feldman *et al.* rounded off to correspond with their measurement uncertainty.

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

References

- Behring, W. E., Seely, J. F., Goldsmith, S., Cohen, L., Richardson, M., and Feldman, U. [1985], *J. Opt. Soc. Am. B* **2**, 886.
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 Feldman, U., Ekberg, J. O., Brown, C. M., and Seely, J. F. [1989], *J. Opt. Soc. Am. B* **6**, 1652.

Ge xxvi

Configuration	Term	J	Level (cm^{-1})
$2s^2 2p^3$	$^4S^{\circ}$	$3/2$	0
$2s^2 2p^3$	$^2D^{\circ}$	$3/2$	233 740
		$5/2$	313 520
$2s^2 2p^3$	$^2P^{\circ}$	$1/2$	426 510
		$3/2$	629 730
$2s 2p^4$	4P	$5/2$	1 081 480
		$3/2$	1 248 820
		$1/2$	1 272 280
$2s 2p^4$	2D	$3/2$	1 504 340
		$5/2$	1 569 240
$2s 2p^4$	2S	$1/2$	1 716 690
$2s 2p^4$	2P	$3/2$	1 787 650
		$1/2$	2 038 470
$2p^5$	$^2P^{\circ}$	$3/2$	2 743 260
		$1/2$	3 016 700
Ge xxvii (3P_0)	Limit		21 500 000

Ge xxvii

Z = 32

C I isoelectronic sequence

Ground state $1s^2 2s^2 2p^2 \ ^3P_0$ Ionization energy $22\ 130\ 000 \pm 220\ 000\ \text{cm}^{-1}$ ($2744 \pm 27\ \text{eV}$)

Two magnetic dipole lines were reported by Denne and Hinnov [1987] from tokamak plasma observations: the $2p^2 \ ^3P_0 - ^3P_1$ and $2p^2 \ ^3P_1 - ^3P_2$ at $454.8 \pm 0.3\ \text{\AA}$ and $1473.7 \pm 0.1\ \text{\AA}$, respectively.

Behring *et al.* [1985] identified two lines of the array $2s^2 2p^2 - 2s 2p^3$ in a laser plasma. Their measurement uncertainty is $\pm 0.005\ \text{\AA}$. We derive level values from their wavelength measurements.

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

References

- Behring, W. E., Seely, J. F., Goldsmith, S., Cohen, L., Richardson, M., and Feldman, U. [1985], *J. Opt. Soc. Am. B* **2**, 886.
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Ge xxvii

Configuration	Term	J	Level (cm^{-1})
$2s^2 2p^2$	3P	0	0
		1	219 880
		2	287 736
$2s 2p^3$	$^3P^\circ$	2	1 516 690
$2s 2p^3$	$^3S^\circ$	1	1 633 620
Ge xxviii ($^2P_{1/2}^\circ$)	Limit		22 130 000

Ge xxviii

Z = 32

B I isoelectronic sequence

Ground state $1s^2 2s^2 2p^2 \ ^2P_{1/2}^\circ$ Ionization energy $23\ 280\ 000 \pm 230\ 000\ \text{cm}^{-1}$ ($2886 \pm 29\ \text{eV}$)

The $^2P^\circ$ ground term splitting is determined by the magnetic dipole line at $339.5 \pm 0.3\ \text{\AA}$ given by Denne and Hinnov [1987]. We derive the $^2P_{3/2}^\circ$ level from their wavelength measurement.

The transition array $2s 2p^2 - 2s 2p 3d$ was reported by Burkhalter *et al.* [1989] at $7.3 - 7.8\ \text{\AA}$. No transitions from the $2s 2p^2$ levels to the ground term are known.

We calculated the value for the ionization energy with the Cowan [1981] HFR code and estimate the uncertainty to be 1%.

References

- Burkhalter, P. G., Newman, D. A., and Knauer, J. P. [1989], *J. Opt. Soc. Am. B* **6**, 1964.
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 Denne, B., and Hinnov, E. [1987], *Phys. Scr.* **35**, 811.

Ge xxviii

Configuration	Term	J	Level (cm^{-1})
$2s^2 2p$	$^2P^\circ$	$1/2$	0
		$3/2$	294 550
Ge xxix (1S_0)	Limit		23 280 000

Ge xxix

Z = 32

Be I isoelectronic sequence

Ground state $1s^2 2s^2 \ ^1S_0$ Ionization energy $24\,750\,000 \pm 250\,000 \text{ cm}^{-1}$ ($3069 \pm 31 \text{ eV}$)

Denne and Hinnov measured the M1 transition $2s2p \ ^3P_1^o - ^3P_2^o$ at $408.7 \pm 0.3 \text{ \AA}$. In addition they reported the $2s^2 \ ^1S_0 - 2s2p \ ^3P_1^o$ and $^1P_1^o$ lines at $199.36 \pm 0.03 \text{ \AA}$ and $92.90 \pm 0.03 \text{ \AA}$, respectively. The levels are derived from these identifications.

Seven lines of the $2s2p - 2s3d$ and $2p^2 - 2p3d$ arrays are reported by Burkhalter *et al.* [1989]. Four of these lines are also classified in the next lower ionization stage and two lines are multiply classified in B-like germanium. We omit these results.

Boiko *et al.* [1977] have given classifications for transitions in the $2p^2 - 2p3d$ array, but these are mostly unresolved.

The value for the ionization energy was calculated with the Cowan [1981] HFR code.

Ge xxix

Configuration	Term	J	Level (cm ⁻¹)
$2s^2$	1S	0	0
$2s2p$	$^3P^o$	0	
		1	501 600
		2	746 280
$2s2p$	$^1P^o$	1	1 076 400
.....			
Ge xxx ($^2S_{1/2}$)	<i>Limit</i>		24 750 000

References

- Boiko, V. A., Pikuz, S. A., Safronova, U. I., and Faenov, A. Ya. [1977], *J. Phys. B* 7, 1253.
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 Denne, B., and Hinnov, E. [1987], *Phys. Scr.* 35, 811.

Ge xxx

Z = 32

Li I isoelectronic sequence

Ground state $1s^2 2s^2 S_{1/2}$ Ionization energy $25\,766\,000 \pm 30\,000 \text{ cm}^{-1}$ ($3194.6 \pm 3.7 \text{ eV}$)

Hinnov *et al.* [1989] report wavelengths for the $2s - 2p$ doublet at $200.290 \pm 0.01 \text{ \AA}$ and $122.705 \pm 0.02 \text{ \AA}$ from a tokamak discharge. Behring *et al.* [1989], using a laser-produced plasma, measured the $2s^2 S_{1/2} - 2p^2 P_{3/2}$ transition at $122.738 \pm 0.015 \text{ \AA}$, and the $2s - 3p$, $2p - 3s$, $2p - 3d$, $3p - 4d$ and $3d - 4f$ transitions from $6 - 20 \text{ \AA}$ with a wavelength uncertainty of $\pm 0.015 \text{ \AA}$. We use the measurements of Hinnov *et al.* for the $2s - 2p$ transitions. The uncertainties of the $2p^2 P_{1/2}$ and $2p^2 P_{3/2}$ levels are $\pm 25 \text{ cm}^{-1}$ and $\pm 140 \text{ cm}^{-1}$, respectively. Using the same laser as Behring *et al.*, Burkhalter *et al.* [1989] identified the transitions $2s - 3p$, $2p - 3d$, $2p - 3s$, and one line of the $2p - 4d$ doublet.

We derive the levels of $3s$, $3p$, $3d$, $4d$, and $4f$ from the measurements of Behring *et al.* with an uncertainty of $\pm 30\,000 \text{ cm}^{-1}$.

The value for the ionization energy was determined by Behring *et al.* from a Dirac-Fock calculation of the $4d$ and $4f$ binding energies and their measured excitation energies.

References

- Behring, W. E., Seely, J. F., Brown, C. M., Feldman, U., and Knauer, J. P. [1989], *J. Opt. Soc. Am. B* **6**, 531.
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 Hinnov, E., TFTR Operating Team, Denne, B., and JET Operating Team [1989], *Phys. Rev. A* **40**, 4357.

Ge xxx

Configuration	Term	J	Level (cm ⁻¹)
$1s^2 2s$	2S	$1/2$	0
$1s^2 2p$	$^2P^o$	$1/2$ $3/2$	499 276 814 963
$1s^2 3s$	2S	$1/2$	14 455 000
$1s^2 3p$	$^2P^o$	$1/2$ $3/2$	14 594 000 14 686 000
$1s^2 3d$	2D	$3/2$ $5/2$	14 756 000 14 781 000
$1s^2 4d$	2D	$3/2$ $5/2$	19 560 000 19 570 000
$1s^2 4f$	$^2F^o$	$5/2$ $7/2$	19 594 000 19 595 000
.....			
Ge xxxI (1S_0)	Limit		25 766 000

Ge xxxi

Z = 32

He I isoelectronic sequence

Ground state $1s^2\ ^1S_0$ Ionization energy $109\ 347\ 260 \pm 5000\ \text{cm}^{-1}$ ($13\ 557.34 \pm 0.62\ \text{eV}$)

MacLaren *et al.* [1992] have measured the $1s^2 - 1s2p$ resonance lines with an EBIT device and obtained the following values:

Transition	Wavelength (\AA)
$1s^2\ ^1S_0 - 1s2p\ ^1P_1^o$	$1.205\ 99 \pm 0.000\ 03$
$1s^2\ ^1S_0 - 1s2p\ ^3P_1^o$	$1.212\ 94 \pm 0.000\ 04$
$1s^2\ ^1S_0 - 1s2p\ ^3P_2^o$	$1.208\ 48 \pm 0.000\ 04$
$1s^2\ ^1S_0 - 1s2s\ ^3S_1$	$1.217\ 76 \pm 0.000\ 06$

The $^1S - ^1P^o$ transition energy is 53 ppm above the value calculated by Drake [1988]. We give Drake's separations for the $1s2s$ and $1s2p$ levels and the experimental values for the $1s2s\ ^3S_1$ and $1s2p\ ^3P_1^o$ levels. Drake's estimated uncertainty is $\pm 130\ \text{cm}^{-1}$. We adopt an uncertainty of 5 parts in 10^5 representing the average difference between his values and the best observations in this region of the sequence (see e.g., Beiersdorfer *et al.* [1989] and Drake [1988] for a summary of the measurements). Drake's $n = 2$ levels are $130 \pm 20\ \text{cm}^{-1}$ lower than those circulated privately by him in 1985, which included levels of the $n = 3$ shell. We give the latter reduced by $130\ \text{cm}^{-1}$.

The ionization energy is the value calculated by Drake [1988].

References

- Beiersdorfer, P., Bitter, M., von Goeler, S., and Hill, K. W. [1989], *Phys. Rev. A* **40**, 150.
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 MacLaren, S., Beiersdorfer, P., Vogel, D. A., Knapp, D., Marrs, R.E., Wong, K., and Zasadzinski, R. [1992], *Phys. Rev. A* **45**, 329.

Ge xxxi

Configuration	Term	J	Level (cm^{-1})
$1s^2$	1S	0	0
$1s2s$	3S	1	82 118 000
$1s2p$	$^3P^o$	0	[82 422 810]
		1	82 444 310
		2	[82 751 280]
$1s2s$	1S	0	[82 445 630]
$1s2p$	$^1P^o$	1	[82 923 470]
$1s3s$	3S	1	[97 353 380]
$1s3p$	$^3P^o$	0	[97 437 160]
		1	[97 443 200]
		2	[97 534 920]
$1s3s$	1S	0	[97 440 830]
$1s3d$	3D	2	[97 577 460]
		1	[97 579 170]
		3	[97 610 380]
$1s3p$	$^1P^o$	1	[97 581 770]
$1s3d$	1D	2	[97 613 790]
Ge xxxii ($^2S_{1/2}$)	Limit		[109 347 260]

Ge xxxii

Z = 32

H I isoelectronic sequence

Ground state $1s^2S_{1/2}$ Ionization energy $113\,880\,880 \pm 50 \text{ cm}^{-1}$ ($14\,119.4348 \pm 0.0060 \text{ eV}$)

No spectral line measurements are available for this ion. We give the theoretical value for the binding energy of the $1s$ ground level and the energies of the $2s$ and $2p$ levels relative to this value published by Johnson and Soff [1985]. They estimate the absolute uncertainties to be $\pm 50 \text{ cm}^{-1}$, the uncertainty in the $2p$ $^2P^\circ$ fine structure as $\pm 3 \text{ cm}^{-1}$, and the uncertainty in the $2s - 2p$ interval as $\pm 9 \text{ cm}^{-1}$.

The energies of the levels with $n = 3 - 5$ were obtained by subtracting the binding energies calculated by Erickson [1977] from that of the ground state given by Johnson and Soff. Assuming that the Lamb-shift scales as $(1/n)^3$, we estimate the error in Erickson's calculations for the ns levels as $(8/n)^3$ times the uncertainty of $\pm 472 \text{ cm}^{-1}$ for $2s$. The results for $3s$, $4s$, and $5s$ are $\pm 140 \text{ cm}^{-1}$, $\pm 60 \text{ cm}^{-1}$, and $\pm 30 \text{ cm}^{-1}$. For the remaining levels with $n \geq 3$ we estimate the uncertainties to be $\pm 30 \text{ cm}^{-1}$.

References

- Erickson, G. W. [1977], J. Phys. Chem. Ref. Data 6, 831.
Johnson, W. R., and Soff, G. [1985], At. Data Nucl. Data Tables 33, 405.

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Configuration	Term	J	Level (cm^{-1})
1s	2S	$1/2$	0
2p	$^2P^\circ$	$1/2$	[85 295 210]
		$3/2$	[85 692 518]
2s	2S	$1/2$	[85 304 498]
3p	$^2P^\circ$	$1/2$	[101 220 487]
		$3/2$	[101 338 318]
3s	2S	$1/2$	[101 223 426]
3d	2D	$3/2$	[101 338 105]
		$5/2$	[101 376 377]
4p	$^2P^\circ$	$1/2$	[106 778 029]
		$3/2$	[106 827 682]
4s	2S	$1/2$	[106 779 278]
4d	2D	$3/2$	[106 827 591]
		$5/2$	[106 843 756]
4f	$^2F^\circ$	$5/2$	[106 843 727]
		$7/2$	[106 851 758]
5p	$^2P^\circ$	$1/2$	[109 343 447]
		$3/2$	[109 368 837]
5s	2S	$1/2$	[109 344 085]
5d	2D	$3/2$	[109 368 789]
		$5/2$	[109 377 067]
5f	$^2F^\circ$	$5/2$	[109 377 052]
		$7/2$	[109 381 167]
5g	2G	$7/2$	[109 381 159]
		$9/2$	[109 383 622]
.....			
	Limit		113 880 880